

High-Fidelity FSI Simulations and V&V of Vertical and Oblique Flexible Plate Slamming

Riccardo Pellegrini^{1,2}, Matteo Diez², Zhaoyuan Wang³, and Frederick Stern³

¹ Visiting postdoc at IIHR-Hydroscience & Engineering, University of Iowa, Iowa City, IA, USA

² National Research Council, Institute of Marine Engineering (INM), Rome, Italy

³ IIHR-Hydroscience & Engineering, University of Iowa, Iowa City, IA, USA

- Introduction
 - Motivation, background, objectives, and approach
- Physical problem and EFD
- Hydroelasticity parameters
- CFD/CSD/FSI methods and setup
- Overview validation
- Stagnation point model derivation and analysis
- Energy equation derivation and analysis
- MDO for reduced slamming response
- Ongoing research
- Conclusions and future research

- Achievement of weight reduction without loss of structural safety in the design of high-speed vessels
 - Slamming phenomenon must be considered
- High-fidelity computational fluid and structural dynamics (CFD/CSD) are used for load prediction in realistic operating conditions
- Challenges include multi-phase flow and fluid structure interaction (FSI) for hydroelastic phenomena
- Simplified geometries are used to decouple physical phenomena and investigate the underlying physics of the slamming

Search and Rescue, Patrolling



<https://www.youtube.com/watch?v=Sx57-LnuuFs>

Commercial activities



<https://www.youtube.com/watch?v=bZSM5ZbdpWw>

- Analytical, numerical, and experimental investigations on slamming of wedges, plates, and other simplified geometries have been performed in the last century
- Main findings are:
 - Self-similar flow models are a reasonable approximation for 2D investigations
 - Gravity is important for low-speed water entry
 - There exists a high-pressure ridge moving along the wet surface during the slam
 - The velocity of the spray root is not constant
 - Impact loads scale with the normal to the body surface velocity
 - Elastic structures can exhibit:
 - Spray separation
 - Increase of the total load
 - Reduction of the local pressure
 - Cavitation and/or ventilation
 - For elastic plates, the significance of the hydroelasticity is a function of the wetting time and the first eigenfrequency

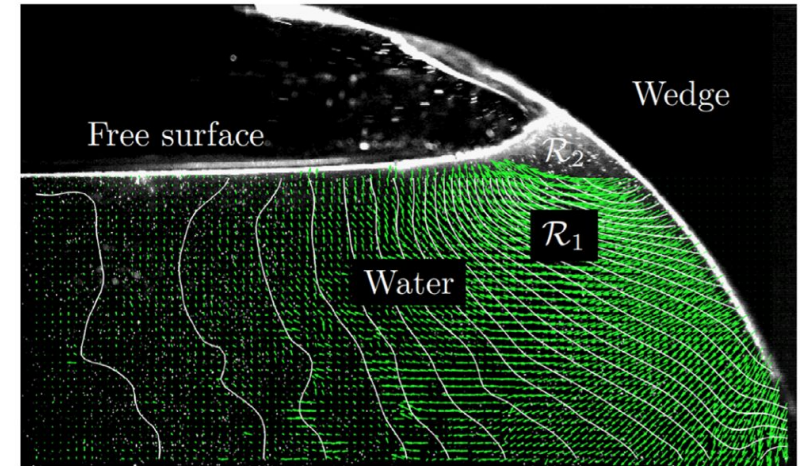
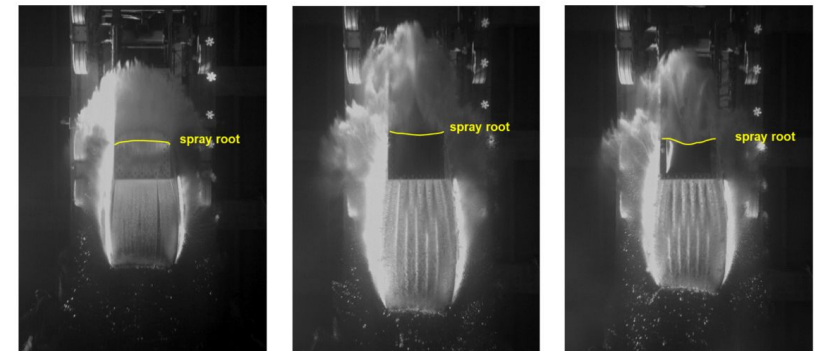


Fig. 4. Example of a high speed image with overlaid PIV measurement and relative contour lines. The wedge considered here has an average deadrise angle $\beta_0 = 35^\circ$ and deadrise angle at the keel of $2\beta_0$.

Pancioli, R., A. Shams, and M. J. O. E. Porfiri. "Experiments on the water entry of curved wedges: High speed imaging and particle image velocimetry." *Ocean Engineering* 94 (2015): 213-222.

Example of water slam with large deflection and air entrapment

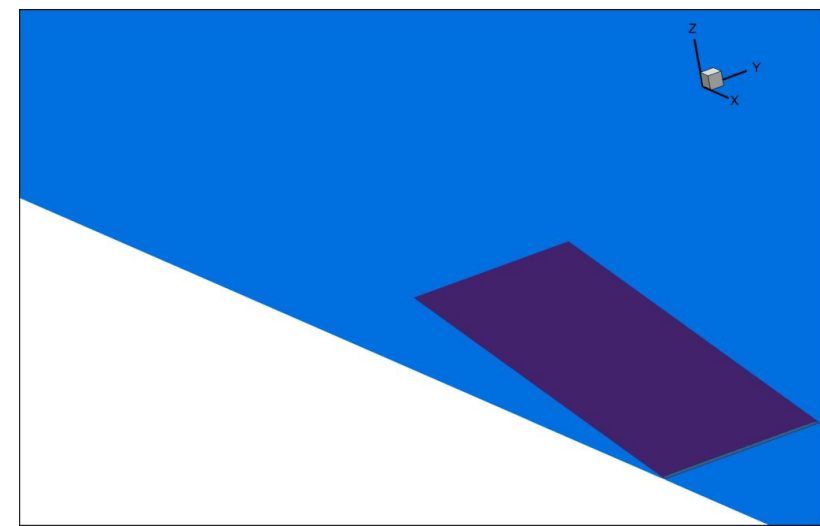


(a) 15.10.30 (thickness 15 mm) (b) 3.10.30 (thickness 3 mm) (c) 08.10.30 (thickness 0.8 mm)

Spinosa, Emanuele, and Alessandro Iafrati. "Experimental investigation of the fluid-structure interaction during the water impact of thin aluminium plates at high horizontal speed." *International Journal of Impact Engineering* 147 (2021): 103673.

Objective

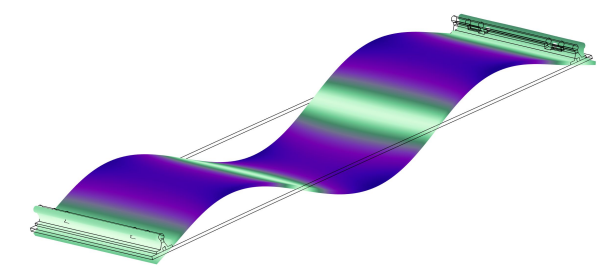
- Development of numerical methods for
 - Resolving complex FSI problem with large impacts and deformations
 - Understanding the physics of complex FSI phenomena
 - Assessing/investigating new advanced concepts, designs, and materials
 - Building capabilities for optimizing/controlling the response of ship structures with the aim of reducing the weight, increasing the structural payload, and personnel safety
- Development of effective analysis methods and approaches
 - Extended Bernoulli equation for spray root dynamic analysis
 - Energy conservation applied to FSI
- Analysis and multidisciplinary optimization of anisotropic structures
- Collaboration with UMD for investigation of elastic flat plate slamming



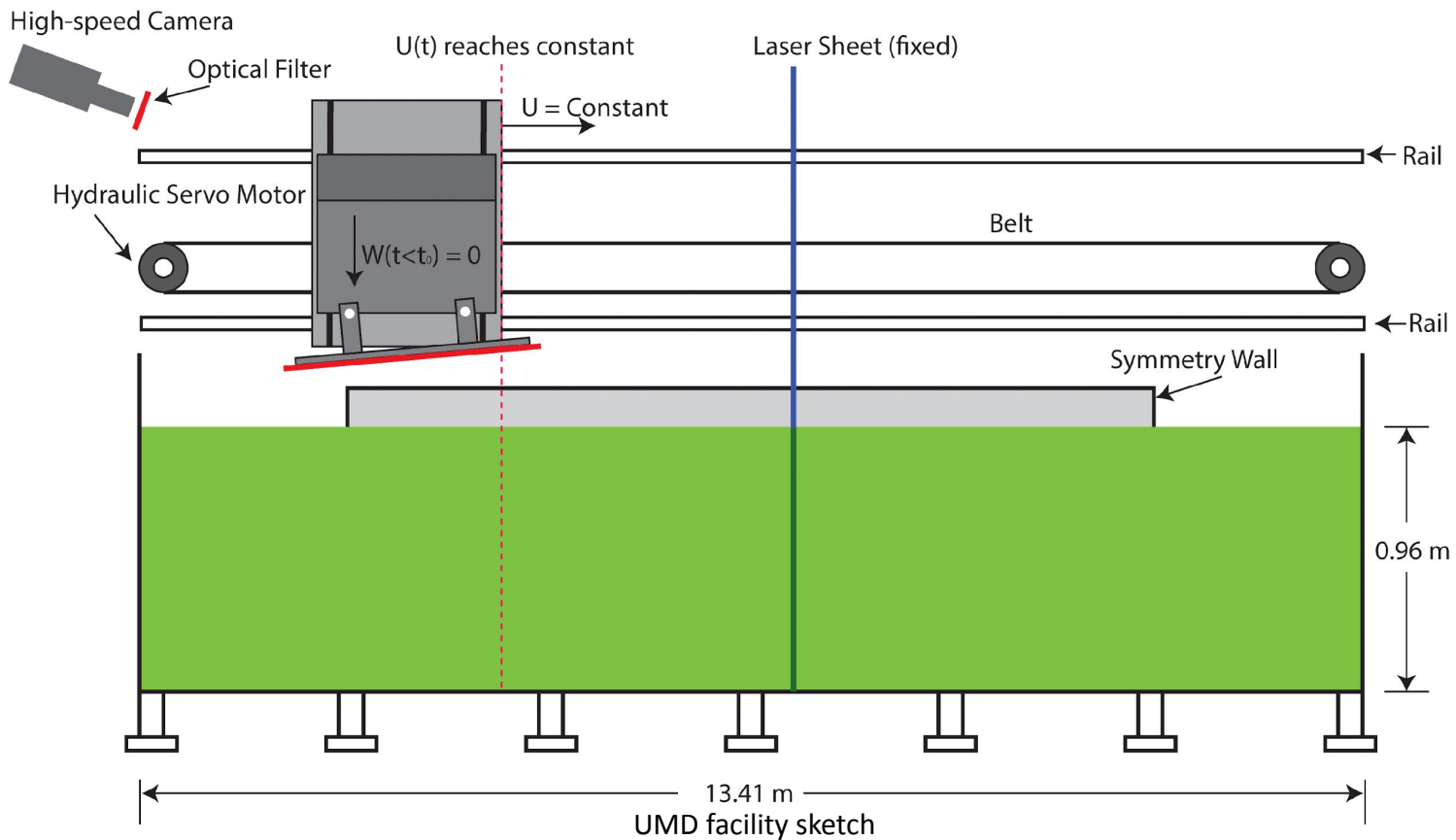
Oblique flat plate slamming

Approach

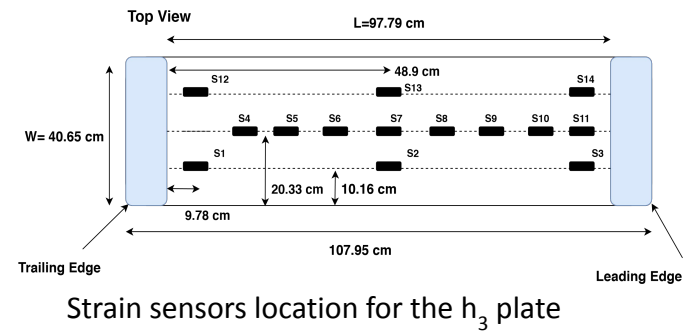
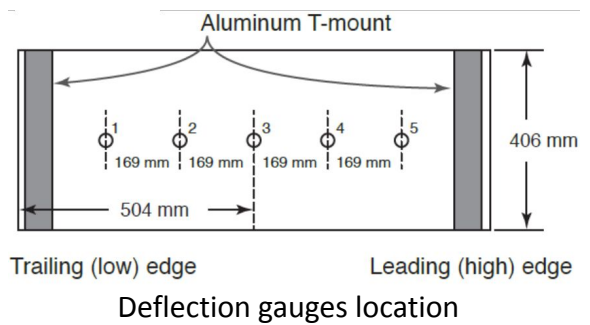
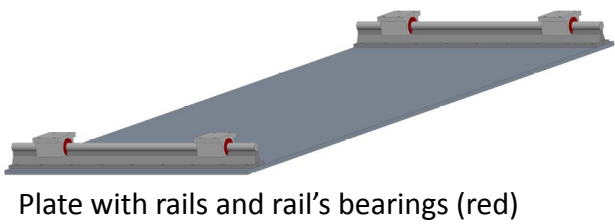
- Identification of a general and effective approach to model geometric nonlinearities
- Lagrangian and Eulerian approaches for the investigation of the physics of slamming
- Machine learning for MDO of anisotropic structures
- Comparison with UMD data from elastic plate slamming experiments



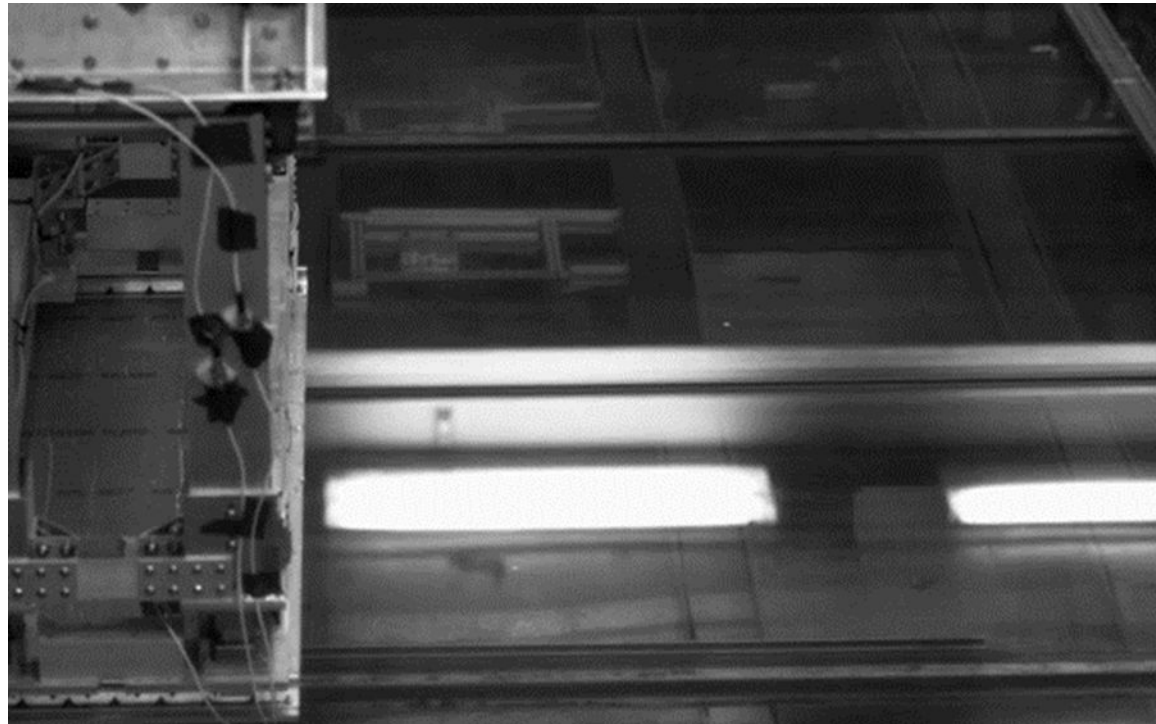
Prestressed modal analysis for nonlinear modal expansion setup



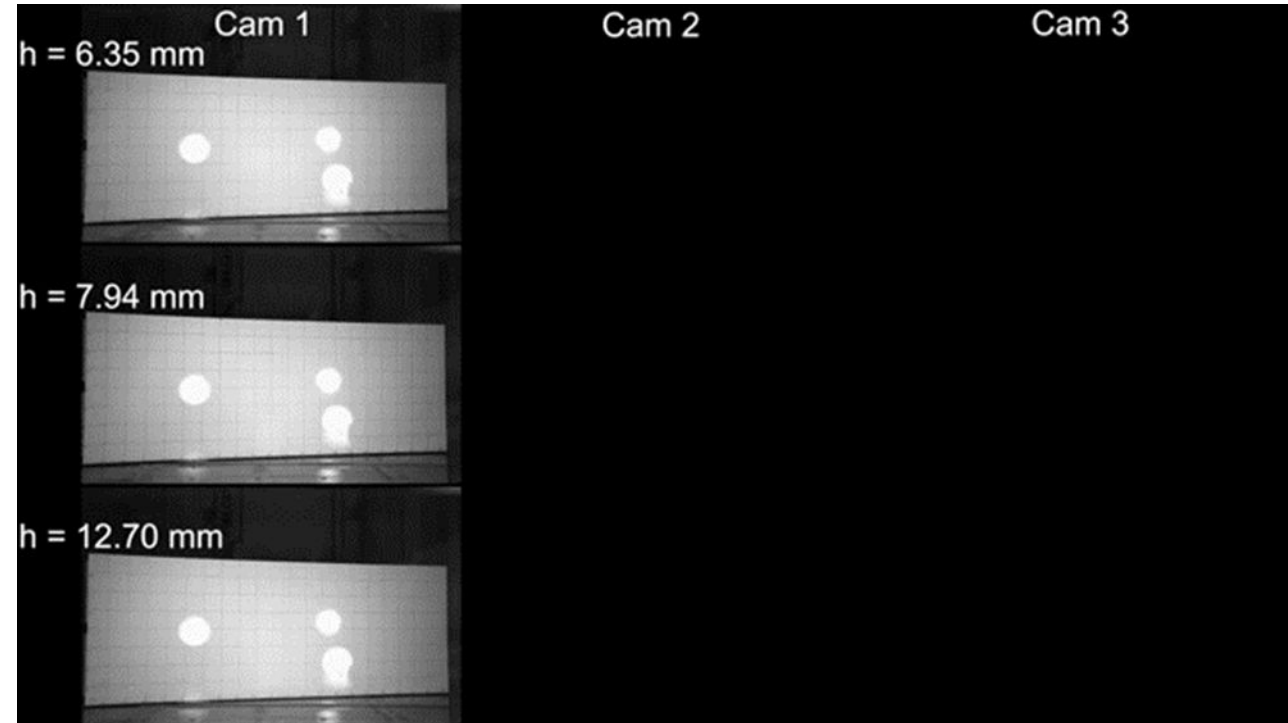
- Validation variables are
 - Normal force and transverse moment
 - Strains
 - Centerplate deflection
 - Pressure
 - Spray root position and velocity



- Videos of the EFD, from the YouTube channel of An Wang from the University of Maryland.
 - $U = 4 \text{ m/s}$, $W = 0.8 \text{ m/s}$
 - Pitch angle = 10 deg

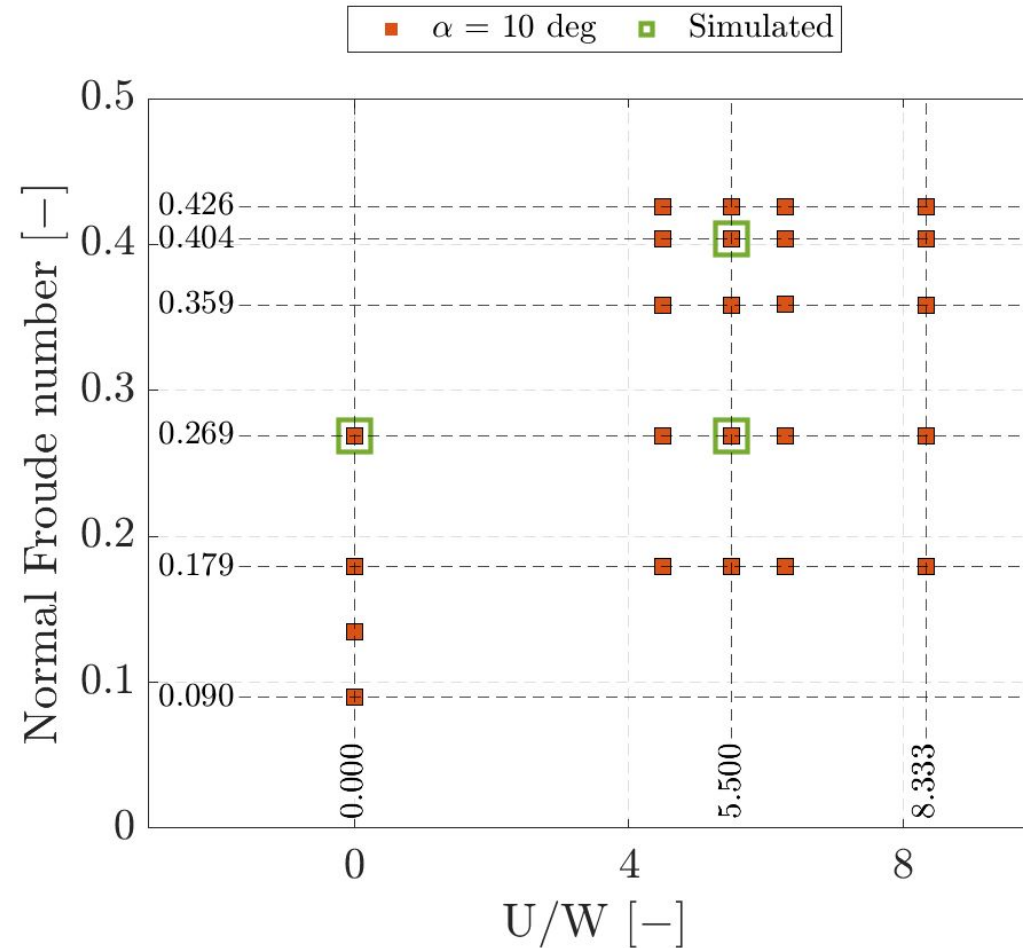


<https://www.youtube.com/watch?v=p7n1u-Hewl>



<https://www.youtube.com/watch?v=eyU-HeDn52Y>

- UMD test matrix covers 24 combinations of vertical and horizontal velocities for 3 plate thicknesses



