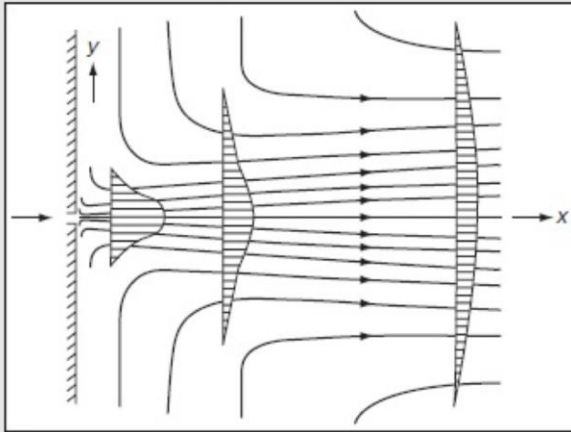


## The planar laminar jet



**FIGURE 4-23**

Definition sketch for a two-dimensional laminar free jet. [After Schlichting (1933a).]

Since jet spreads at  $p = \text{constant}$  w/o bounding walls

$$u_x + v_y = 0$$

$$uu_x + vu_y = \nu u_{yy}$$

$$J = \rho \int_{-\infty}^{\infty} u^2 dy = \text{constant}$$

$$\psi = \nu^{1/2} x^{1/3} f(\eta) = a(x) f(\eta) \quad \eta = \frac{y}{3\nu^{1/2} x^{2/3}} = \frac{y}{b(x)} \quad \text{scaling assumed}$$

$$\eta_y = b^{-1}, \quad \eta_x = -\frac{y}{b^2} b_x = -\frac{\eta b_x}{b}$$

$$u = \psi_y = a f' b^{-1} = \frac{x^{-1/3}}{3} f' \quad \frac{\nu^{1/2} x^{1/3}}{3\nu^{1/2} x^{2/3}} = \frac{x^{-1/3}}{3}$$

$$v = -\psi_x = -a_x f + a f' \frac{\eta b_x}{b}$$

$$v = -\frac{v^{1/2}}{3}x^{-2/3}f + \frac{x^{-1/3}}{3}(2v^{1/2}x^{-1/3})\eta f'$$

$$v = -\frac{1}{3}v^{1/2}x^{-2/3}[f - 2\eta f']$$

$$a_x = \frac{v^{1/2}}{3}x^{-2/3} \quad b = 3v^{1/2}x^{2/3} \quad b_x = \frac{2}{3}3v^{1/2}x^{-1/3} = 2v^{1/2}x^{-1/3}$$

$$(a_x b - a b_x)f'^2 - a_x b f f'' - v f''' = 0$$

$$a_x b = \frac{v^{1/2}}{3}x^{-2/3} \cdot 3v^{1/2}x^{2/3} = v$$

$$a b_x = v^{1/2}x^{1/3} \cdot 2v^{1/2}x^{-1/3} = 2v$$

$$\therefore f''' + \underbrace{f f'' + f'^2}_{\frac{d}{d\eta}(f'f)} = 0$$

$$\text{BC: } u_y = 0, \quad v = 0, \quad y = 0 \quad u = 0, \quad y = \pm\infty$$

$$\therefore f''(0) = 0, \quad f(0) = 0 \quad y = 0$$

$$f'(\pm\infty) = 0$$

$$f'' + f f' = c_1 = 0 \quad f' = f'' = 0 \quad \eta = \pm\infty$$

$$\text{let } \beta = \alpha\eta, \quad f = 2\alpha F(\beta)$$

$$\beta' = \alpha \quad f' = 2\alpha F' \alpha = 2\alpha^2 F'$$

$$f'' = 2\alpha^3 F''$$

$$2\alpha^3 F'' + 2\alpha^2 F' 2\alpha F = 0$$

$$F'' + 2FF' = 0 \quad \frac{d}{d\eta}(F^2) = 2FF'$$

$$\beta = 0, \quad F = 0, \quad \beta = \pm\infty, \quad F' = 0$$

Also  $f'(0) = 1$       $u = \frac{x^{-1/3}}{3} f'$

$u(0) = u_{max} f'(0) \quad \therefore f'(0) = 1$  and  $F'(0) = 1$

$F' + F^2 = c_2 = 1$  Riccati type

$$\beta = \int_0^F \frac{dF}{1-F^2} = \frac{1}{2} \ln \left( \frac{1+F}{1-F} \right) = \tanh^{-1} F$$

or  $F = \tanh \beta = \frac{1-\exp(-2\beta)}{1+\exp(-2\beta)}$

$$\frac{dF}{d\beta} = 1 - \tanh^2 \beta$$

$$u = \frac{2}{3} \alpha^2 x^{-1/3} (1 - \tanh^2 \beta) = \frac{2}{3} \alpha^2 x^{-1/3} \operatorname{sech}^2 \beta$$

$J = \rho \int_{-\infty}^{\infty} u^2 dy = \text{constant}$

$dy = b(x) d\eta = 3v^{1/2} x^{2/3} d\eta$

$$J = \rho \int_{-\infty}^{\infty} \left( \frac{2\alpha^2}{3x^{1/3}} \operatorname{sech}^2(\alpha\eta) \right)^2 3v^{1/2} x^{2/3} d\eta$$

$$= \frac{16}{9} \rho v^{1/2} \alpha^3$$

i.e.,  $\alpha = \left( \frac{9J}{16\sqrt{\rho\mu}} \right)^{1/3} \approx 0.8255 \frac{J^{1/3}}{(\rho\mu)^{1/6}}$

$$u_{max} = u(0) = \frac{2\alpha^2}{3} x^{-1/3} = \frac{2}{3} \left( \frac{9J}{16\sqrt{\rho\mu}} \right)^{2/3} x^{-1/3}$$

$$= \frac{2}{3} \left( \frac{9}{16} \right)^{2/3} \frac{J^{2/3}}{(\rho\mu x)^{1/3}} \approx 0.4544 \left( \frac{J^2}{\rho\mu x} \right)^{1/3} \propto x^{-1/3}$$

$$u = u_{max} \operatorname{sech}^2(\alpha\eta) = u_{max} \operatorname{sech}^2 \left[ \underbrace{0.2752 \left( \frac{\rho J}{\mu^2 x^2} \right)^{1/3}}_{\alpha\eta} y \right]$$

note  $\operatorname{sech}^2 3 \approx 0.01$

**b = width jet** =  $2y$  distance where  $u = 0.01u_{max}$

$$= 2y|_{1\%} \approx 21.8 \left( \frac{\mu^2}{\rho J} \right)^{1/3} x^{2/3} \propto x^{2/3}$$

$$\dot{m} = \rho \int_{-\infty}^{\infty} u \, dy = (36\mu\rho Jx)^{1/3} \approx 3.302 (\mu\rho Jx)^{1/3} \propto x^{1/3}$$

Jet drags ambient fluid along its path as it develops. Solution not valid near  $x = 0$ , as BL approximation not valid since  $Re = \frac{\dot{m}}{\mu} \sim \left( \frac{\rho Jx}{\mu^2} \right)^{1/3}$  per unit depth not large. Profile S shape with inflection points  $\therefore$  unstable transition turbulence  $Re_{crit} \sim 30$ .