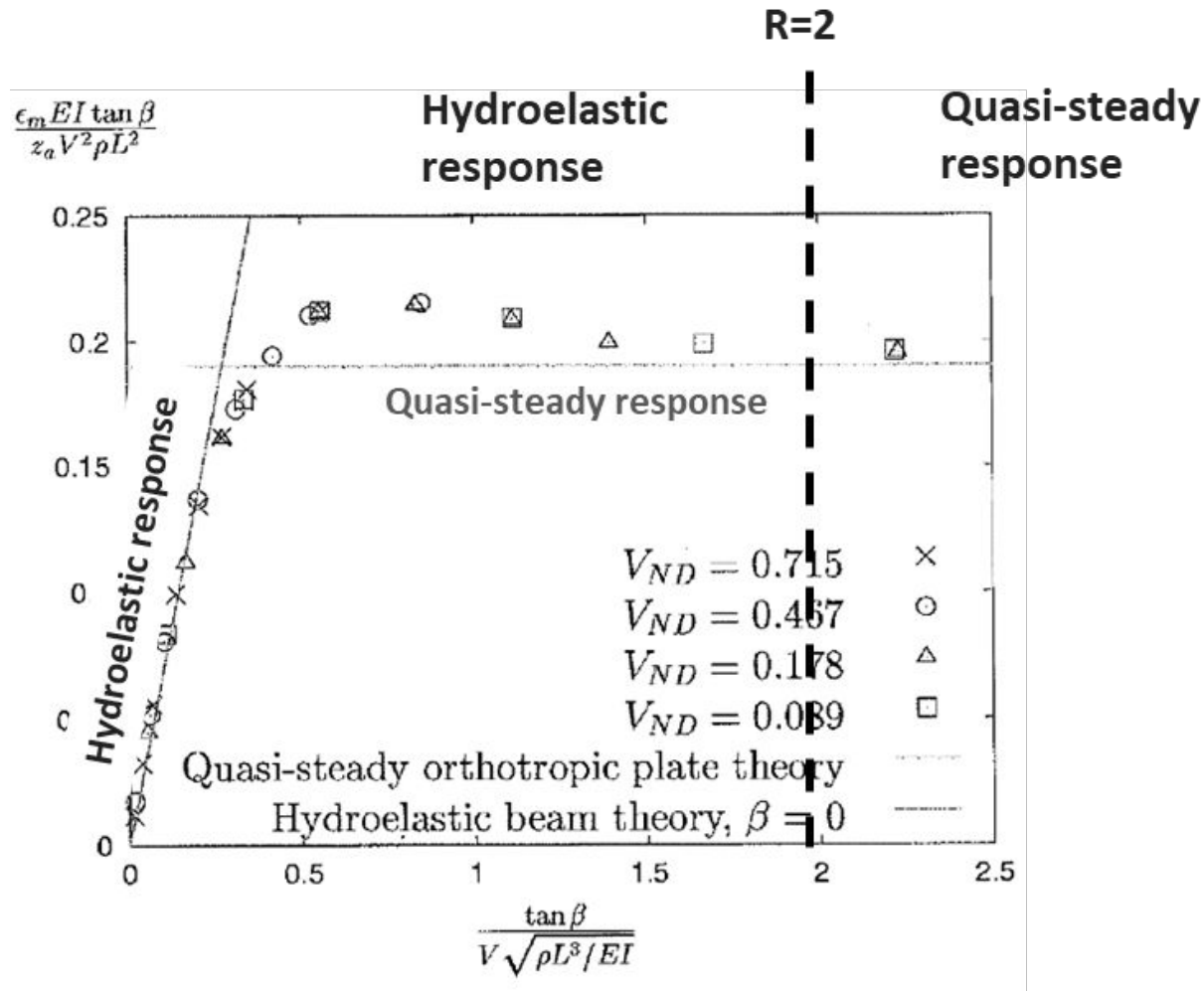
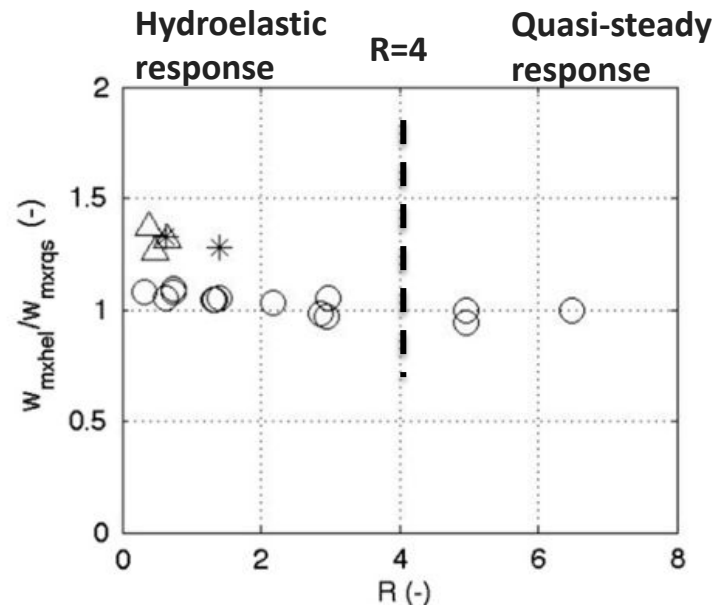


Figure from Faltinsen, O. M. (1999)

- Ongoing studies focus on the scaling of deflections and strains. Faltinsen, 1999 suggested two scaling for the maximum strains
- Quasi-steady response better use v_n^2 : this results from the solution of a boundary value problem for quasi-steady response (Faltinsen, 1999)
- Hydroelastic response better use v_n : this results from the solution of a boundary value problem of a free-vibrating beam slamming on water (Faltinsen 1997, Kvalsvold and Faltinsen 1995)

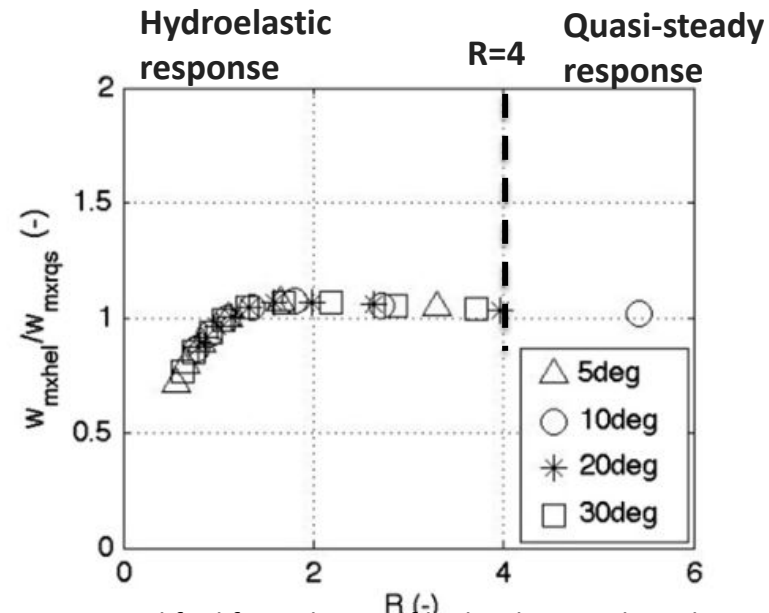
Reference	Slam type	Structure	R definition	R interpretation	R threshold	Scaling
2007 Stenius et al. Explicit FE-modelling of hydroelasticity in panel-water impacts	Vertical	Beam		Ratio between two times the wetting time of the panel and the first natural wet period	R<4 hydroelasticity matters	The hydroelastic deflection and strain are scaled with the equivalent quasi-steady deflection and strains
2010 Stenius et al. Hydroelastic Interaction in Panel-Water Impacts of High-Speed Craft	Vertical	Plate				

- This formulation considers the boundary conditions of the structure
- The range of R is different as the threshold for hydroelastic response
- Deflections are nondimensionalized with respect to the quasi-steady response
- There is a significant difference between simulations and the analytical model
- Many analytical models consider the slam until the spray root exits the structure thus not considering that the deformation of the structure may increase after



Results of numerical simulations

Figures from Stenius et al. 2010

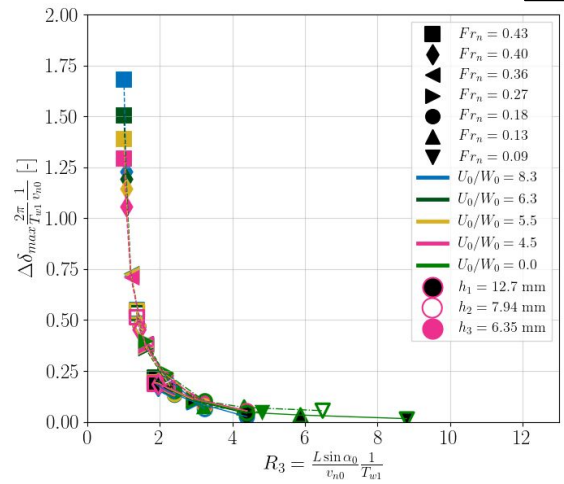


From simplified formulation of hydroelasticity based on engineering beam theory and potential theory
Similar to Faltinsen

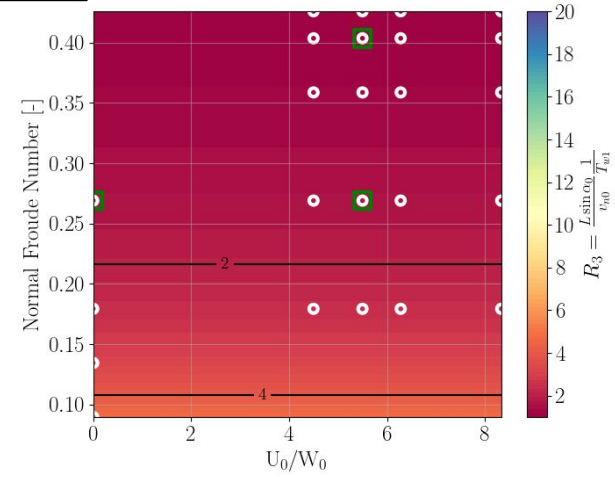
Reference	Slam type	Structure	R definition	R interpretation	R threshold	Scaling
1999 Faltinsen Water Entry of a Wedge	Vertical	Plate reinforced with stiffeners		Ratio between the wetting time of the rigid wedge and the first natural wet period of the stiffener	$R < 2$ hydroelasticity matters	
2001 Berezniński Slamming: the role of hydroelasticity	Vertical	Beam (as the lower part of a wedge)		As per definition	$R < 2$ hydroelasticity matters	The hydroelastic deflection is scaled with the equivalent quasi-steady deflection
2007 Bogaert and Kaminski Hydro-elastic criterion for practical design	Vertical	2D cone (wedge)		As per definition	$R < 2$ hydroelasticity matters	The hydroelastic deflection is scaled with the equivalent quasi-steady deflection
2007 Stenius et al. Explicit FE-modelling of hydroelasticity in panel-water impacts	Vertical	Beam		Ratio between two times the wetting time of the panel and the first natural wet period	$R < 4$ hydroelasticity matters	The hydroelastic deflection and strain are scaled with the equivalent quasi-steady deflection and strains
2010 Stenius et al. Hydroelastic Interaction in Panel-Water Impacts of High-Speed Craft	Vertical	Plate		Ratio between two times the wetting time of the panel and the first natural wet period	$R < 4$ hydroelasticity matters	The hydroelastic deflection and strain are scaled with the equivalent quasi-steady deflection and strains
2015 Panciroli, Porfiri Analysis of hydroelastic slamming through particle image velocimetry	Vertical	Plate		Not discussed	$R < 2$ hydroelasticity matters	Not discussed

- Two definitions for the hydroelasticity parameter are used
 - R_3 is effective in identifying slamming conditions that have similar structural responses in terms of deflection
 - R_4 is effective to distinguish the different slam conditions and supports the choice of selecting only three slamming conditions to completely investigate the validation variables variability as per the experiments
- Deflections scaled with v_n
- Hydroelasticity parameter (R) effectively to scale the maximum deflection (from EFD)
 - R is a reliable metric for the hydroelastic effects
- The V.269 and O.403 are characterized by significant hydroelastic effects
- The O.269 case shows limited hydroelastic effects
- The analysis based on the hydroelasticity parameter shows that the three cases selected for CFD and CFD/FSI numerical investigations are adequate to investigate the EFD trends

$R_3(v_n)$

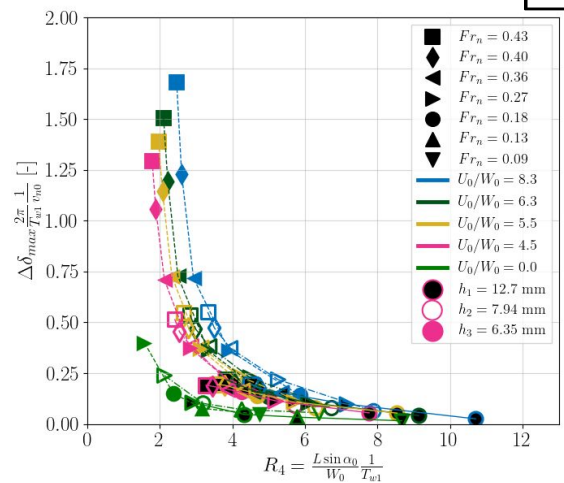


R_3 , hydrodynamic scaling of the differential deflection

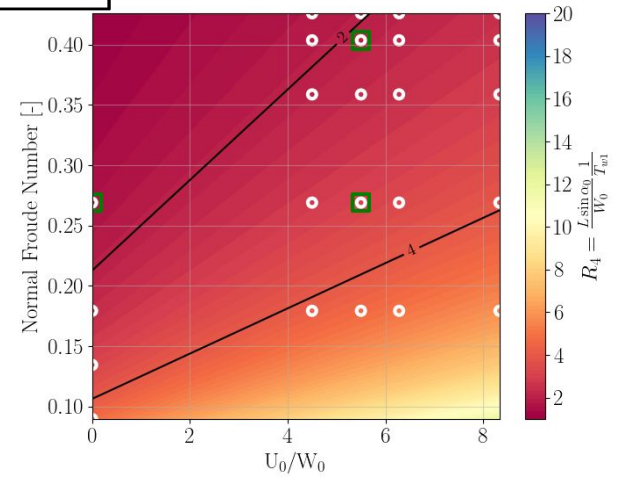


UMD test matrix with contour plot based on R_3 for 6.35 mm thick plate

$R_4(W_0)$

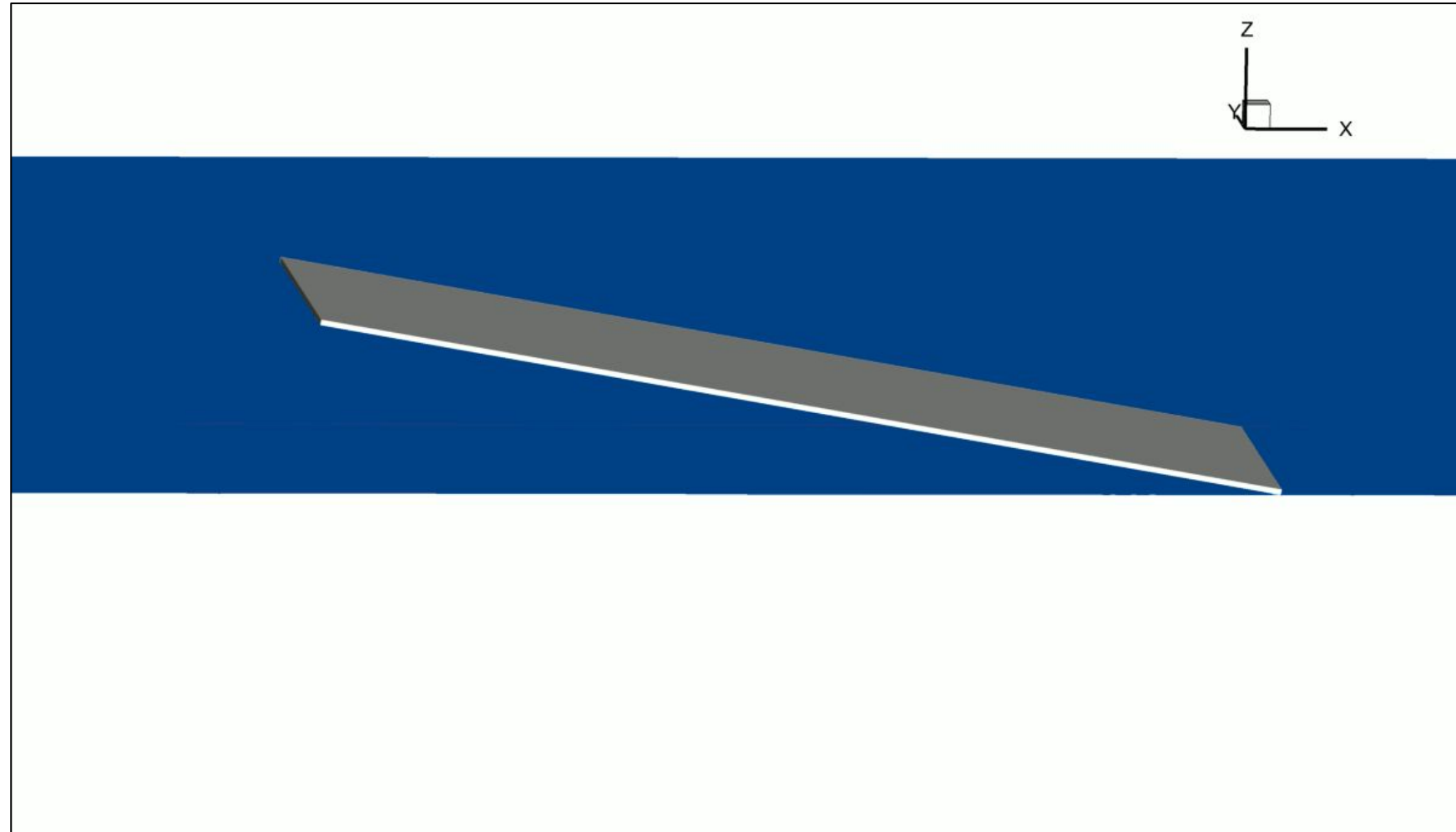
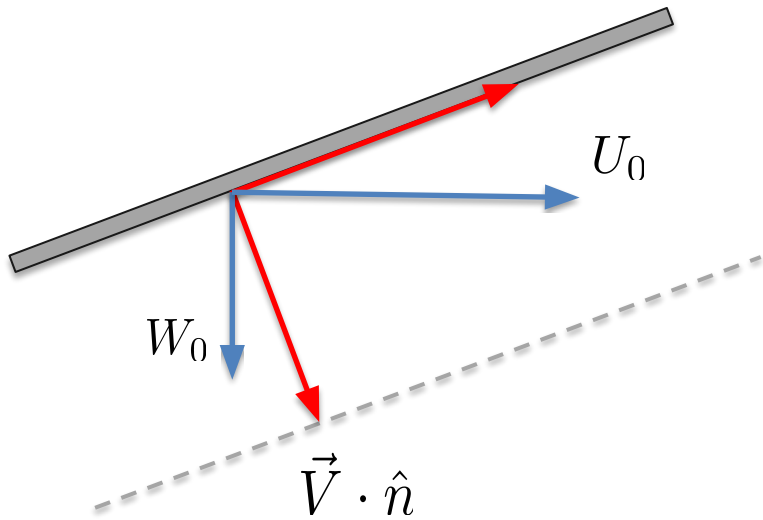


R_4 , hydrodynamic scaling of the differential deflection



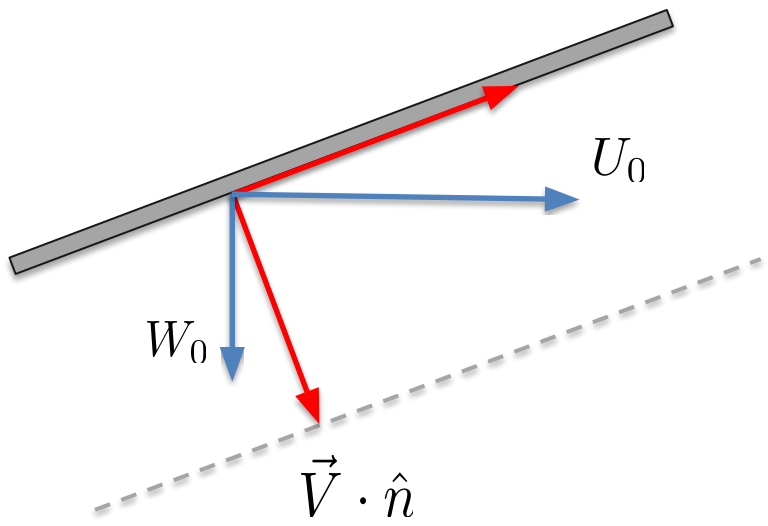
UMD test matrix with contour plot based on R_4 for 6.35 mm thick plate

- This is the reference case.
- All the analysis will be performed in the direction normal to the plate (normal forces, normal velocity, etc...)

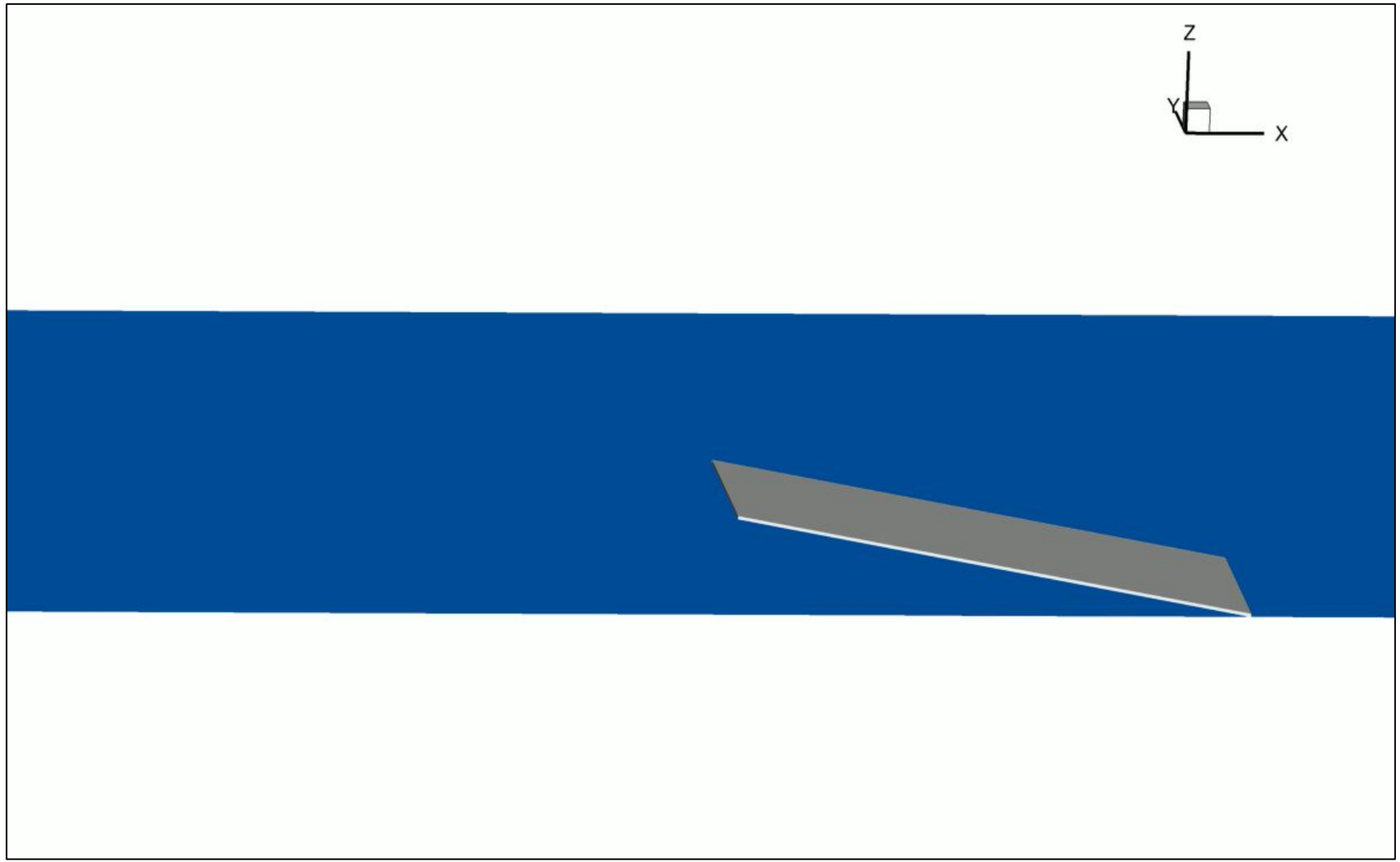


- V.269 – Vertical impact:
- $W = 0.889 \text{ m/s}$
- $U = 0$
- $U/W = 0$
- Normal impact $Fr = 0.269$

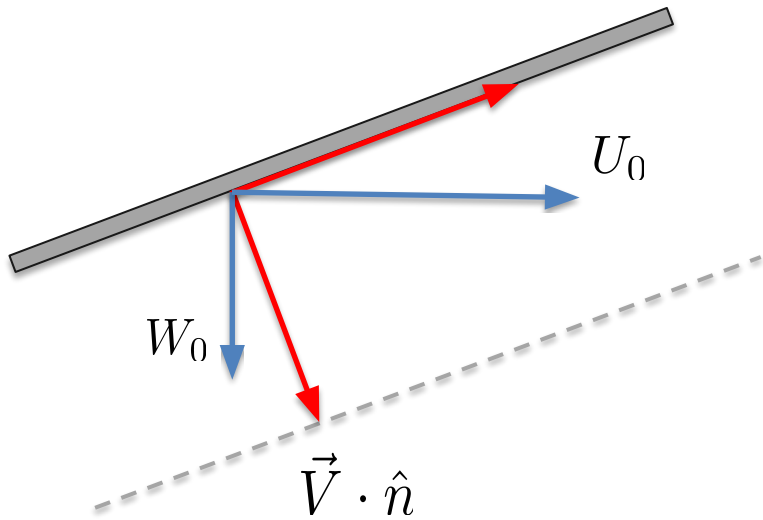
- What happen if a horizontal component of the velocity is considered without changing the normal impact Fr ?



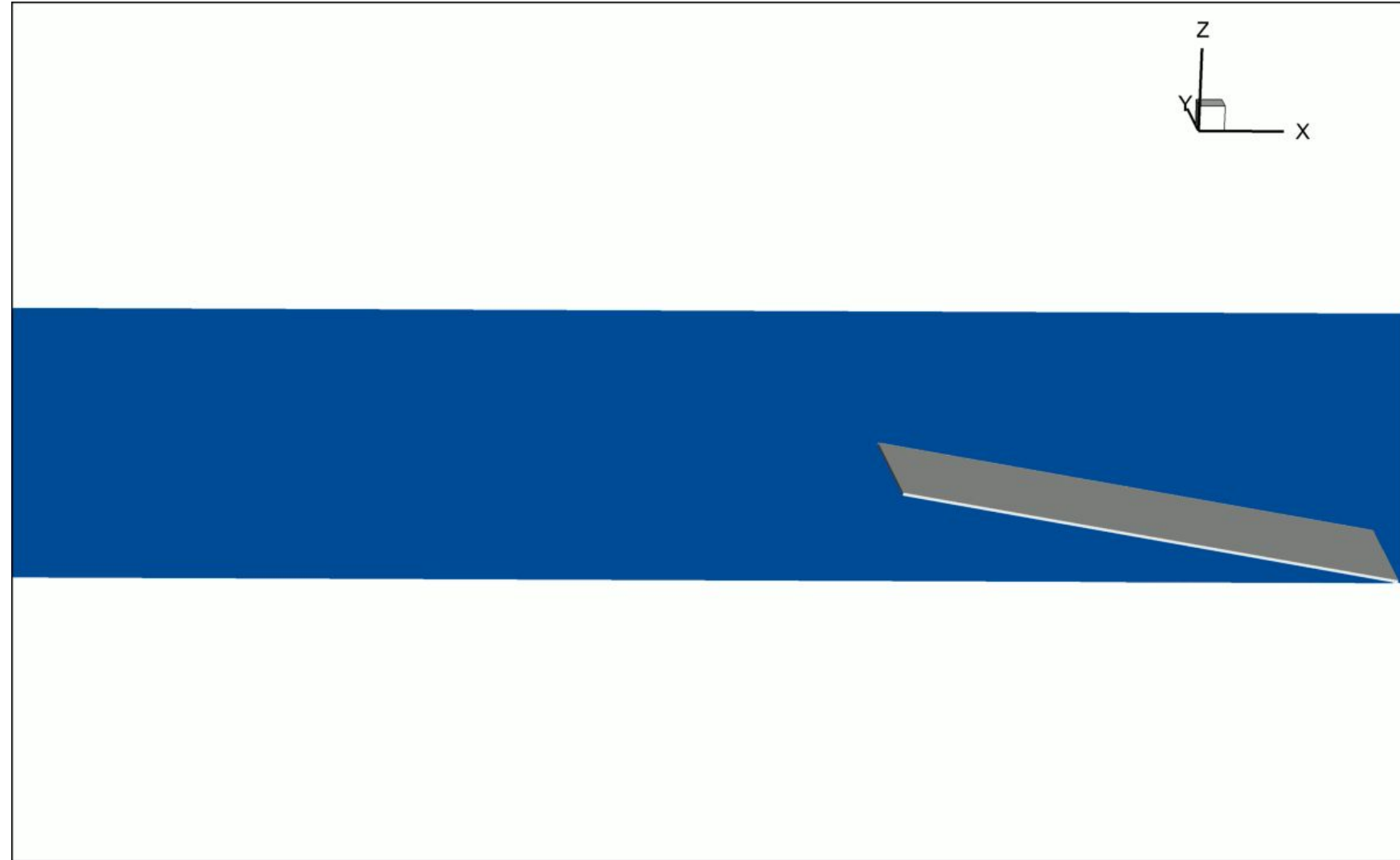
- 0.269 – Slow oblique impact:
- $W=0.451$ m/s
- $U=2.482$ m/s
- $U/W = 5.5$
- Normal impact $Fr = 0.269$



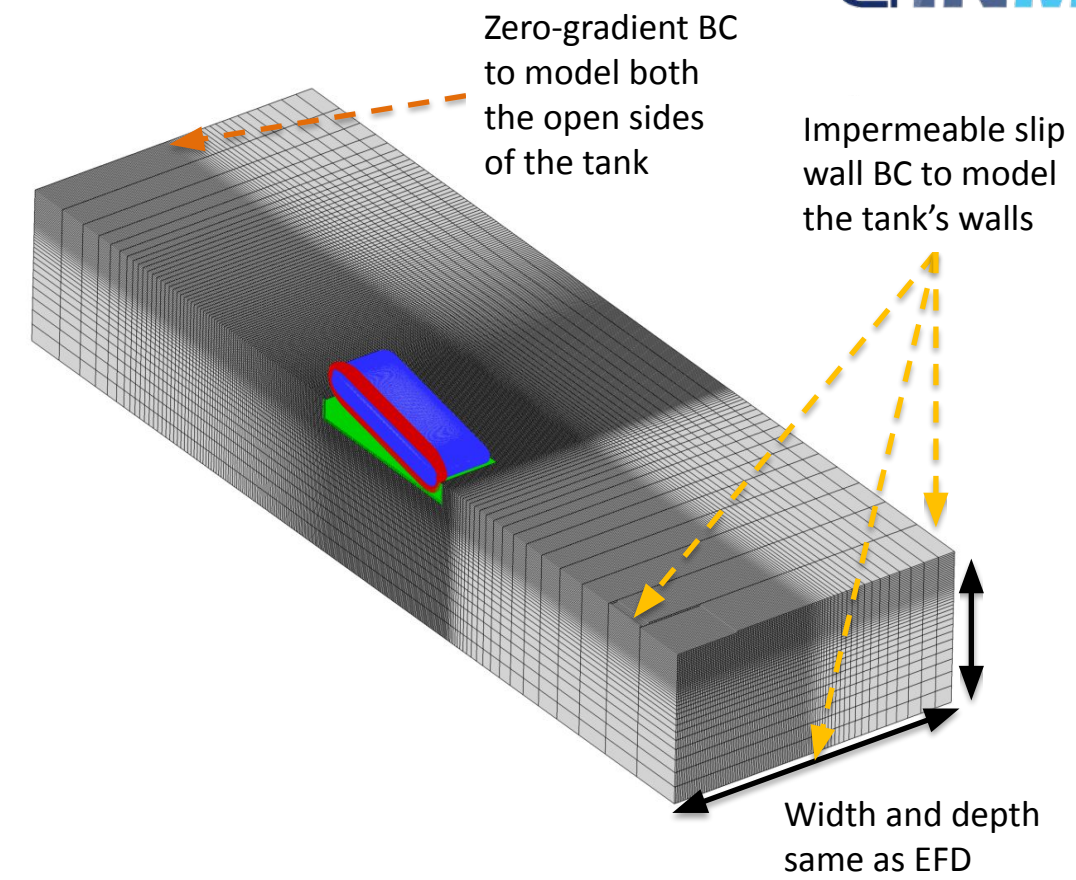
- What happen if the normal impact Fr is increased without changing the ratio of the velocity components?



- 0.403 – Fast oblique impact:
- $W=0.677$ m/s
- $U=3.723$ m/s
- $U/W = 5.5$
- Normal impact Fr = 0.403

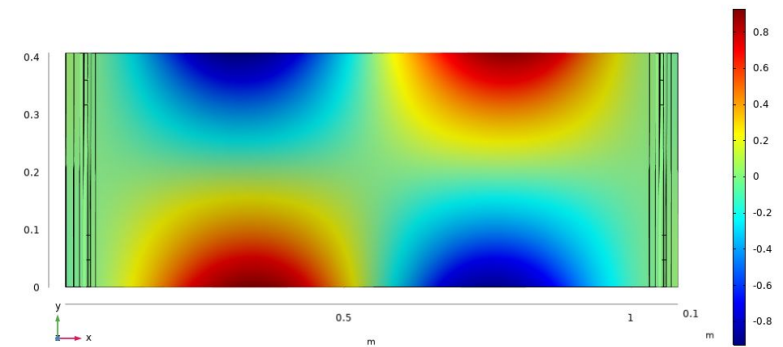
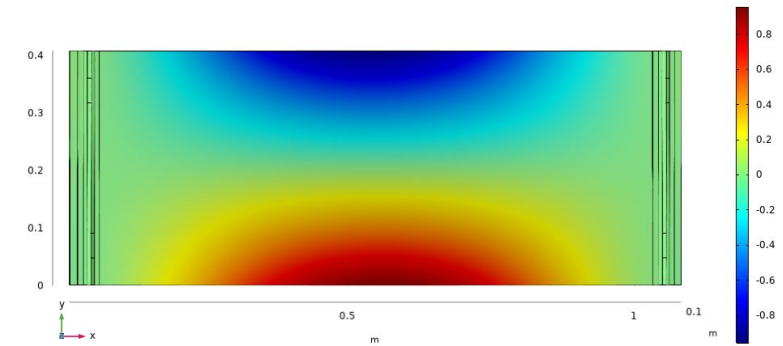
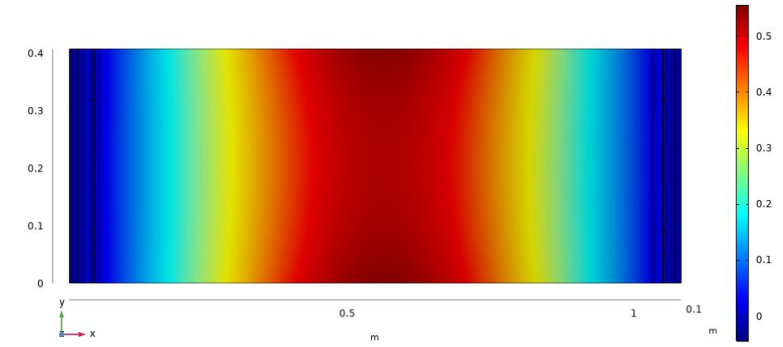
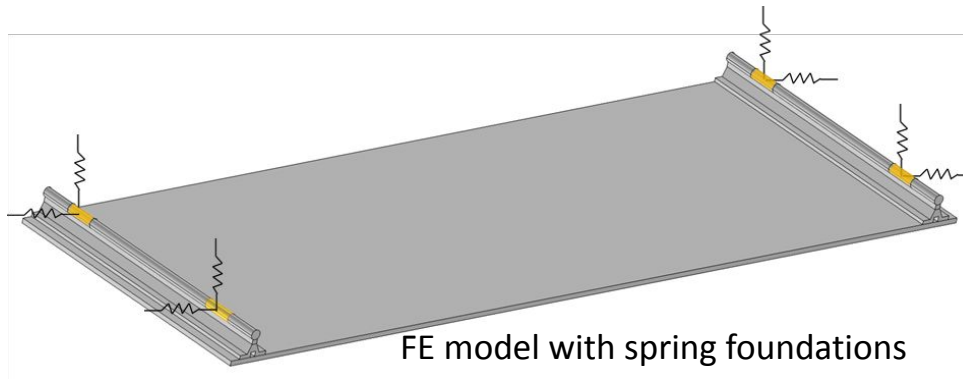


- CFDShip Iowa V4.5
- Simulation conditions selected to investigate the effects of Δh , U/W and Fr_n
- One and two-way FSI simulations are performed using both single and two-phase CFD models
- The 1-way results are compared with those for the thickest plate $h_1 = 12.7$ mm
- The grid is designed to investigate the fluid-structure interaction during the slamming and not yet refined for splash studies
- Width and depth of the domain are the same of the tank used for the experiments, including the gap between the plate and the inner “symmetry” wall
- For all the simulations, the crossing of the plate through the still water level is discretized with 334 time steps.
- FSI simulations are performed using 25 dry natural modes for the modal expansion with Rayleigh damping. The stiffness dumping coefficient is $\beta = 0.0001$

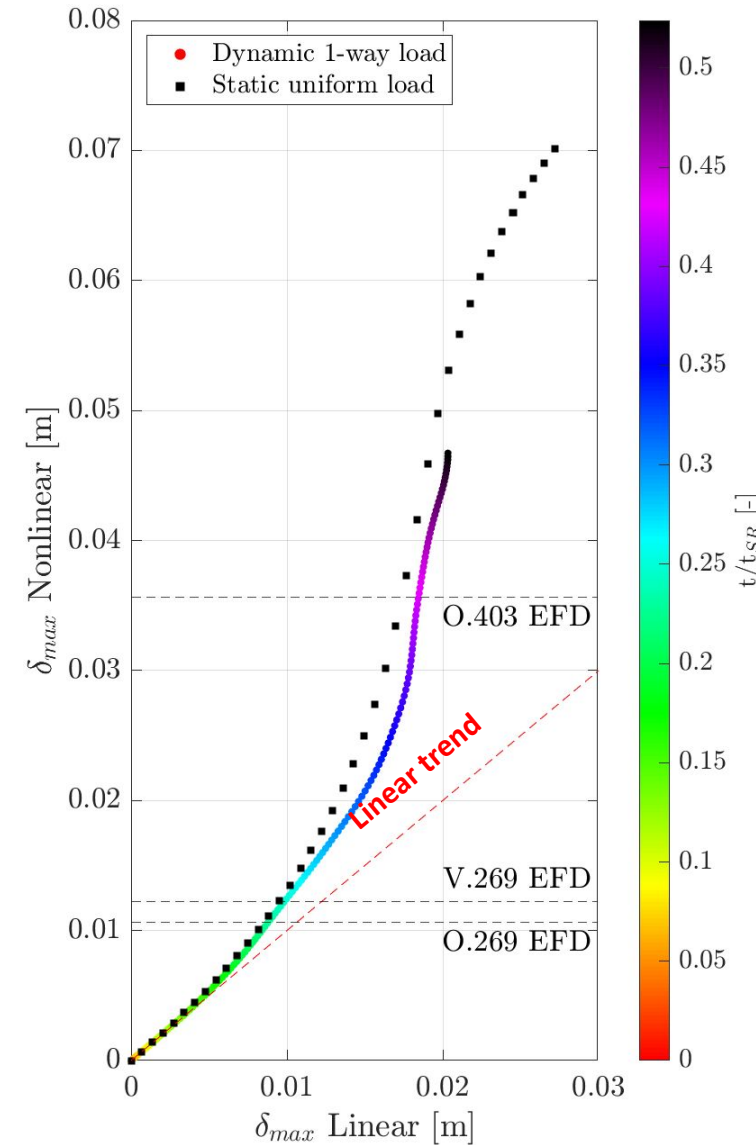


Block	Numbers of grid points [M]
Background (grey)	3.7
Plate's body fitted (blue)	28.4
Gap (red)	1.5
Refinement (green)	2.7

- COMSOL is used to model the plate by FEM
 - Spring foundations are used to calibrate the model
- Elastic constant of the springs is defined to match the first eigenfrequency of the pinned plate



Deflection, nonlinear versus linear evaluation



$$\bar{M}\ddot{\bar{u}} + \bar{K}(\bar{u})\bar{u} = \bar{f}_e$$

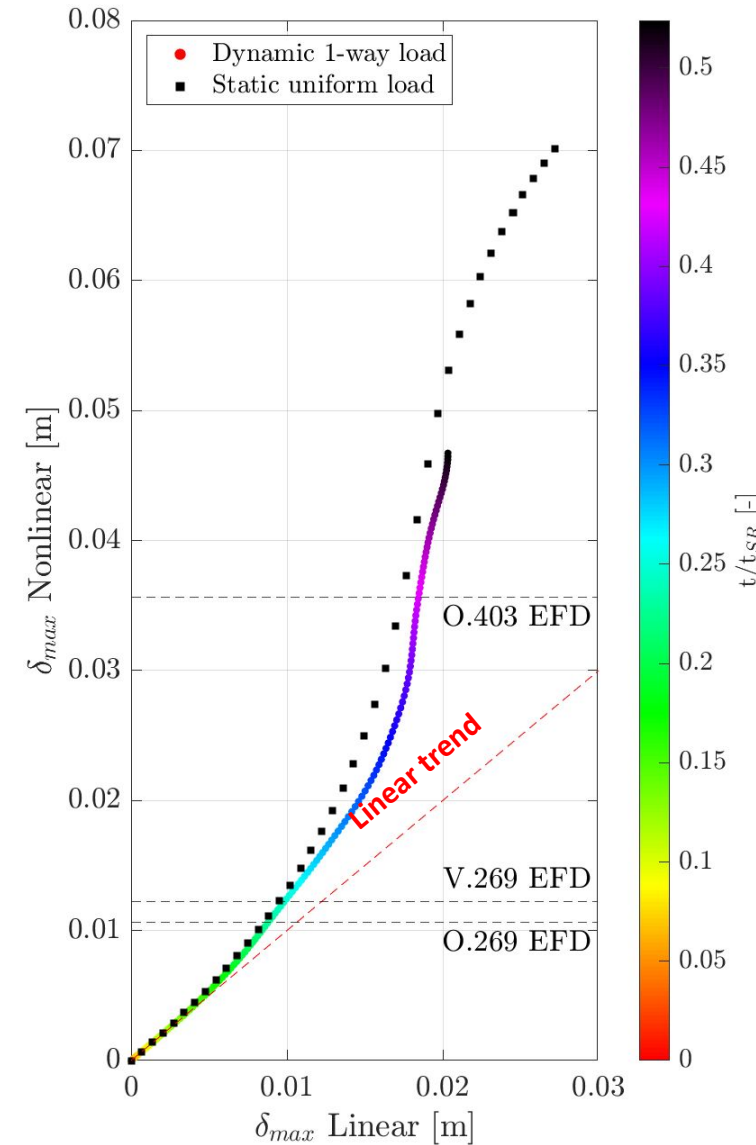
$\bar{K}(\bar{u})$

Modal expansion

$$\Delta\ddot{q}_i + 2\omega_i^2(\bar{q}_i)\xi\Delta\dot{q}_i + \omega_i^2(\bar{q}_i)\Delta q_i = \Delta f_i$$

The stiffness operator is nonlinear since it is dependent from the displacement

Deflection, nonlinear versus linear evaluation



$$\bar{M}\ddot{\bar{u}} + \bar{K}(\bar{u})\bar{u} = \bar{f}_e$$

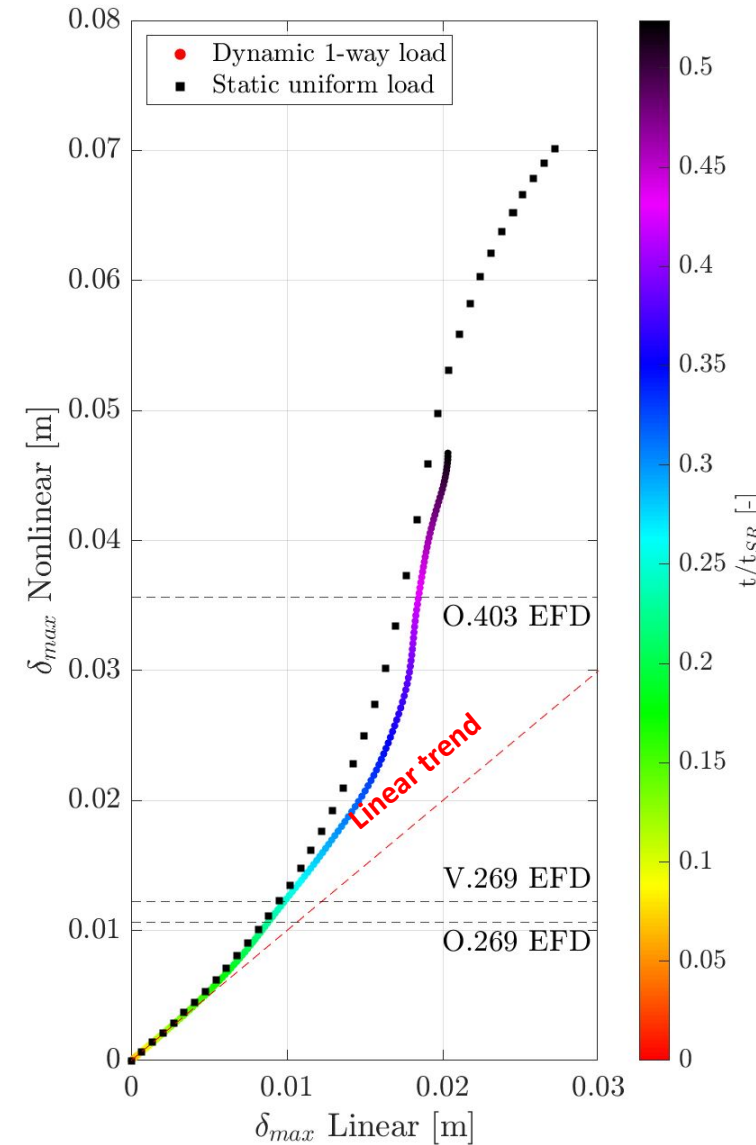
Modal expansion

$$\Delta\ddot{q}_i + 2\omega_i^2(\bar{q}_i)\xi\Delta\dot{q}_i + \omega_i^2(\bar{q}_i)\Delta q_i = \Delta f_i$$

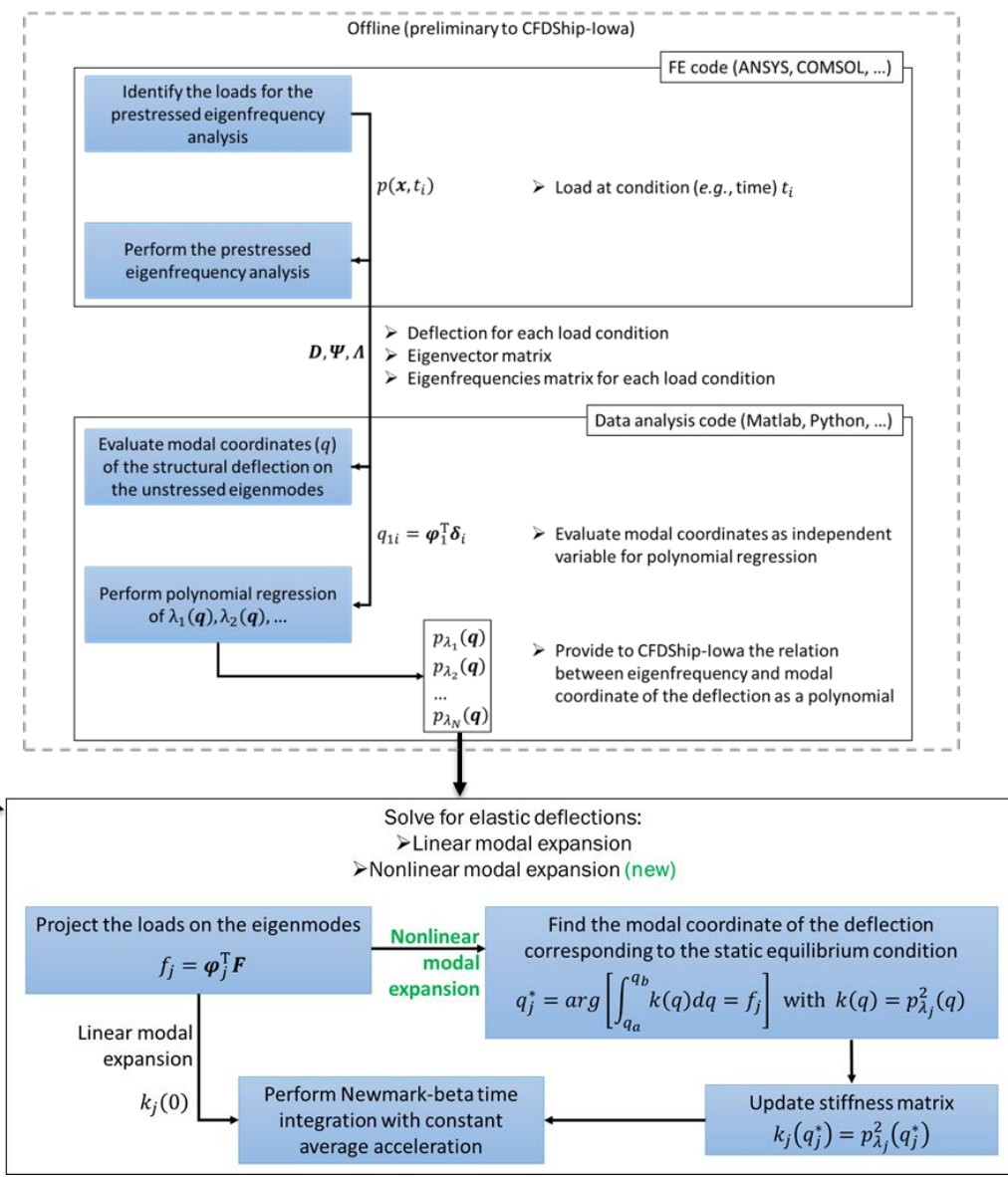
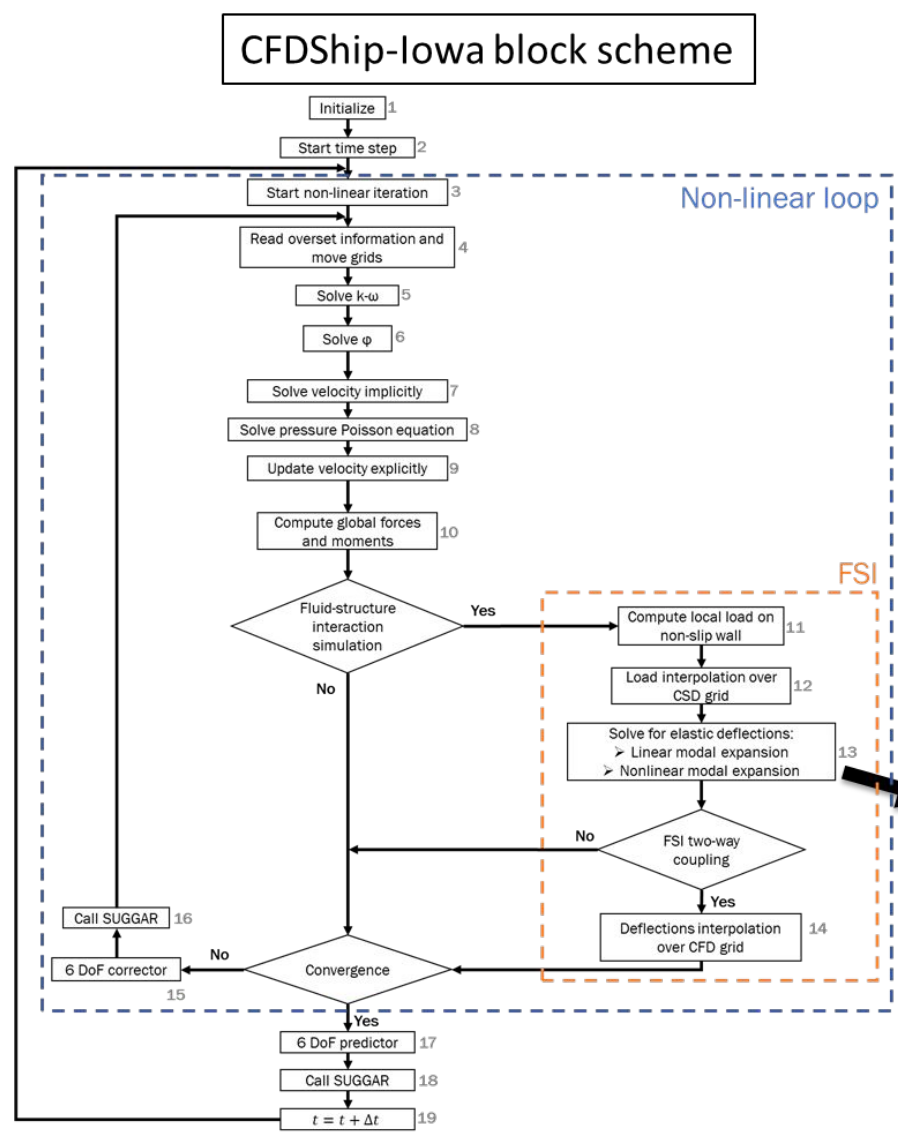
The stiffness operator is nonlinear since it is dependent from the displacement

- Non-intrusive reduced order methods (ROM of discretized PDEs)
 - Any structure
 - Low computational cost
 - Easy to couple with other solvers for multidisciplinary analysis
- CFDShip-lowa now uses a nonlinear modal expansion approach
 - The basis is computed as eigenvectors in vacuum
 - Water added mass effects are provided by direct computation of the fluid
 - Stiffness coefficients are determined by pre-stressed eigenfrequency analysis
 - **A polynomial model of the stiffness coefficient as a function of the modal coordinates is then realized**

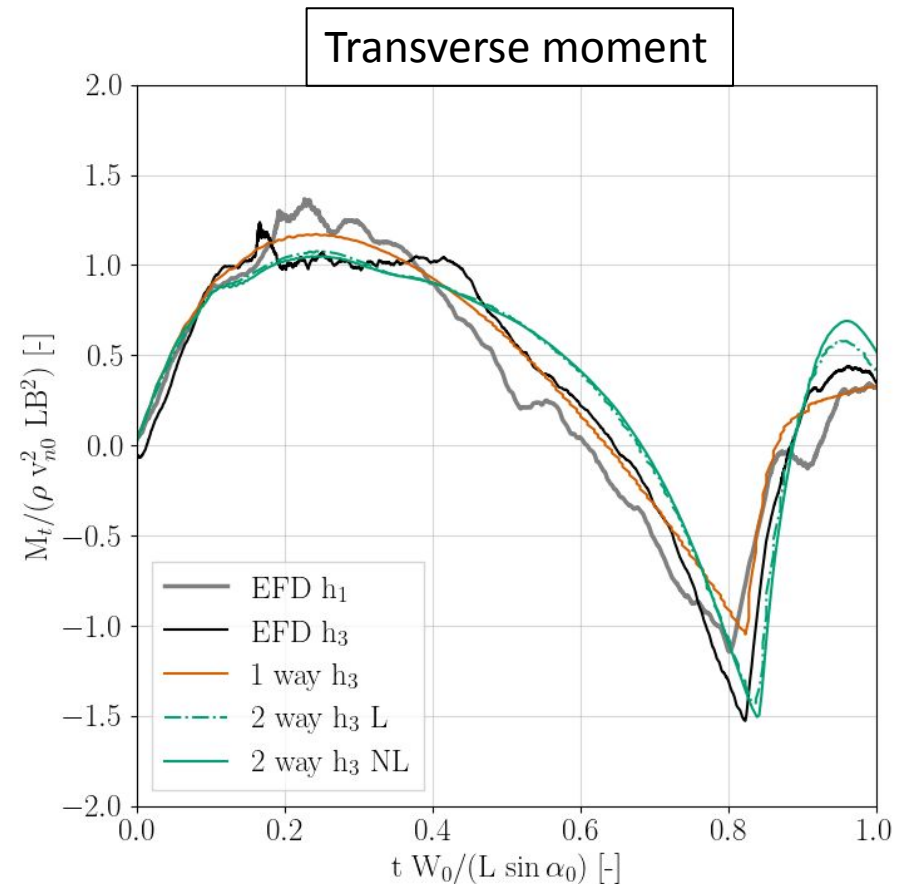
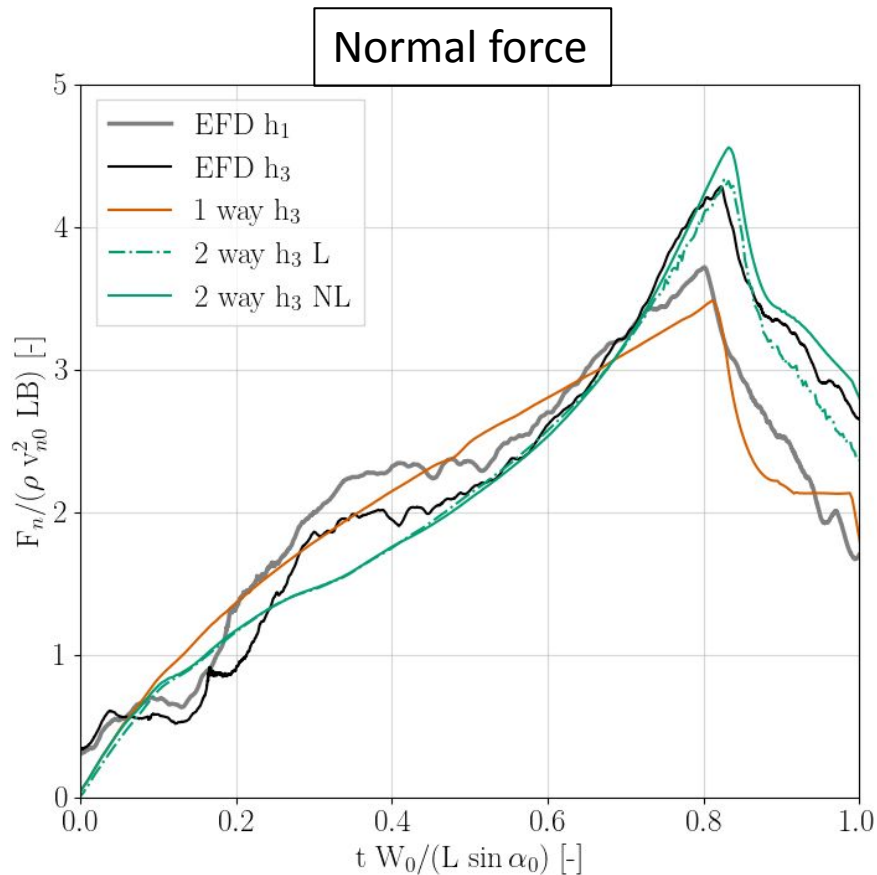
Deflection, nonlinear versus linear evaluation



- CFDSHIP-IOWA V4.5
 - Single phase
- Non-intrusive reduced order methods (ROM of discretized PDEs)
 - Any structure
 - Low computational cost
 - Easy to couple with other solvers for multidisciplinary analysis
- Any CSD software can be used
- The ROM can be easily trained for the specific FSI phenomena with a limited computational cost



- Vertical slam
- Flexible plate shows initially a reduced load that significantly grows greater than with the rigid plate in the second part of the slam
- The peak values for the flexible plate are also postponed w.r.t. the rigid plate
- For the flexible plate the force grows concave whereas it is linear for the rigid plate
- The predictions are reasonably accurate, especially in the concave growth phase



- Vertical slam
- Prediction accuracy is probe dependent, better agreement with EFD for probes on the edges of the plate
- Overall, the deflection distribution of the plate is well captured

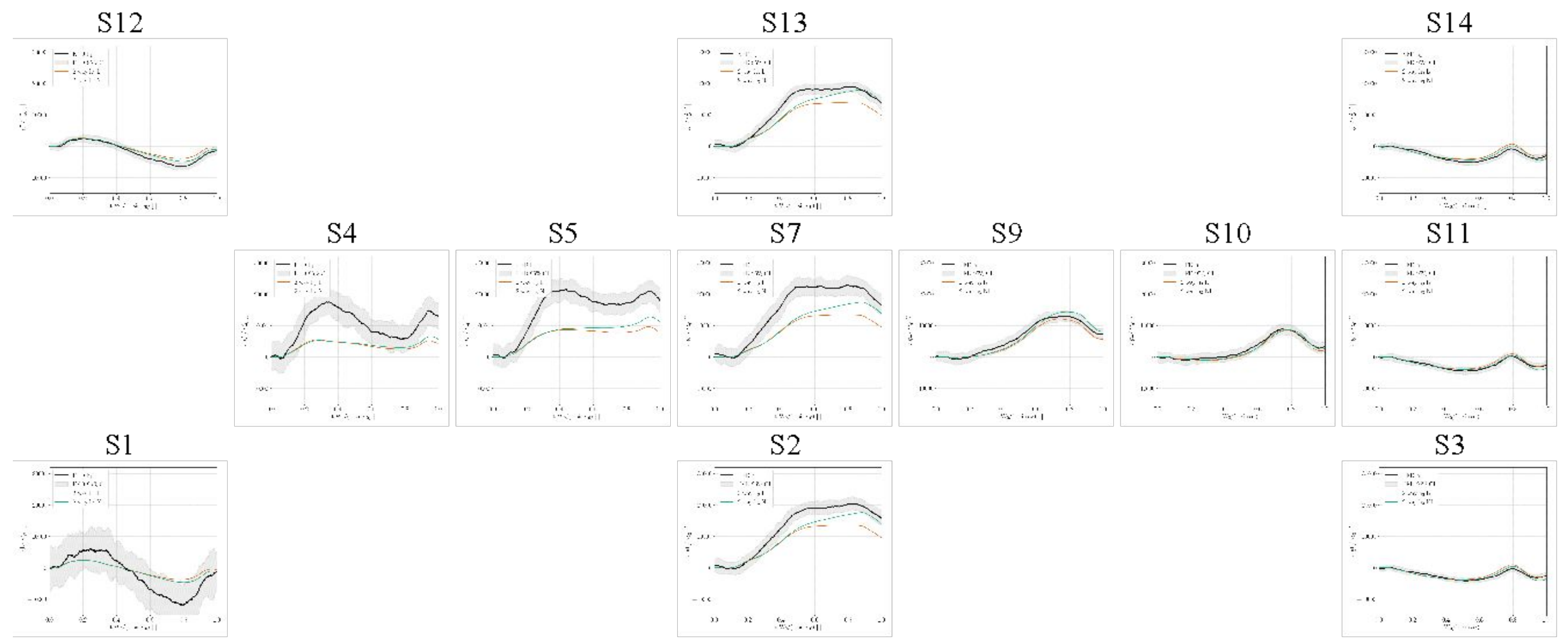
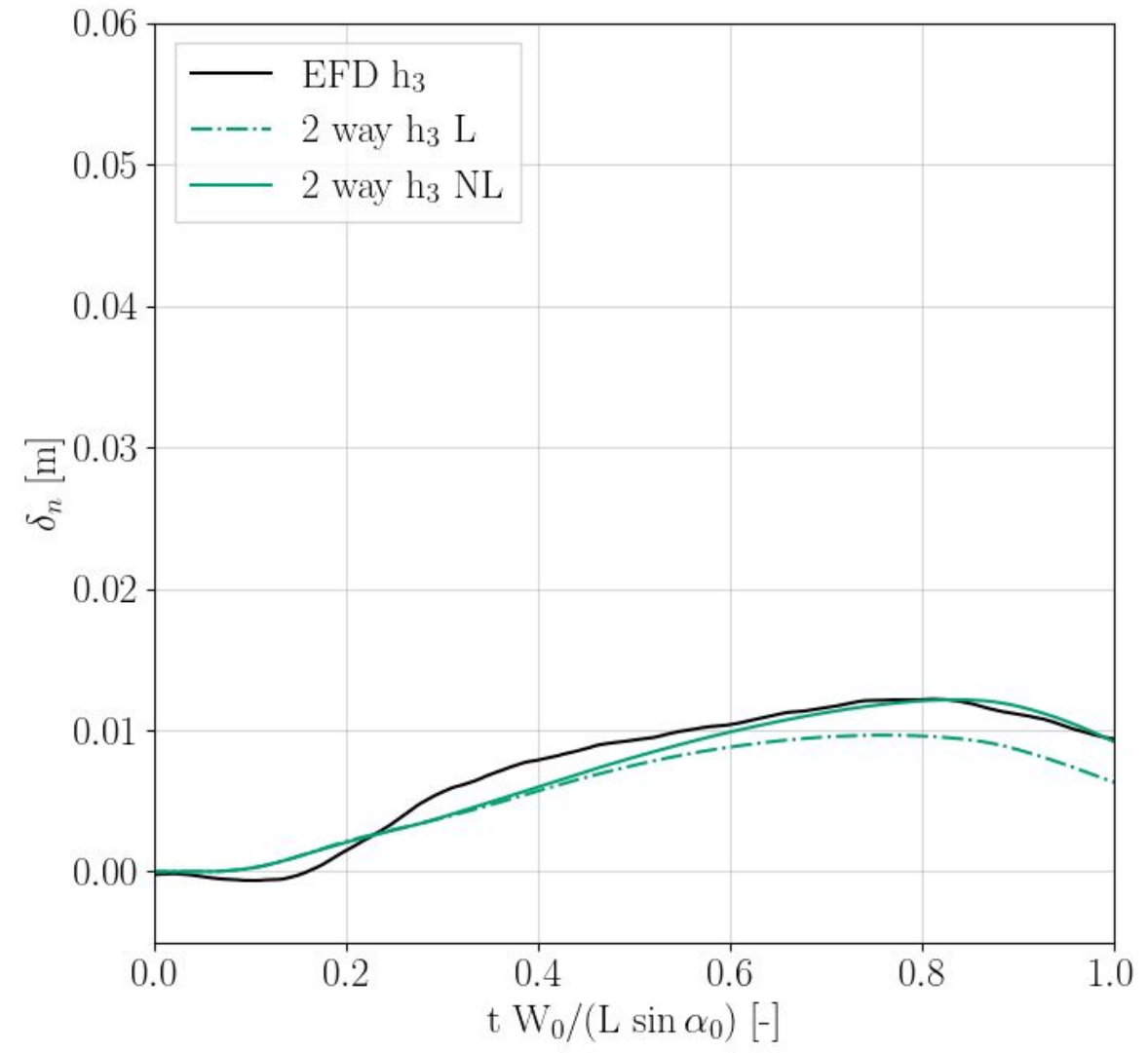
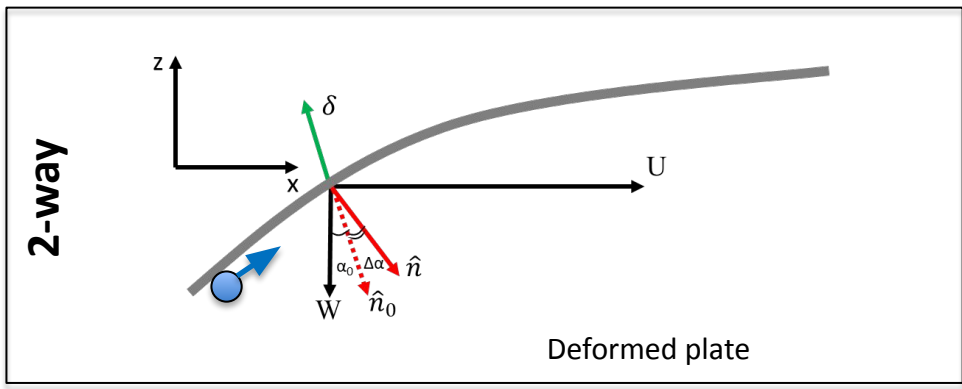
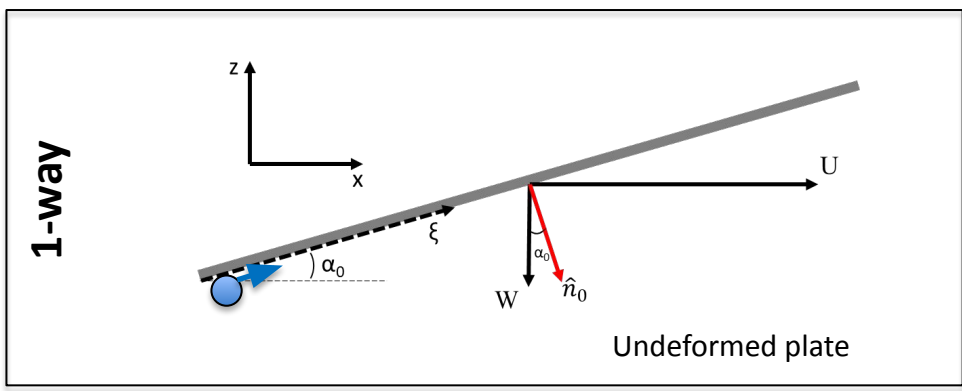


Figure 7.2: V.269 – h_3 , strain comparison between experimental and numerical results at probes location from the trailing (left) to the leading edge (right).

- Vertical slam
- Normal deflection prediction with nonlinear solver is accurate, especially in predicting the peak value
- The linear solver can predict the time when the peak value occurs but underpredict the peak value
- For small deformations the linear and nonlinear solvers predict the same values



- Vertical slam
- Significant differences exist between the rigid and flexible plates
- Qualitative analysis suggests an interaction of the high pressure ridge with the plate deformation



Rigid plate

Pressure



Flexible plate

Pressure



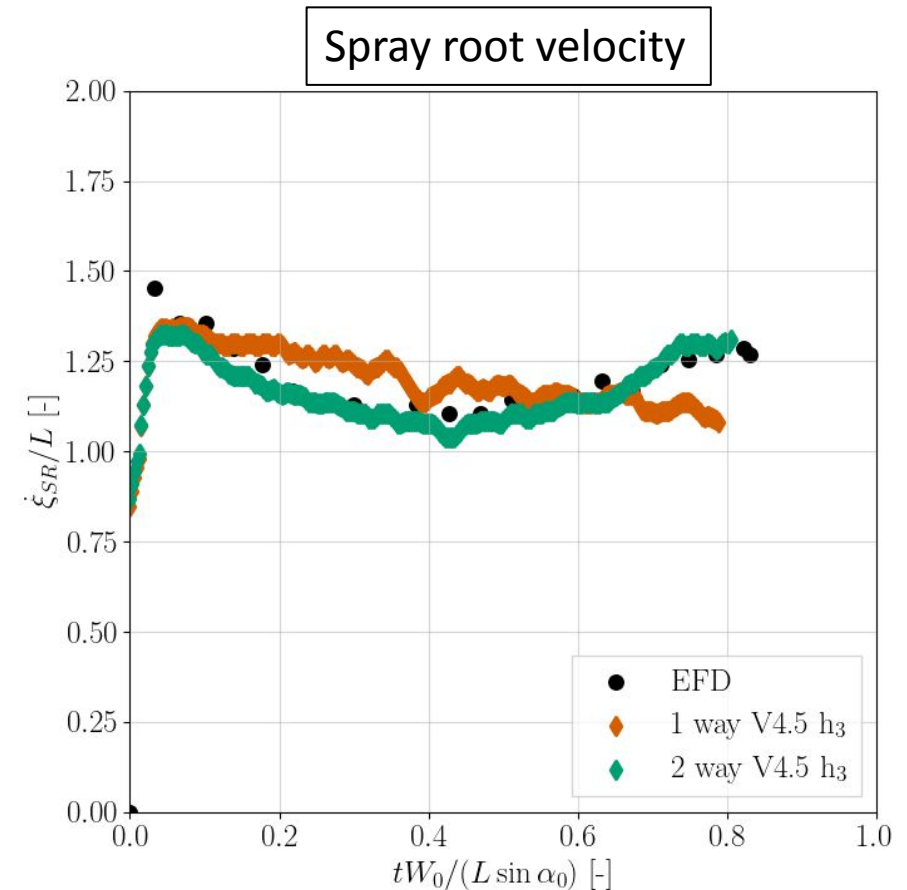
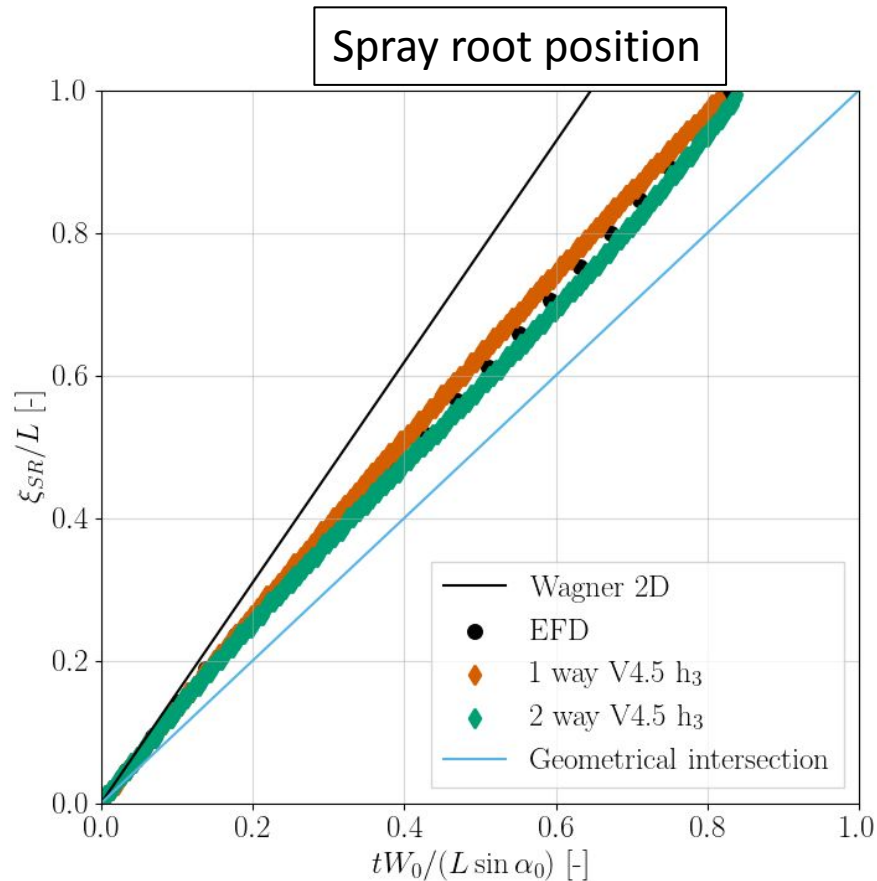
Displacement



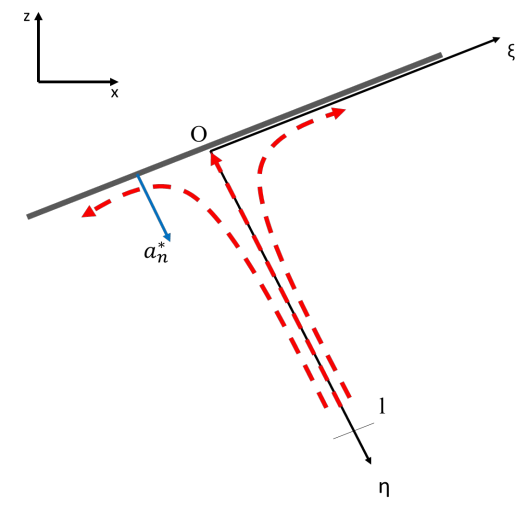
Local variation of the impact angle



- Vertical slam
- EFD and one- and two-way FSI simulations lies between the geometrical intersection and 2D Wagner prediction
- FSI predictions are reasonably accurate
- Rigid plate show a monotonic decrease of the spray root speed
- Flexible plate show a decrease and then increase of the spray root speed
- The spray root moves slower for the flexible plate in the first part of the slam and then accelerates in the final part, when enters the region with a smaller pitch angle

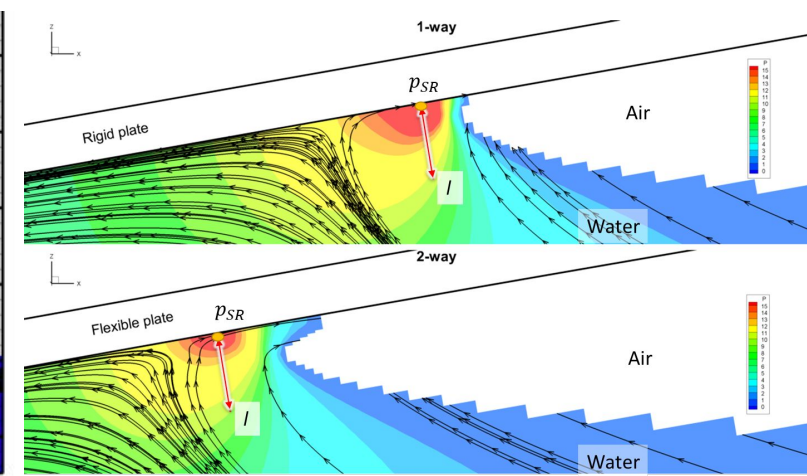
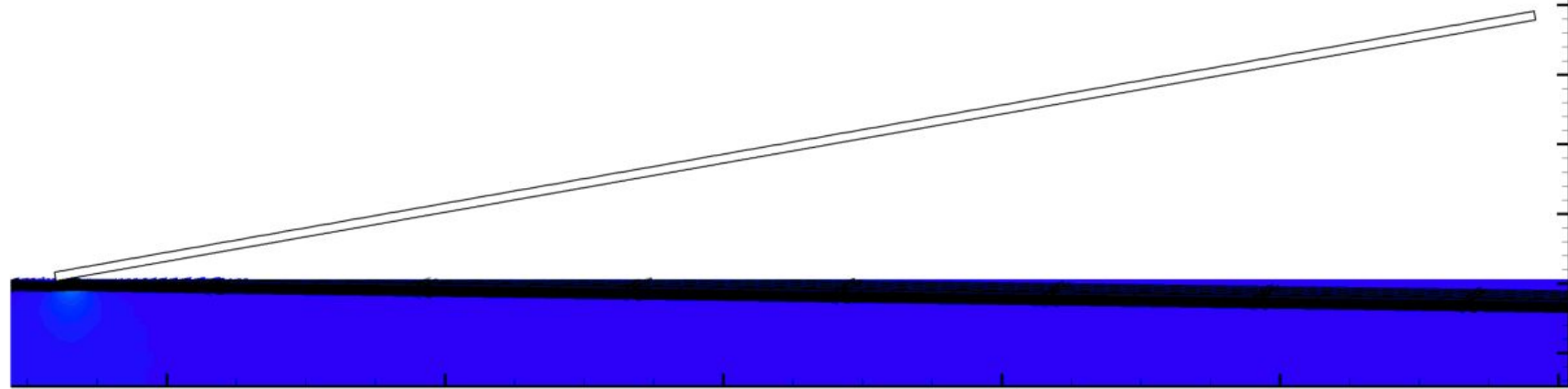


- To investigate the role of the plate kinematics on the resulting forces an investigation of the spray root along the centerplate is performed
- The fluid is modeled as an inner wall jet boundary layer embedded in an outer stagnation flow
- The outer flow spray root physics can be modeled leveraging an analogy with an inviscid quasi-steady stagnation flow with a constant normal acceleration



Moving reference system

Streamlines evaluated w.r.t. the spray root velocity



- The mono dimensional (along n , direction normal to the plate surface) Euler equations, in a moving reference frame with origin at the spray root point, with body force constant normal acceleration a_n^* with potential $\psi(n)$ are

$$\bar{u} \cdot \frac{\partial \bar{u}}{\partial \eta} = -\frac{1}{\rho} \frac{\partial p}{\partial \eta} + \bar{a}$$

$$\bar{a} = \nabla \psi, \text{ with } \psi(\eta) = a_n^* \eta$$

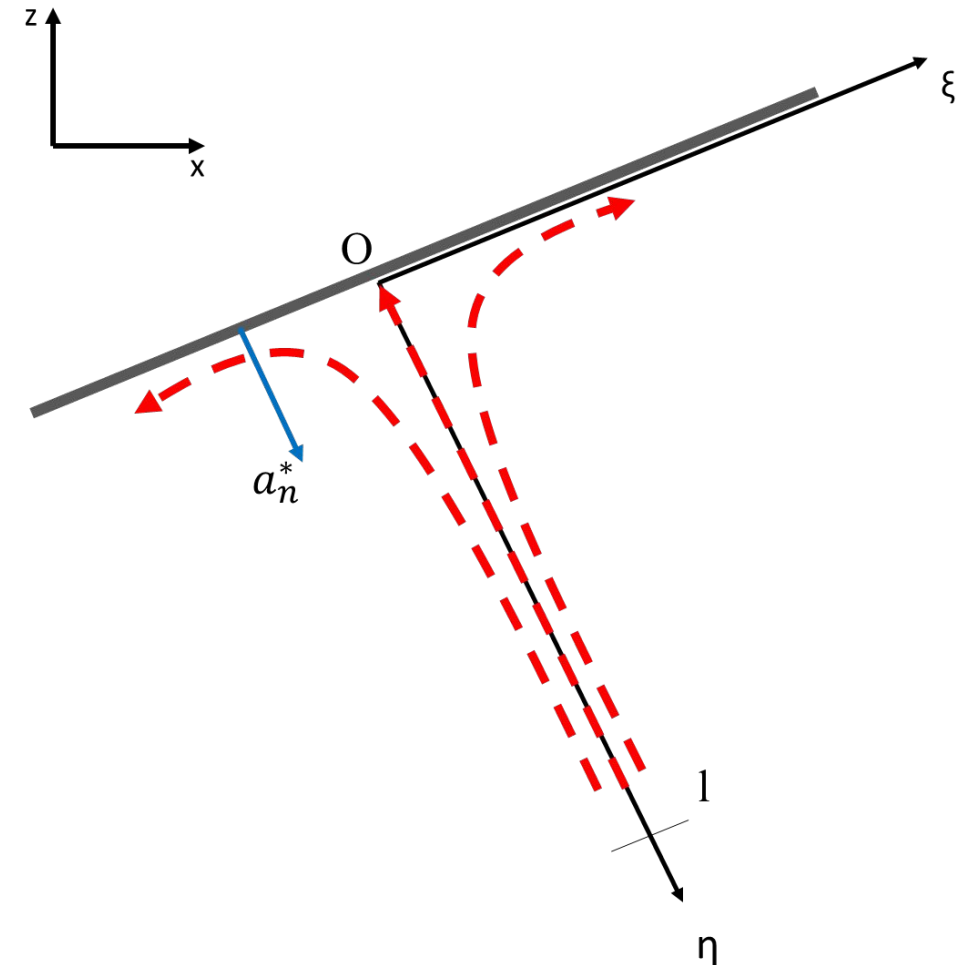
$$\frac{\partial}{\partial \eta} \left(\frac{u^2}{2} + \frac{1}{\rho} p + \psi \right) = 0$$

Integrating between $\eta=0$ and $\eta=l$ with the following boundary conditions:

$$\eta = 0 \rightarrow \begin{cases} u(0) = 0 \\ \psi(0) = 0 \\ p(0) = p_{SR} \end{cases} \quad \eta = l \rightarrow \begin{cases} u(l) = -v_n \\ \psi(l) = a_n^* l \\ p(l) = 0 \end{cases}$$

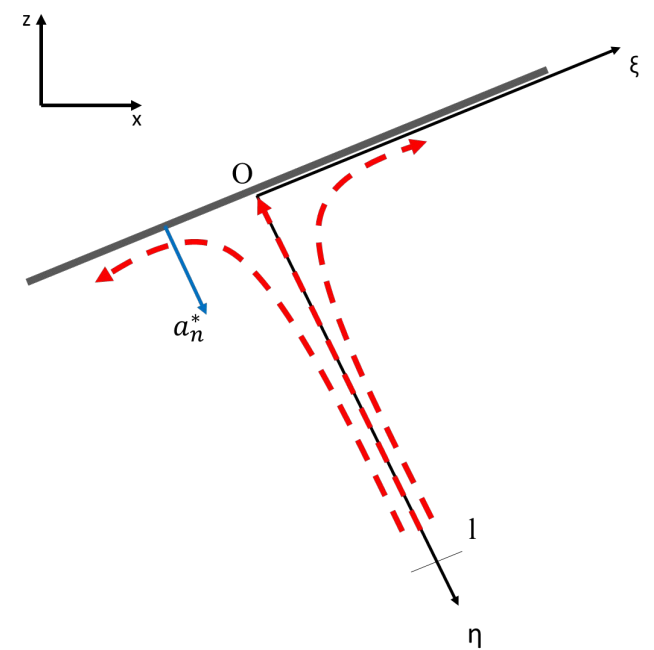
Yields:

$$p_{SR} = \rho \left(\frac{v_n^2}{2} - a_n^* l \right)$$

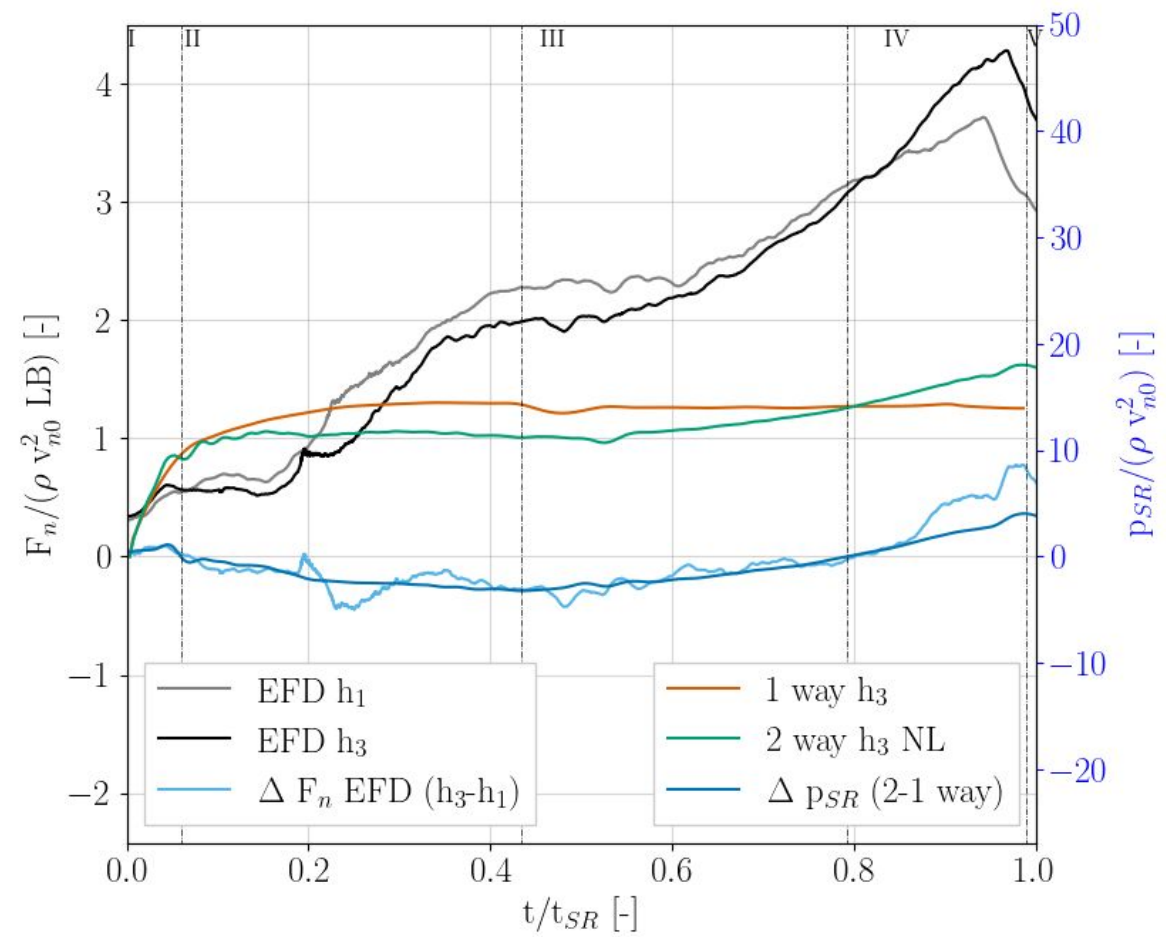


- The resulting pressure difference between a case with acceleration (flexible plate) and without acceleration (rigid plate) can be approximated with the following formula
- In this work the proportional sign is used because the current problem is not a stagnation point but an impact flow with free surface
- The analogy with a stagnation-point flow is used only as a model to explain the trends of rigid versus flexible plates

$$\Delta p_{SR} = \rho \left(\frac{\Delta v_n^2}{2} - a_n^* l \right) = p_{SR}|_{2way} - p_{SR}|_{1way} \propto \rho \left(\frac{\Delta v_n^2}{2} - l a_n^* \right)$$

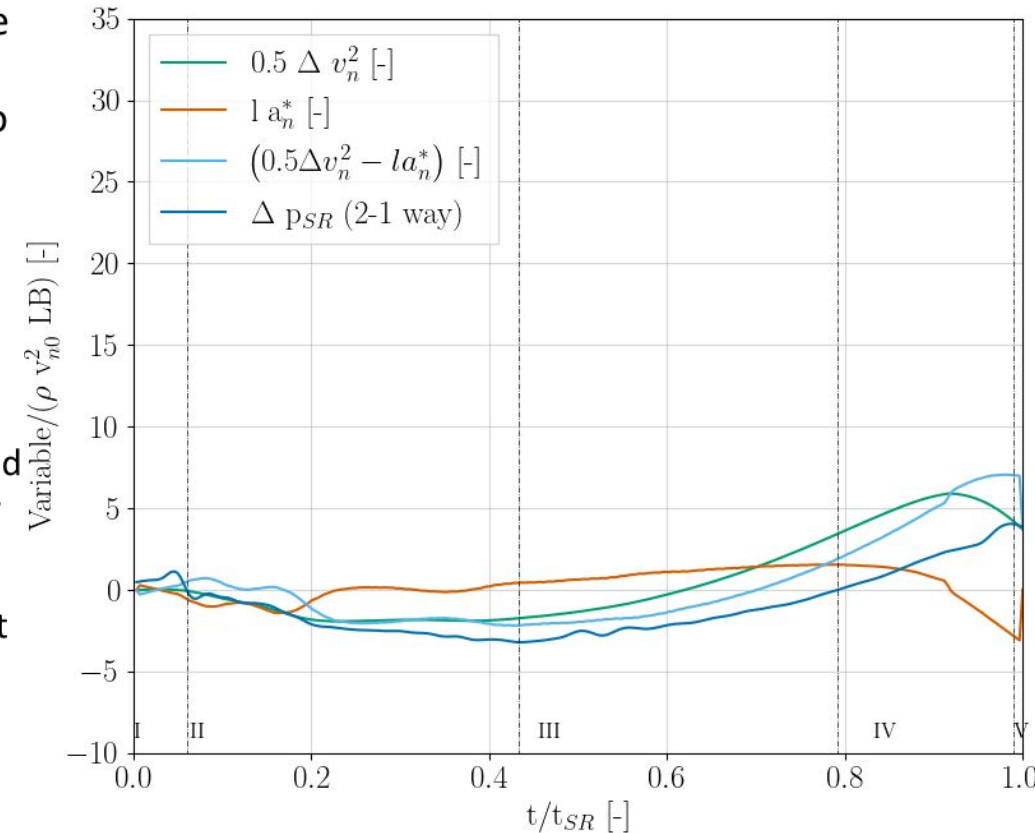


- Differences in rigid versus flexible deformation, loads, and spray root velocity explain how the interplay of global/local loads and deformation is an important factor in the physics of the hydroelastic response
- Δp_{SR} shows a good correlation with ΔF_n
- All three slams cases can be divided into five phases based on Δp_{SR}
 - I. $\Delta p_{SR} \approx 0$
 - II. Δp_{SR} decreases to its minimum
 - III. Δp_{SR} increases to zero
 - IV. Δp_{SR} increases to its maximum
 - v. the spray root exits the plate, i.e. $t = t_{SR}$
- For phases II-III Δp_{SR} is negative, therefore the deformation reduces the load, whereas for phase IV Δp_{SR} is positive, meaning the deformation dynamics increases the load



$$\Delta p_{SR} = p_{SR}|_{2way} - p_{SR}|_{1way} \propto \rho \left(\frac{\Delta v_n^2}{2} - l a_n^* \right), l = 0.01$$

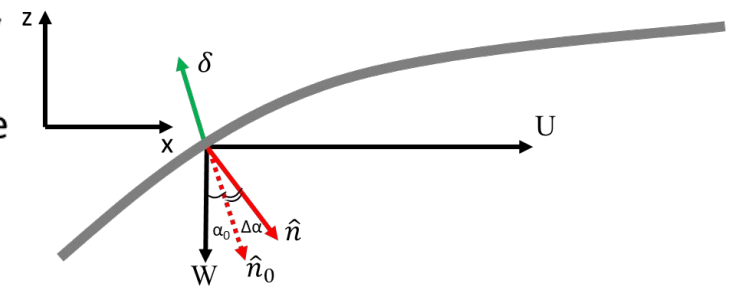
- The proportional terms of the extended Bernoulli equation model reproduce quite well Δp_{SR} trend, although it overpredicts Δp_{SR}
- The simplified model can explain the physics and the dynamics associated to the peak value of the pressure:
 - I. Δv_n^2 is almost zero and decreases, the acceleration decreases reaching its minimum, normal impact velocity and acceleration have opposite effects on Δp_{SR}
 - II. Δv_n^2 is negative and minimum, the acceleration is negative and increases to zero; in this phase velocity and acceleration have opposite effects
 - III. Δv_n^2 goes from negative to positive reaching its maximum, the acceleration increases and achieves its maximum, the two terms both contribute to the load reduction in the first part of this phase; the 2-way FSI load is larger than the 1-way and reaches its maximum along with (positive) Δp_{SR}
 - IV. Δv_n^2 is positive, and the acceleration significantly decreases, both terms contribute to the load increase, especially the acceleration in the very last part where the pressure peak occurs
 - v. both velocity and acceleration terms go to zero
- The load-reduction phase corresponds to minimum impact velocity and positive acceleration
- The load-increase phase corresponds to an increase of the impact velocity along with a significant negative minimum value for the deformation acceleration



- To understand the role of the plate kinematics on the impact force differences between the rigid and flexible plate the quantities that concur to the evaluation of the extended Bernoulli equation are investigated

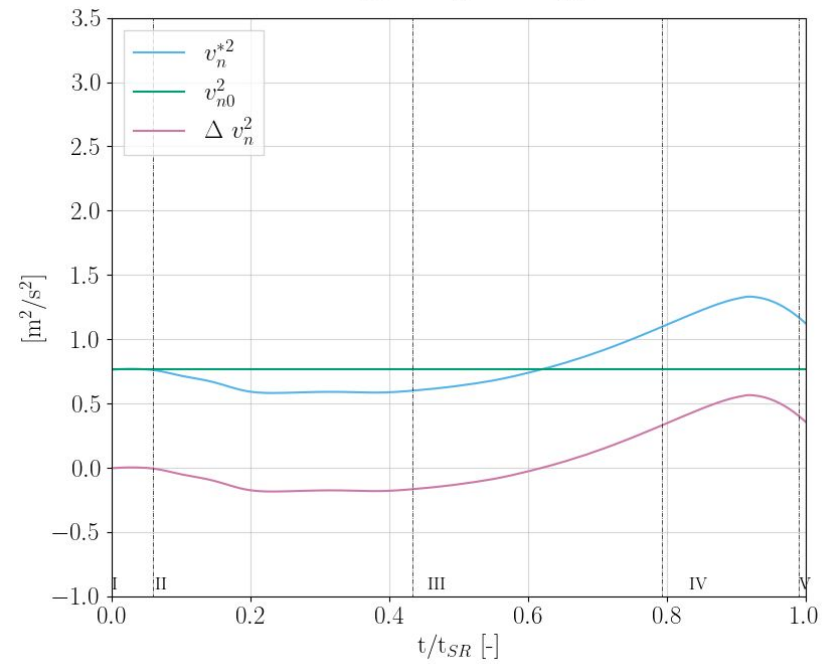
- $\Delta v_n^2 = v_n^{*2} - v_{n0}^2$ ← The difference of the normal impact velocity due to the plate deformation squared
- $v_{n0} = (W \cos \alpha_0 + U \sin \alpha_0)$ ← The normal impact velocity of the undeformed plate
- $v_n^* = \frac{D\delta_n}{Dt} + v_n$ ← The normal impact velocity at the spray root considering the local impact angle and deformation velocity
- $\frac{D\delta_n}{Dt} = \frac{\partial\delta_n}{\partial t} + v_{SR} \frac{\partial\delta_n}{\partial x}$ ← The material derivative of the plate deformation
- $v_n = (W \cos \alpha + U \sin \alpha)$ ← The normal impact velocity at the spray root considering only the local impact angle
- $\alpha = \alpha_0 + \Delta\alpha$ ← The impact angle
- $\Delta\alpha = -\frac{\partial\delta_n}{\partial x}$ ← The local variation of the impact angle at the spray root
- $a_n^* = \frac{Dv_n^*}{Dt} = \frac{D}{Dt} \left(\frac{D\delta_n}{Dt} + v_n \right) = \frac{D^2\delta_n}{Dt^2} + \frac{Dv_n}{Dt}$ ← The acceleration at the spray root
- $\frac{D^2\delta_n}{Dt^2} = \left(\frac{\partial^2\delta_n}{\partial t^2} + \frac{\partial v_{SR}}{\partial t} \frac{\partial\delta_n}{\partial x} + v_{SR} \frac{\partial^2\delta_n}{\partial t\partial x} \right) + \left(v_{SR} \frac{\partial^2\delta_n}{\partial t\partial x} + v_{SR}^2 \frac{\partial^2\delta_n}{\partial x^2} + v_{SR} \frac{\partial v_{SR}}{\partial x} \frac{\partial\delta_n}{\partial x} \right)$
- $\frac{Dv_n}{Dt} = \left(-W \frac{\partial\alpha}{\partial t} \sin \alpha + U \frac{\partial\alpha}{\partial t} \cos \alpha \right) + \left(-W v_{SR} \frac{\partial\alpha}{\partial x} \sin \alpha + U v_{SR} \frac{\partial\alpha}{\partial x} \cos \alpha \right)$

- Δv_n^2 is first negative and then positive meaning that the impact velocity is first reduced and then increased
- The time derivative of the deflection ($\partial\delta_n/\partial t$) provides the reduction of the impact velocity in the first part of the slam, the plate is moving away from the spray root
- In the second part $v_{SR}\partial\delta_n/\partial x$ provides the greater contribution meaning that the curvature of the plate plays a significant role in this phase in increasing the impact velocity and thus the pressure
- The local pitch angle changes significantly during the slam, achieving its minimum when the normal force is at its maximum



The difference of the normal impact velocity due to the plate deformation squared

$$\Delta v_n^2 = v_n^{*2} - v_{n0}^2$$

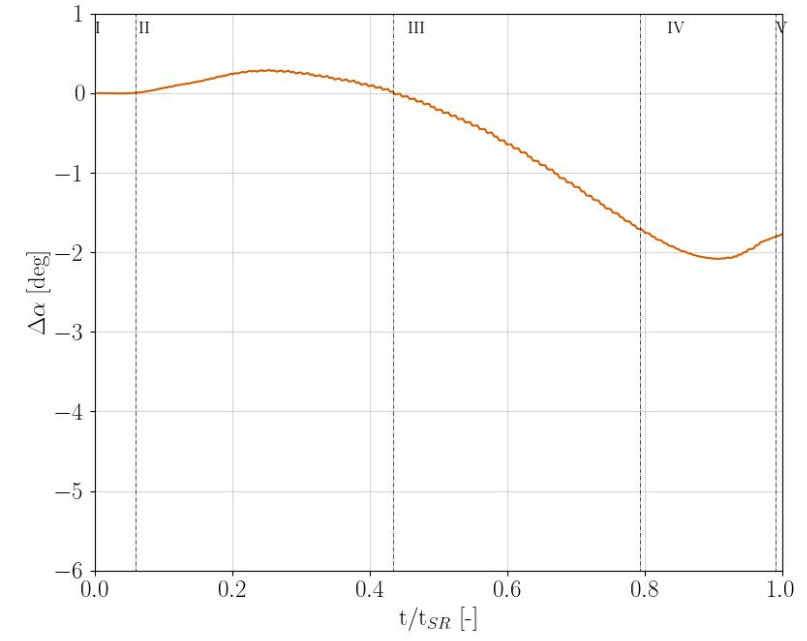
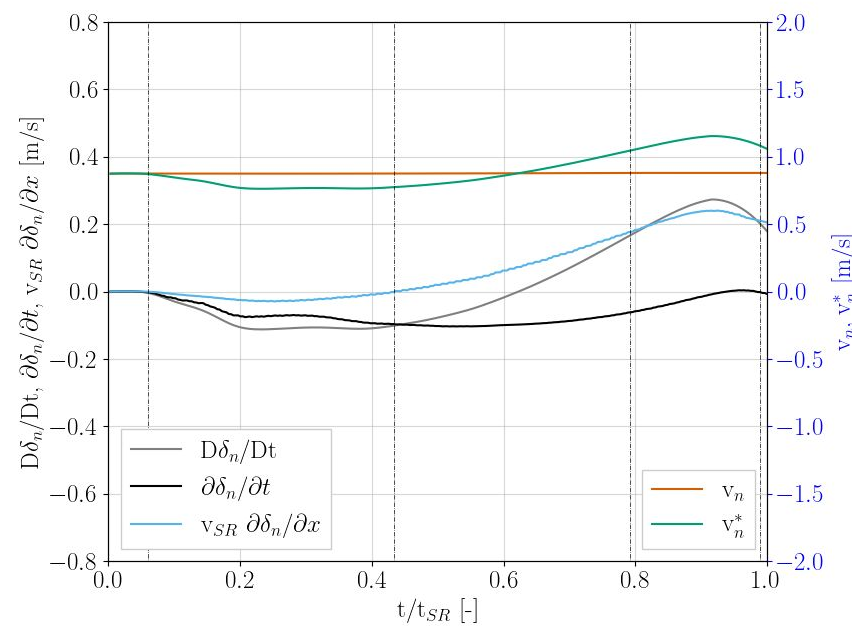


$$v_n^* = \frac{D\delta_n}{Dt} + v_n$$

The normal impact velocity at the spray root considering the local impact angle and deformation velocity

$$\frac{D\delta_n}{Dt} = \frac{\partial\delta_n}{\partial t} + v_{SR} \frac{\partial\delta_n}{\partial x}$$

The material derivative of the plate deformation

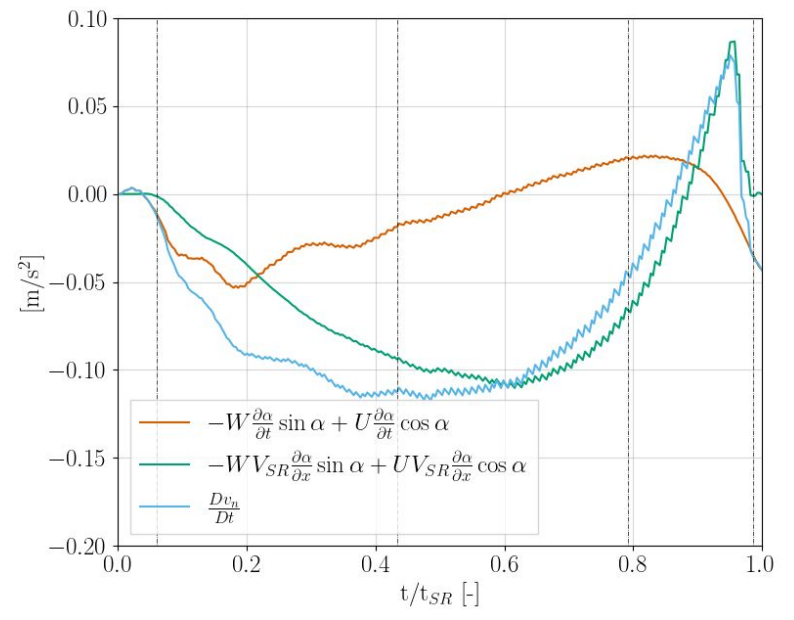
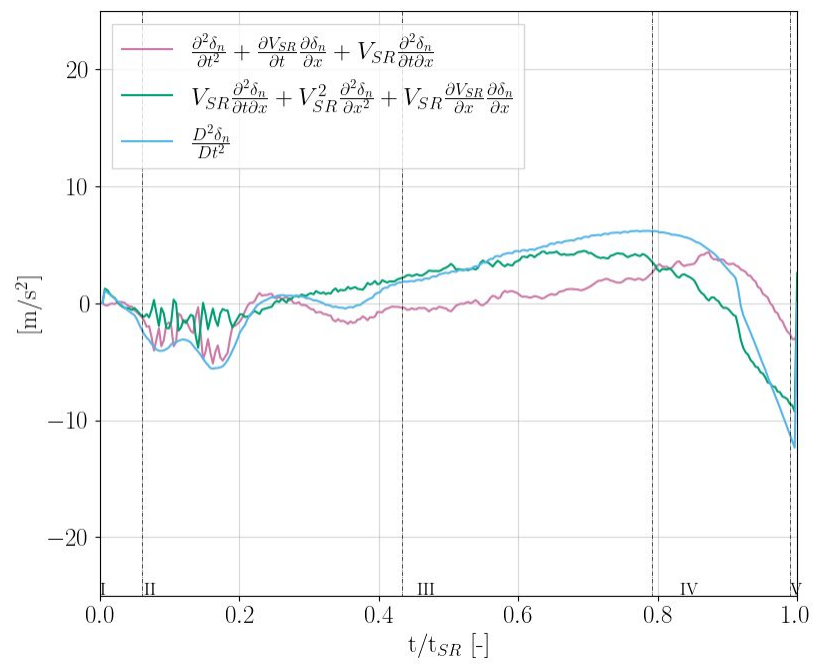
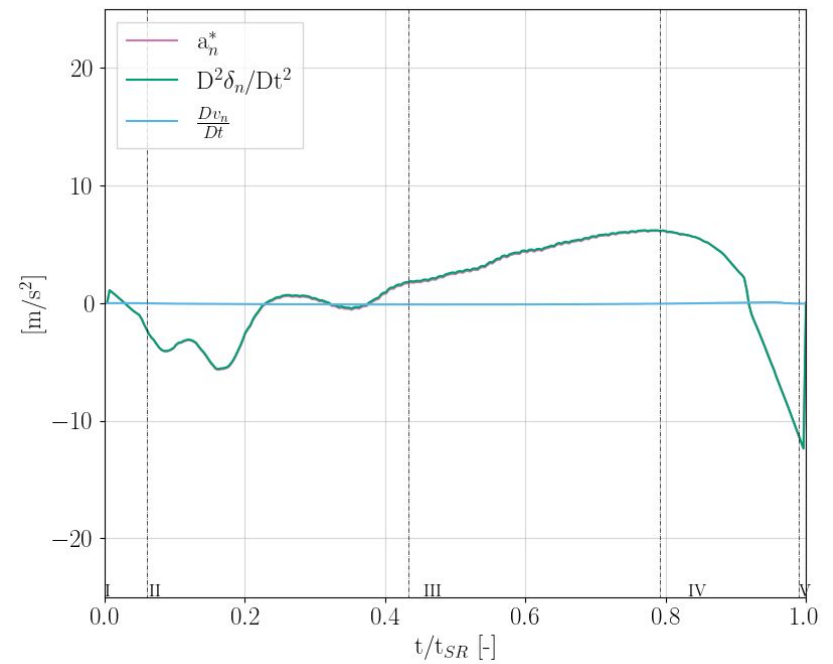


- The term $\frac{D^2 \delta_n}{Dt^2}$ provides the greatest contribution to a_n^* meaning that the geometric deflection provides a greater contribution to the acceleration than the local pitch angle variation
- The components of $\frac{D^2 \delta_n}{Dt^2}$ have similar contribution to their resulting sum
- The convective term of $\frac{Dv_n}{Dt}$ have a greater influence than the time varying, confirming the significant role of the geometric deflection

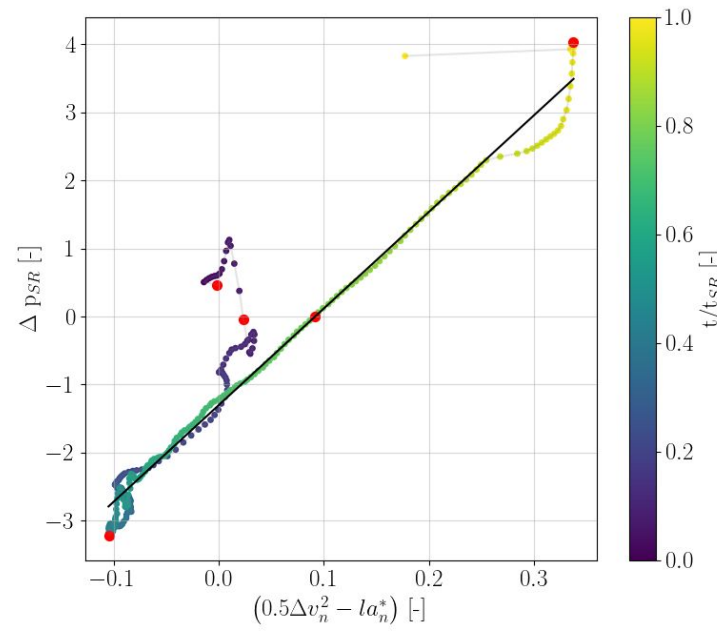
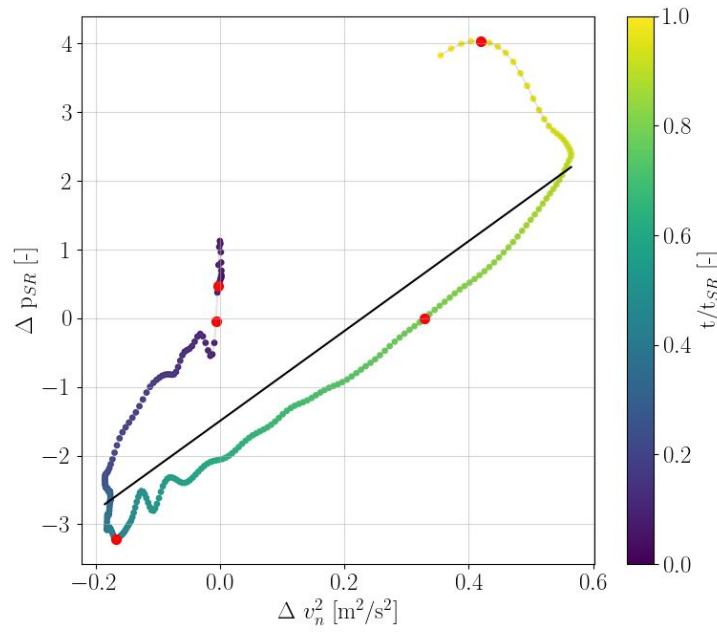
$$a_n^* = \frac{D^2 \delta_n}{Dt^2} + \frac{Dv_n}{Dt}$$

$$\frac{D^2 \delta_n}{Dt^2} = \left(\frac{\partial^2 \delta_n}{\partial t^2} + \frac{\partial V_{SR}}{\partial t} \frac{\partial \delta_n}{\partial x} + V_{SR} \frac{\partial^2 \delta_n}{\partial t \partial x} \right) + \left(V_{SR} \frac{\partial^2 \delta_n}{\partial t \partial x} + V_{SR}^2 \frac{\partial^2 \delta_n}{\partial x^2} + V_{SR} \frac{\partial V_{SR}}{\partial x} \frac{\partial \delta_n}{\partial x} \right)$$

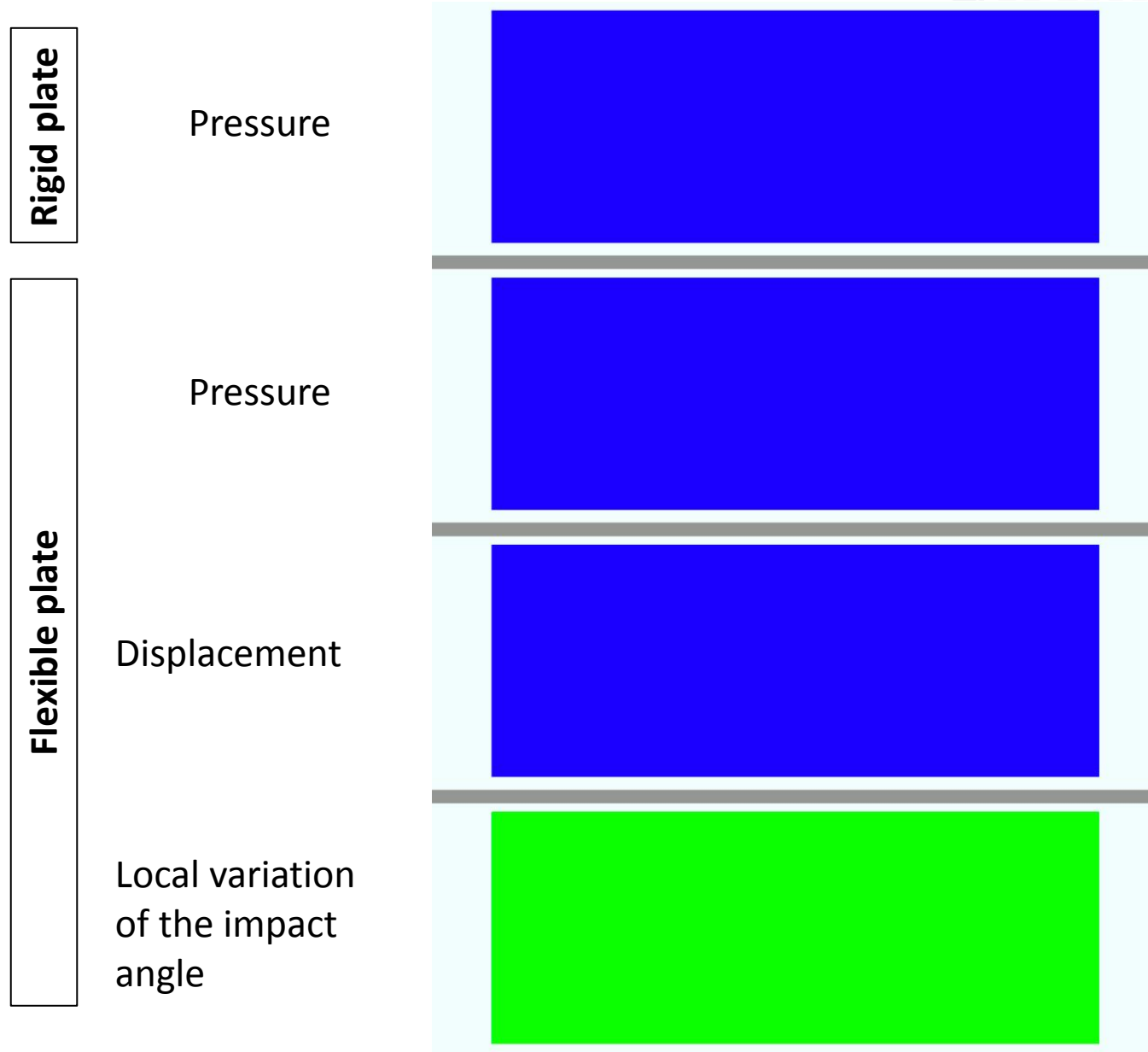
$$\frac{Dv_n}{Dt} = \left(-W \frac{\partial \alpha}{\partial t} \sin \alpha + U \frac{\partial \alpha}{\partial t} \cos \alpha \right) + \left(-W V_{SR} \frac{\partial \alpha}{\partial x} \sin \alpha + U V_{SR} \frac{\partial \alpha}{\partial x} \cos \alpha \right)$$



- The correlation between Δp_{SR} and Δv_n^2 exists only phase wise in the first part of the slam
- The correlation of Δp_{SR} with Δv_n^2 cannot explain the peak phase (IV)
- The correlation of Δp_{SR} with $\left(\frac{\Delta v_n^2}{2} - la_n^*\right)$ is significantly better than using Δv_n^2 only. The proportional terms correlate very well with phase IV of Δp_{SR}
- Although the extended Bernoulli equation is a simplified model, it predicts accurately Δp_{SR} during phase IV, explaining the contribution of the deformation velocity and acceleration to the increased pressure peak
- The acceleration experienced by the spray root is fundamental to the final pressure/force increase
- A reduction of the acceleration experienced by the spray root would lead to a smaller peak in the final part of the slam

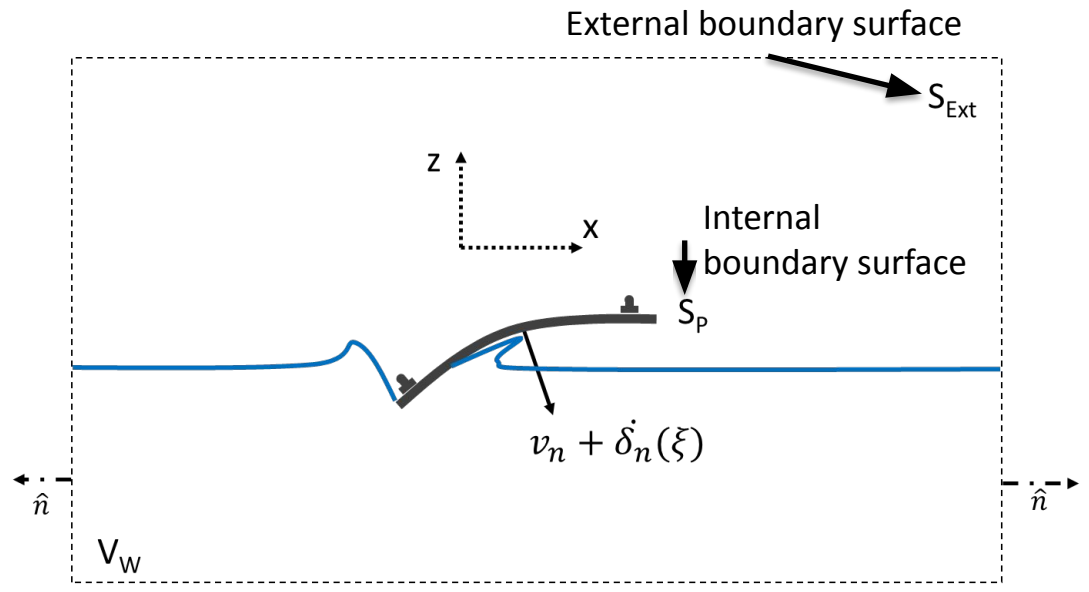


- During the slam the plate undergoes significant deformations that interact with the spray root
- The analysis of the energy conservation during the slam can help in understanding the role of the structure on the spray root dynamics and therefore on the pressure acting on the structure itself



Water (subscript w)

- the energy flux terms are zero due to the rigid wall and adiabatic boundary conditions
- neglecting air and considering only the rate of pressure work done across the wetted surface of the plate S_{wet}



\mathbf{u}_w = the fluid velocity

\dot{W}_{pR} = power provided by the pressure due to rigid body motion

$\dot{W}_{p\delta_n}$ = power provided by the pressure due to the plate's flexibility

$$e = k_e + p_{eg}$$

$$v_n = \mathbf{v} \cdot \hat{n}$$

$$\delta_n = \boldsymbol{\delta} \cdot \hat{n}$$

$$\frac{dE}{dt} = \frac{\partial}{\partial t} \iiint_{V_w} \rho \left(gz + \frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = -\frac{\delta W}{dt} = \iint_{S_{wet}} p [v_n + \dot{\delta}_n] dS = \dot{W}_{pR} + \dot{W}_{p\delta_n}$$

Structure (subscript s)

- the energy flux terms are zero due to the impermeability and adiabatic boundary conditions
- neglecting air and considering only the rate of pressure work done across the wetted surface of the plate S_{wet}

$$\frac{\partial}{\partial t} \iiint_{V_s} \rho e dV = - \iint_{S_{wet}} p[(\mathbf{v}_n + \delta_n) \cdot \hat{n}] dS + \mathbf{F}_M \cdot \mathbf{v}_n$$

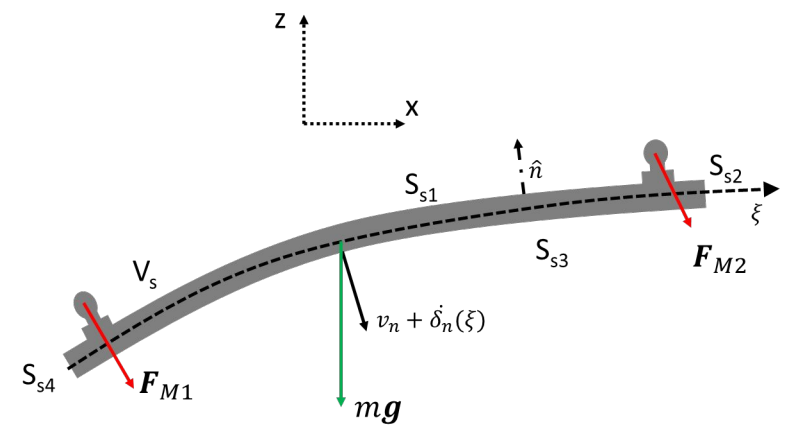
$$e = k_e + p_{e\varepsilon} + p_{eg} = \rho \frac{\|\mathbf{v}_n + \delta_n\|^2}{2} + \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 \sigma_{ij} \varepsilon_{ij} + g(z_{s0} + \delta \cdot \hat{k})$$

Elastic potential energy
 Kinetic potential energy
 Gravitational potential energy

$$v_n = \mathbf{v} \cdot \hat{n}$$

$$\delta_n = \delta \cdot \hat{n}$$

z-coordinate of the undeformed plate



\dot{K}_{δ_n} = time derivative of kinetic energy
 \dot{U}_ε = time derivative of elastic potential energy
 $\dot{U}_{\delta_n z}$ = time derivative of gravitational potential energy
 \dot{W}_M = power provided by the carriage through the mount
 \hat{n} is pointing outward for the structure but inward for the fluid

$$\dot{K}_{\delta_n} + \dot{U}_\varepsilon + \dot{U}_{\delta_n z} = -(\dot{W}_{pR} + \dot{W}_{p\delta_n}) + \dot{W}_M$$

Structure (subscript s)

- Considering an isotropic homogeneous material

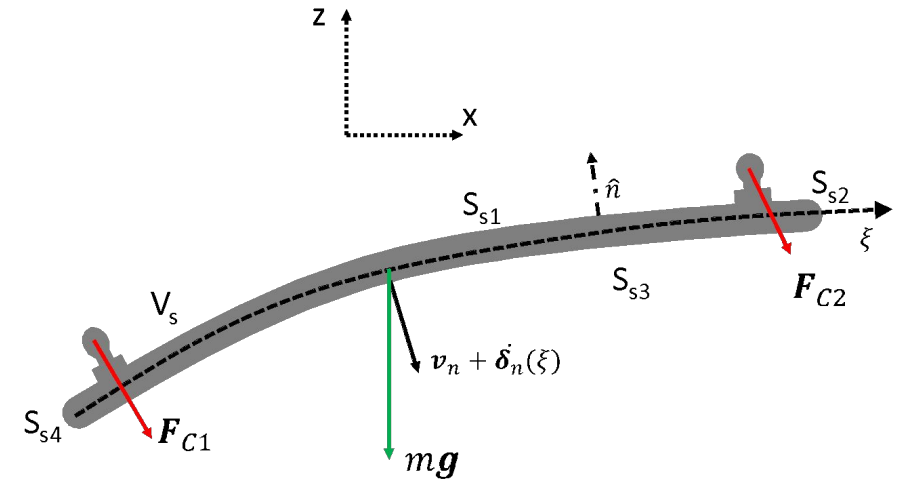
$$\sigma_{ij} = \lambda k_{ij} \varepsilon_{kk} + \mu (2\varepsilon_{ij}) \quad v_n = \mathbf{v} \cdot \hat{n}$$

$$\delta_n = \boldsymbol{\delta} \cdot \hat{n}$$

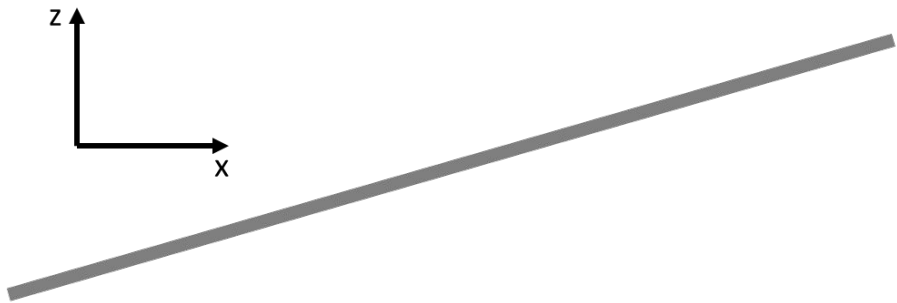
$$e = \rho \frac{\|v_n + \dot{\delta}_n\|^2}{2} + \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 [\lambda k_{ij} \varepsilon_{kk} + \mu (2\varepsilon_{ij})] \varepsilon_{ij} + g(z_{s0} + \boldsymbol{\delta} \cdot \hat{k})$$

$$\lambda = \frac{Ev}{(1+\nu)(1-2\nu)} \quad \text{and} \quad \mu = \frac{E}{2(1+\nu)}$$

$$p_{e\varepsilon} = \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 [\lambda k_{ij} \varepsilon_{kk} + \mu (2\varepsilon_{ij})] \varepsilon_{ij} = \frac{1}{2} (\lambda + 2\mu) (\varepsilon_{11}^2 + \varepsilon_{22}^2 + \varepsilon_{33}^2) + \lambda (\varepsilon_{11}\varepsilon_{22} + \varepsilon_{11}\varepsilon_{33} + \varepsilon_{22}\varepsilon_{33}) + \mu (\varepsilon_{12}^2 + \varepsilon_{13}^2 + \varepsilon_{23}^2)$$



1 is x-direction
2 is y-direction
3 is z-direction



- Rigid plate

- The mount transfers energy directly to the fluid since the plate is rigid

Water

$$\frac{\partial}{\partial t} \iiint_{V_w} \rho_w \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}(t)} p(v_n) dS = \dot{W}_{pR}$$

Structure

$$-\dot{W}_{pR} = \dot{W}_M$$

- Flexible plate

- The mount transfers energy to the fluid and to the plate (kinetic and potential elastic energy) since the plate is deforming

Water

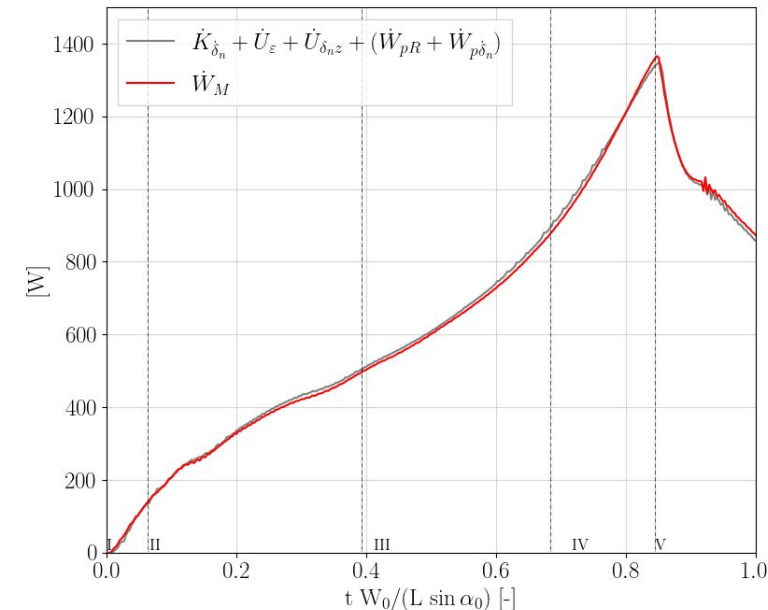
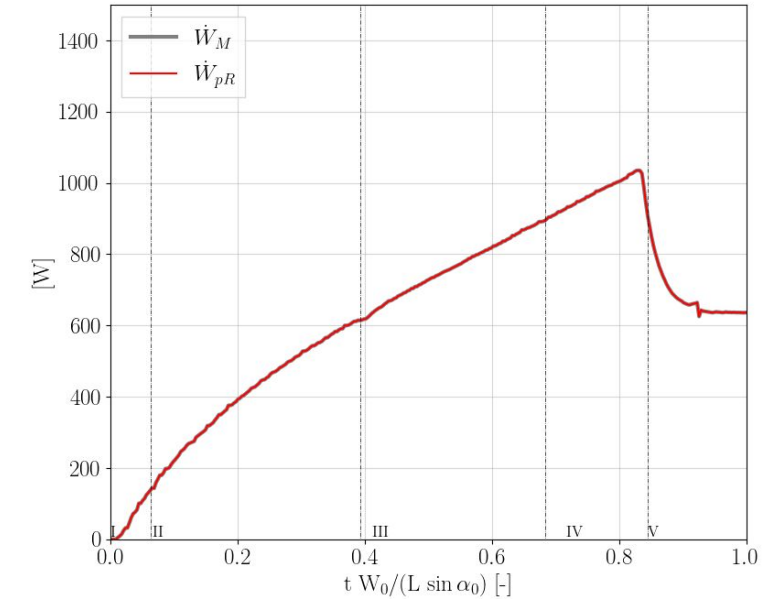
$$\frac{\partial}{\partial t} \iiint_{V_w} \rho \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}} p[v_n + \delta_n] dS = \dot{W}_{pR} + \dot{W}_{p\delta_n}$$

Structure

$$\dot{K}_{\delta_n} + \dot{U}_\varepsilon + \dot{U}_{\delta_n z} + (\dot{W}_{pR} + \dot{W}_{p\delta_n}) = \dot{W}_M$$

$$v_n = \mathbf{v} \cdot \hat{\mathbf{n}}$$

$$\delta_n = \boldsymbol{\delta} \cdot \hat{\mathbf{n}}$$



- Subtracting the equation for the energy conservation for the flexible and the rigid cases allows to study the dynamics of the difference in the force acting on the plate between the flexible and the rigid case

Water

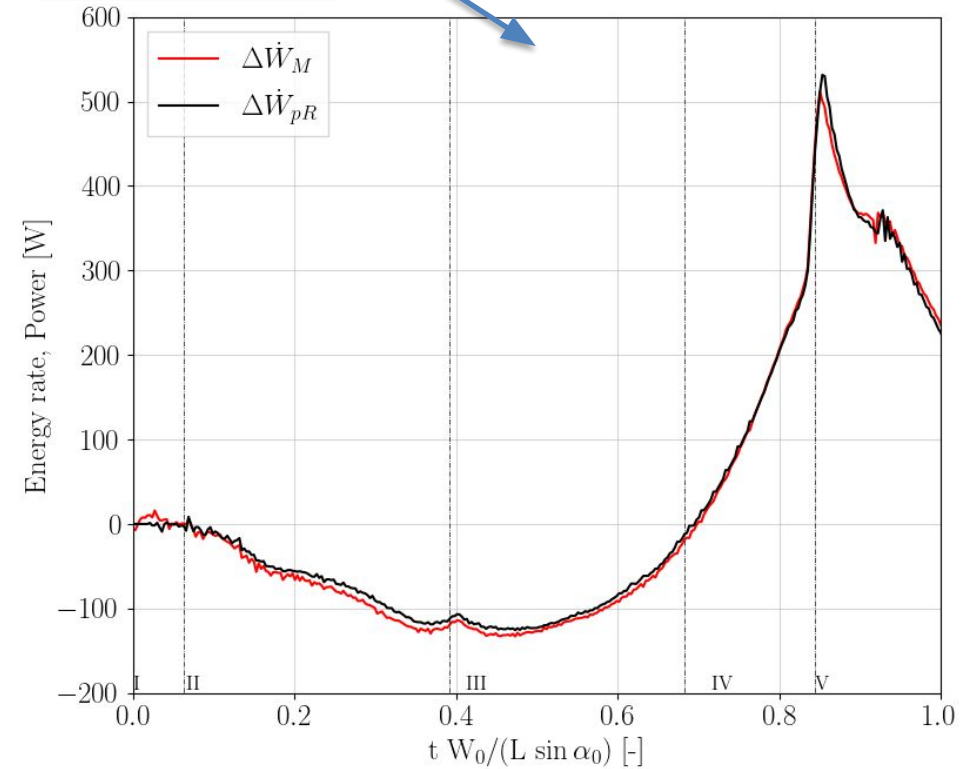
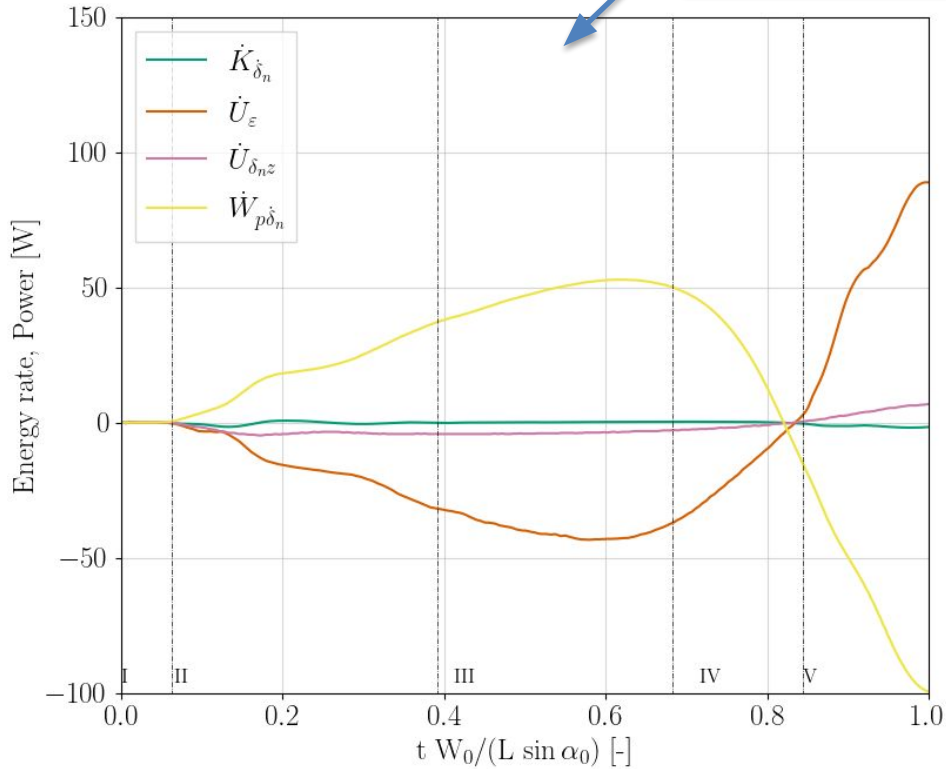
$$\begin{aligned}
 & \left[\frac{\partial}{\partial t} \iiint_{V_w} \rho \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w \right]_{Flexible} - \left[\frac{\partial}{\partial t} \iiint_{V_w} \rho \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w \right]_{Rigid} = \\
 & = \left[- \iint_{S_{Wet}(t)} p(v_n) dS \right]_{Flexible} - \left[\iint_{S_{Wet}(t)} p(v_n) dS \right]_{Rigid} + \left[- \iint_{S_{Wet}(t)} p(\dot{\delta}_n) dS \right]_{Flexible} = \\
 & = \dot{W}_{p\dot{\delta}_n} + \Delta \dot{W}_{pR}
 \end{aligned}$$

Structure

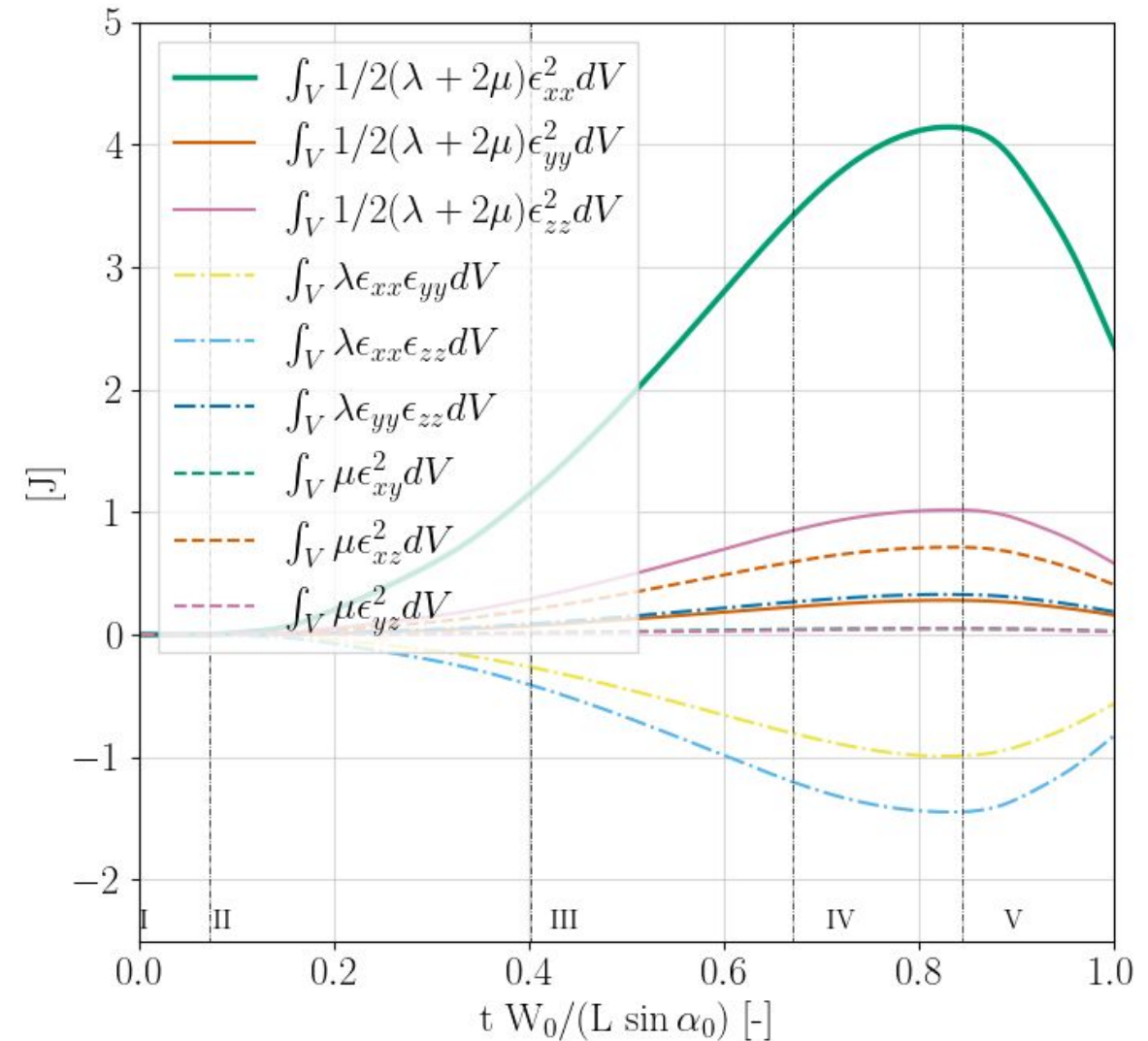
$$\dot{K}_{\dot{\delta}_n} + \dot{U}_\varepsilon + \dot{U}_{\delta_{nz}} + \dot{W}_{p\dot{\delta}_n} = \Delta \dot{W}_M - \Delta \dot{W}_{pR}$$

- The difference equation is satisfied
- The kinetic energy content is the smallest
- The elastic potential energy almost equals the power provided by the pressure due to the flexibility suggesting that most of the work done from the fluid to the plate goes into the plate's deformation

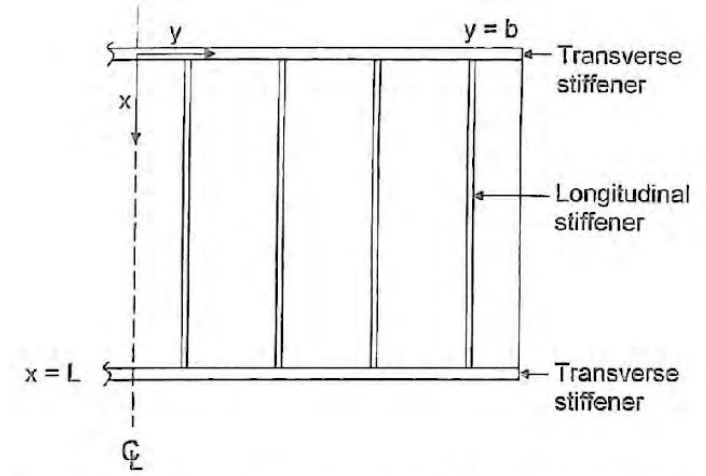
$$\dot{K}_{\delta_n} + \dot{U}_\varepsilon + \dot{U}_{\delta_n z} + \dot{W}_{p\delta_n} = \Delta\dot{W}_M - \Delta\dot{W}_{pR}$$



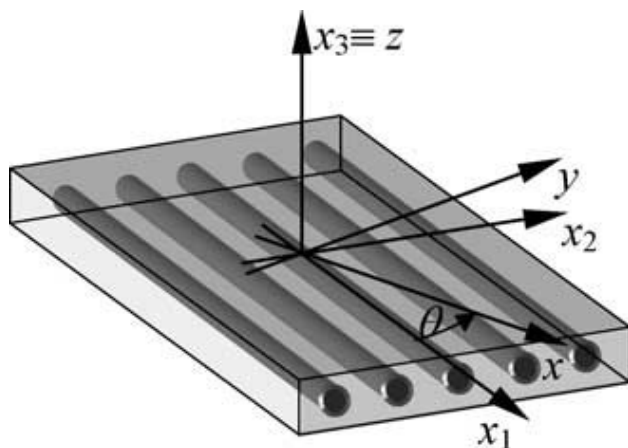
- For all the cases the longitudinal (ϵ_{xx}^2) strain contribution is the most important
- This suggests that the plate rigidity should be increased in this direction to reduce the total deflection
- The other directions are less solicited during this specific slam and therefore their rigidity could be reduced in order to reduce the structural weight



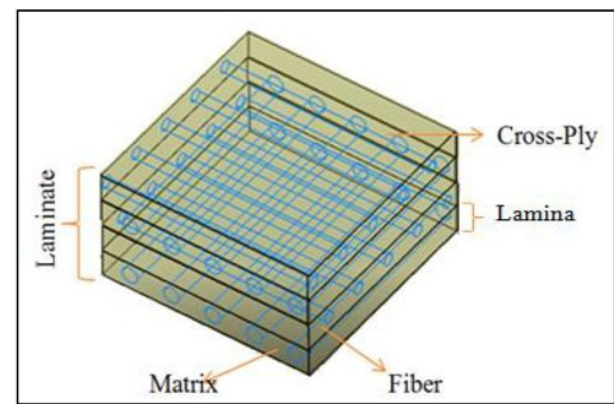
- Materials that exhibit macroscopically non-isotropic response are considered
 - Orthotropic material
 - Material properties are defined according to three principal (orthogonal) directions
 - e.g. Plates with stiffeners, single layer composites
 - Anisotropic material
 - Material properties changes with the direction
 - e.g. Laminated composites



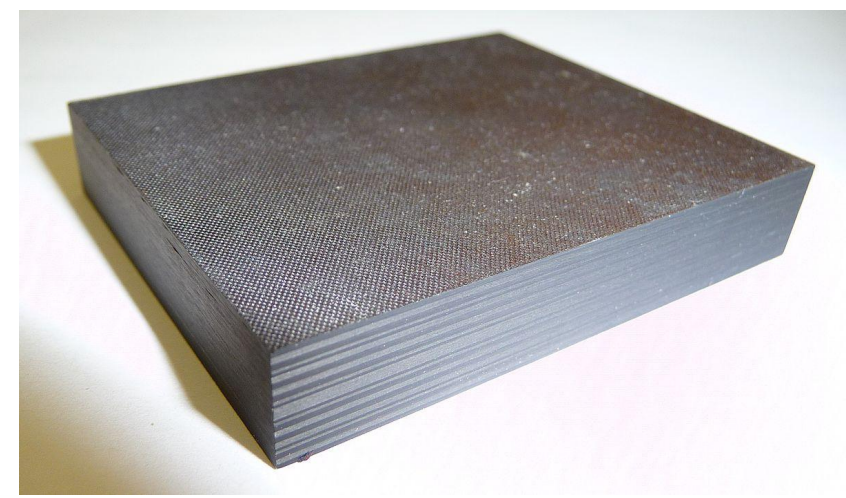
Orthotropic plate by reinforcing a simple plate with stiffeners. Faltinsen, 1999, "Water Entry of a Wedge by Hydroelastic Orthotropic Plate Theory". Journal of Ship Research 49: 180-193



General sketch of an orthotropic layer: case of an isotropic matrix reinforced by fibres. Vannucci, 2002, "A Special Planar Orthotropic Material". Journal of Elasticity 67: 81-96

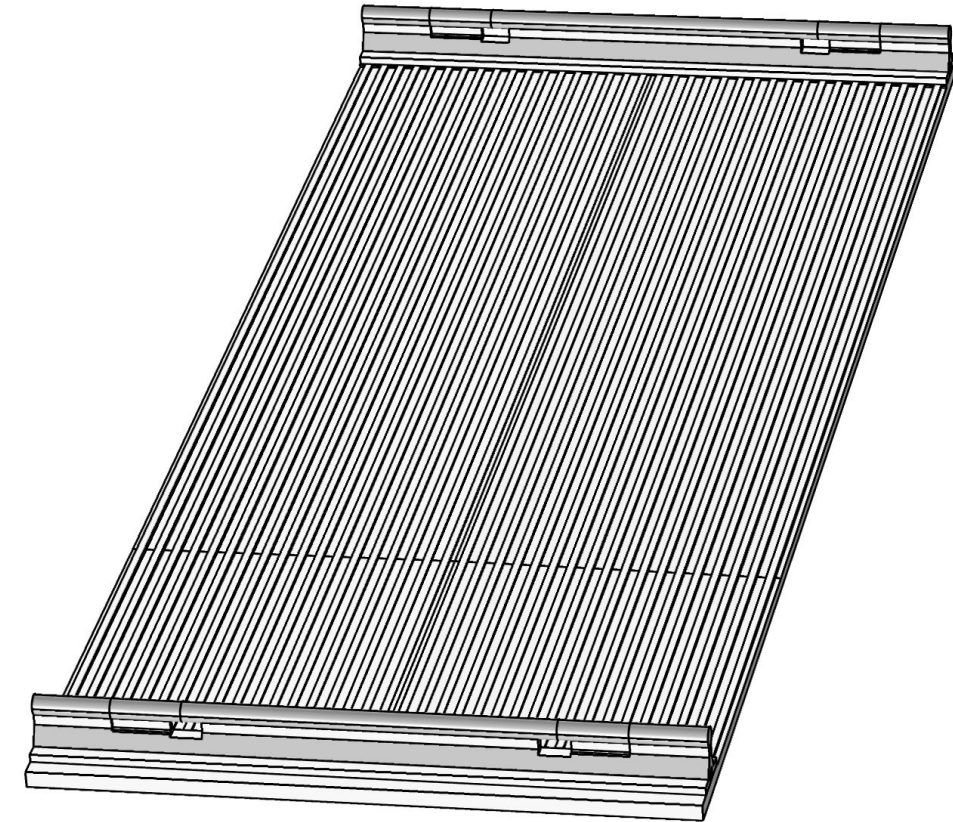


Schematic diagram of a [0/90/0/90] cross-ply laminate. Saxena, M. and Sushen Kirtania. "Stiffness analysis of symmetric cross-ply laminated composite plates." ADBU Journal of Engineering Technology (AJET) 4 (2016)



This is a photo of a small piece of laminated uni-directional Carbon Fibre. Simon.white.1000 (2012)

- MDO problem formulation using orthotropic materials
 - structural weight (W) and the strain/stress/normal force peak
 - industrial interest (weight/stress)
 - fundamental research perspective (strain/stress/normal force peak)
- Bernoulli analysis showed that the local pitch angle directly affects the spray root pressure and thus the normal force
- Controlling the plate deformation to control the local pitch angle can achieve the desired hydroelastic response
- Plate deformation can be controlled by changing the flexural rigidity along the longitudinal direction
- Use of grooves and non-uniform thickness allow to vary the longitudinal flexural rigidity

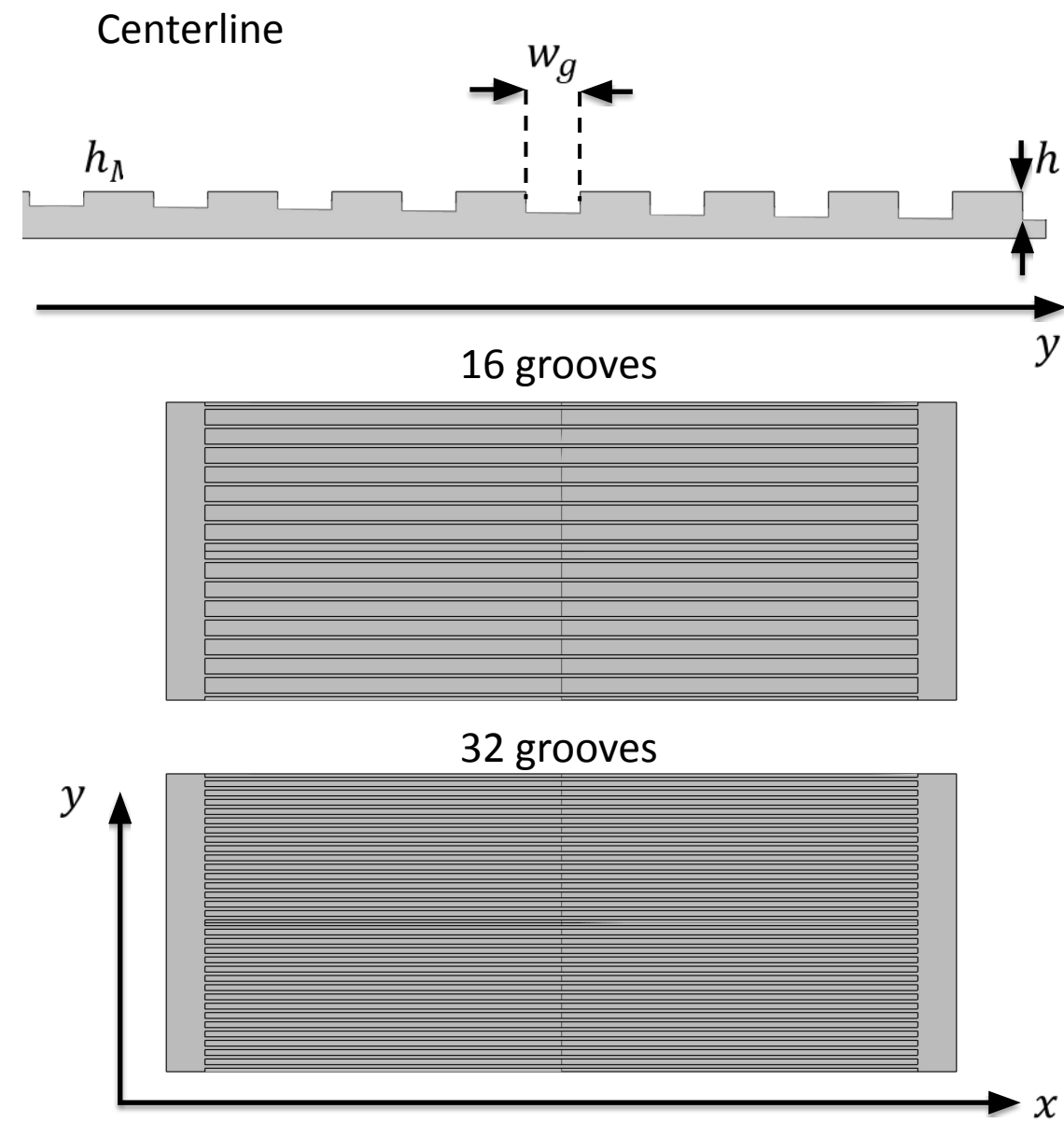


$$\text{minimize } f(\mathbf{x}) = \{W(\mathbf{x}), \max[\sigma_{Mises}(\mathbf{x}), \varepsilon_{Mises}(\mathbf{x}), \Delta F_n(\mathbf{x}), \delta_n(\mathbf{x})]\}$$

$$\text{subject to } \begin{cases} \max(\sigma) < \sigma_{max} \\ \max(\tau) < \tau_{max} \end{cases}$$

- Longitudinal thickness variations of the stiffeners is initially not allowed to obtain a plate that is equivalent to the original rectangular isotropic plate
- Domain center is defined to have a weight close to the original
- Variable ranges are defined in order to provide large design variability while maintaining a certain degree of manufacturability

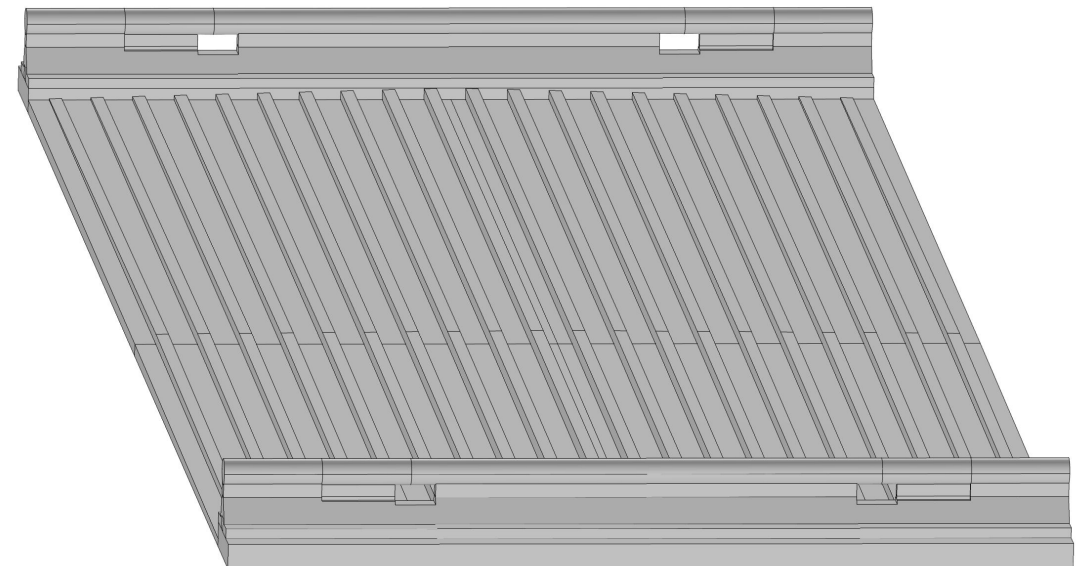
Variable	Description
1	
2	
3	
4	
5	
6	
7	



- Stiffeners thickness variations in the longitudinal direction are allowed
- Domain center is defined to have a weight close to the original

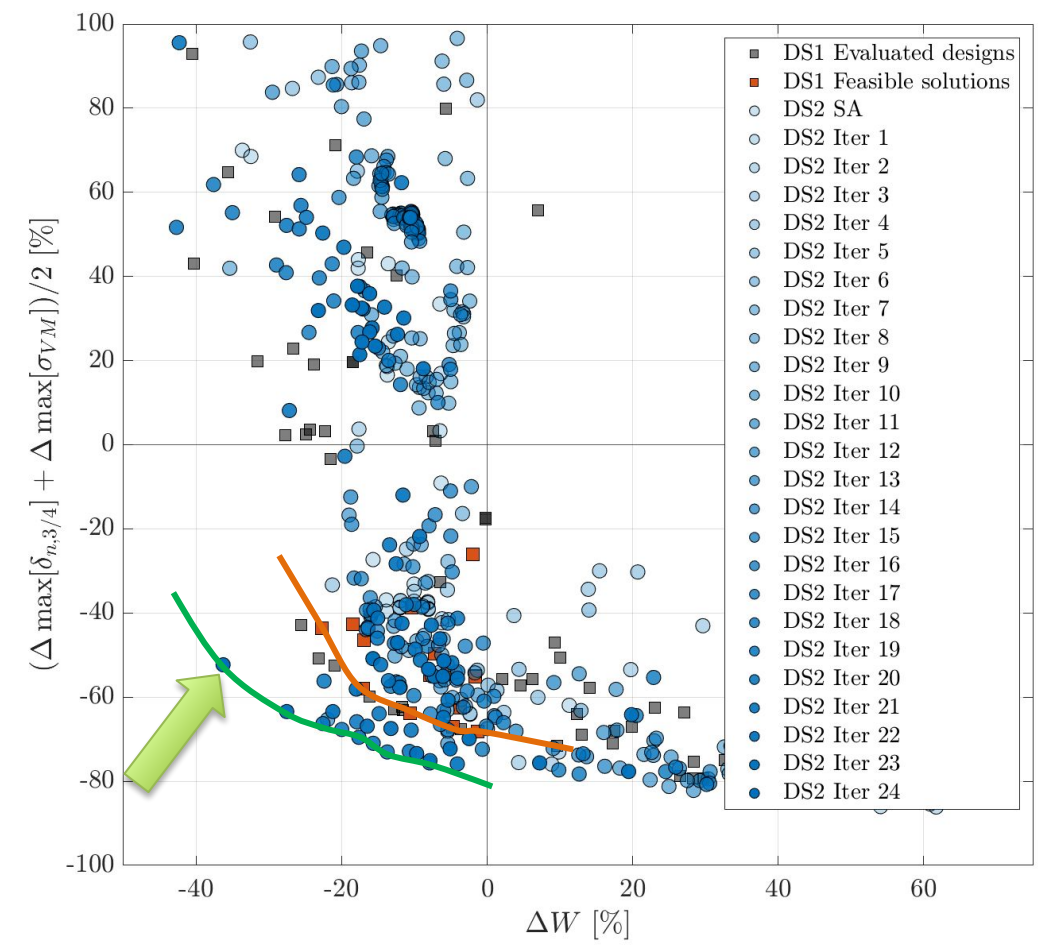
Variable	Description
1	Plate thickness – midline - xMinThick
2	Plate thickness – edge - xMinThick
3	Plate thickness – midline – TE
4	Plate thickness – edge – TE
5	Plate thickness – midline – LE
6	Plate thickness – edge – LE
7	Stiffeners thickness – midline - xMinThick
8	Stiffeners thickness – edge – xMinThick
9	Stiffeners thickness – midline – TE
10	Stiffeners thickness – edge – TE
11	Stiffeners thickness – midline – LE
12	Stiffeners thickness – edge – LE

Variable	Description
13	Longitudinal position of the minimum plate thickness (xMinThick)
14	Stiffeners width
15	Overall plate thickness
16	Number of stiffeners

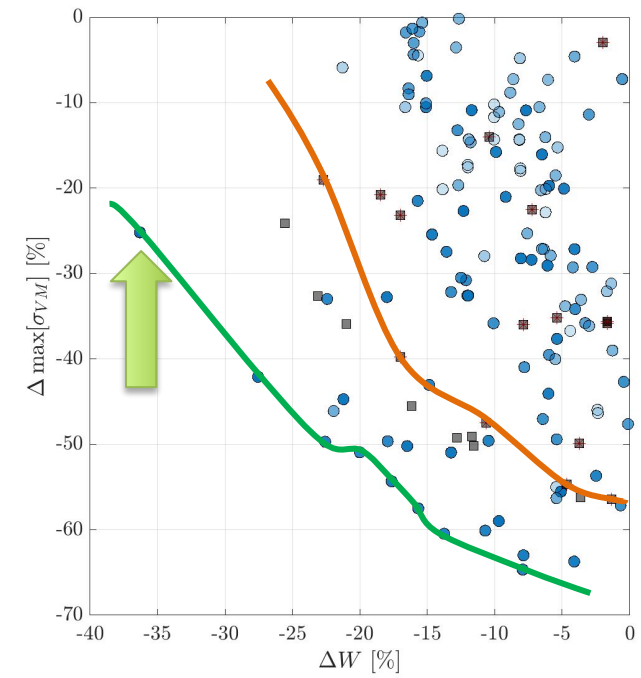


- The optimization on Design Space 2 achieve better optimal results than Design Space 1
 - 489 one-way simulations have been performed
 - Using a larger design space allowed a significant improvement of the plate performance
 - The entire Pareto set has been improved

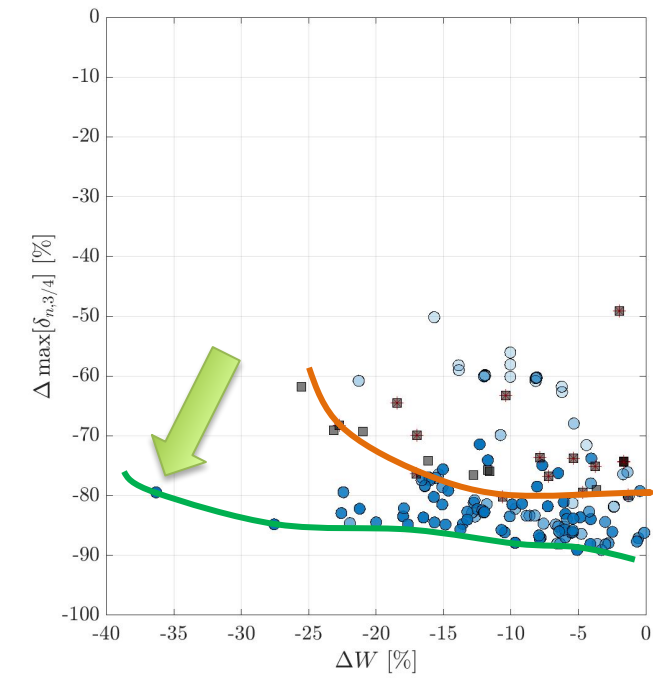
Aggregated results



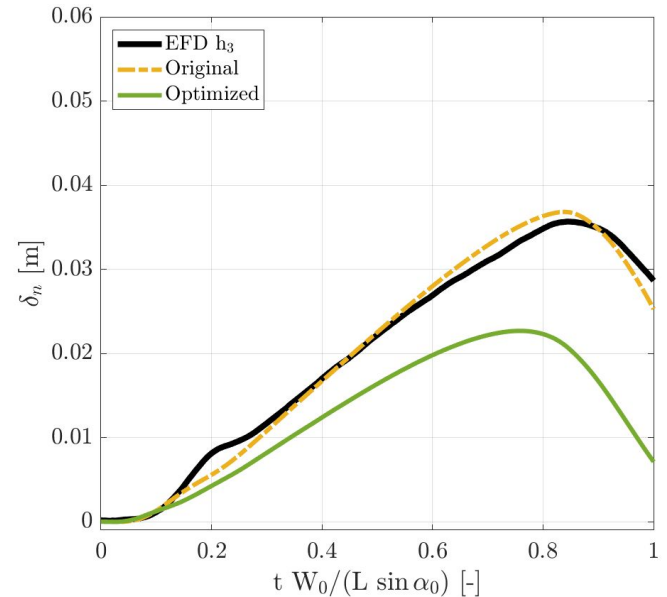
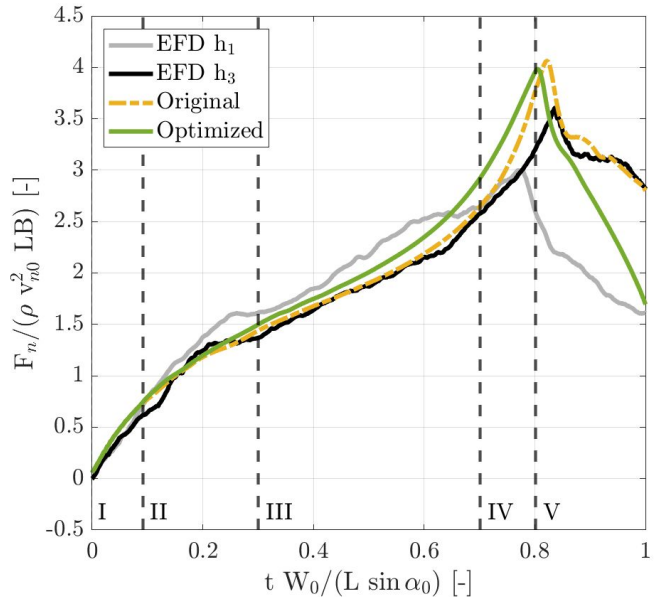
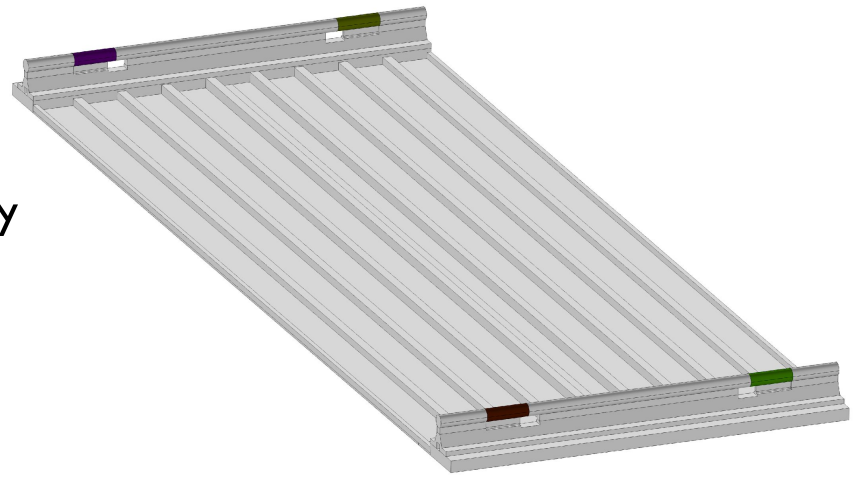
Maximum normal deflection



Maximum Von Mises stress



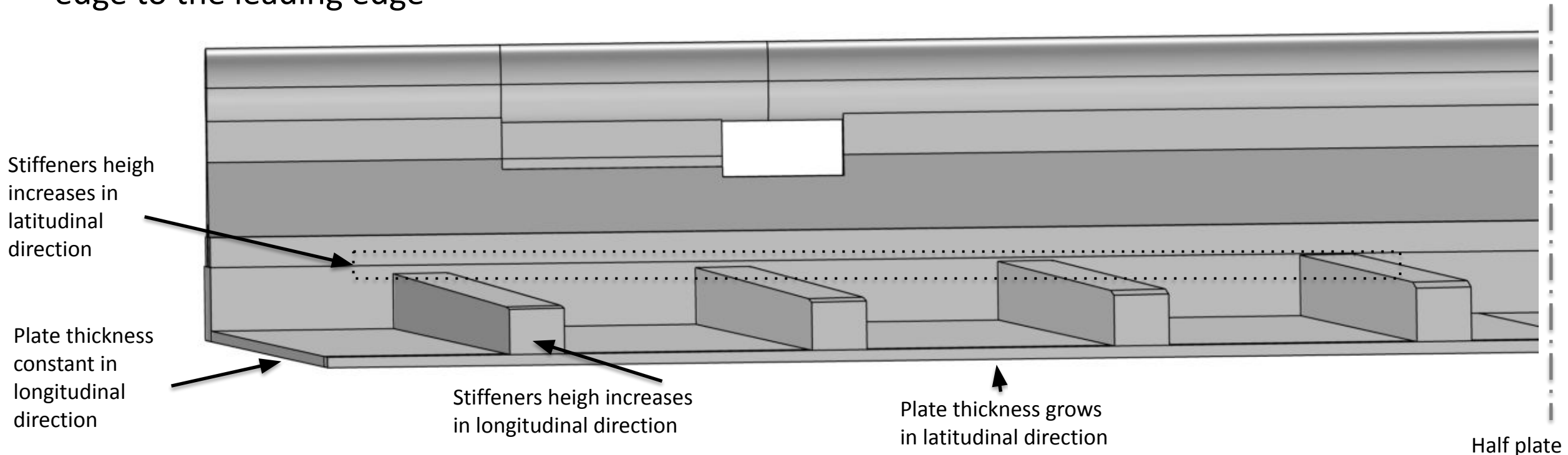
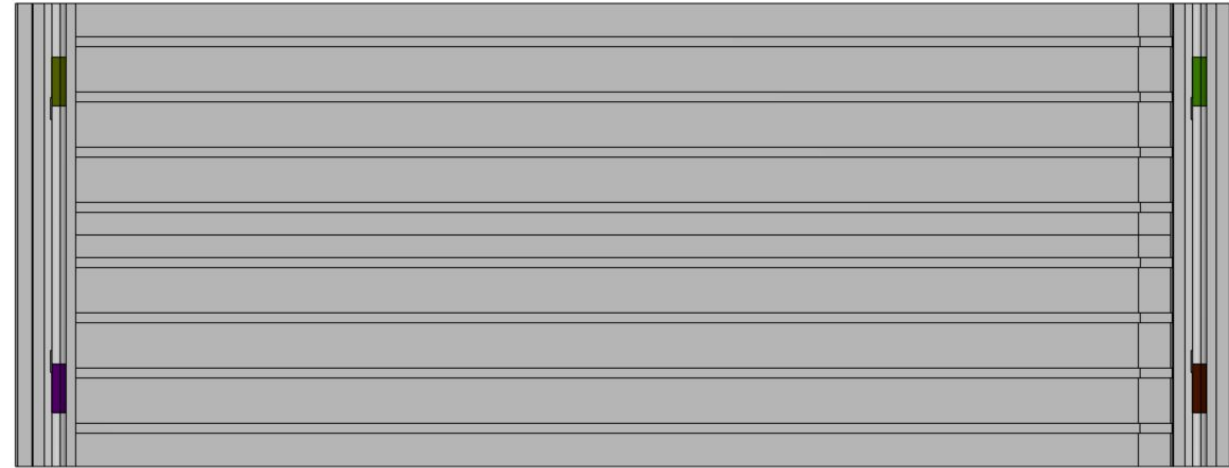
- ❑ Two-way nonlinear FSI simulation of the optimized plate performed with CFDShip V4.5
- ❑ The normal force peak is 1.2% smaller than the original plate and slightly anticipated
- ❑ The normal deflection is reduced by 38% w.r.t. to the original plate



Variations with respect to the original plate

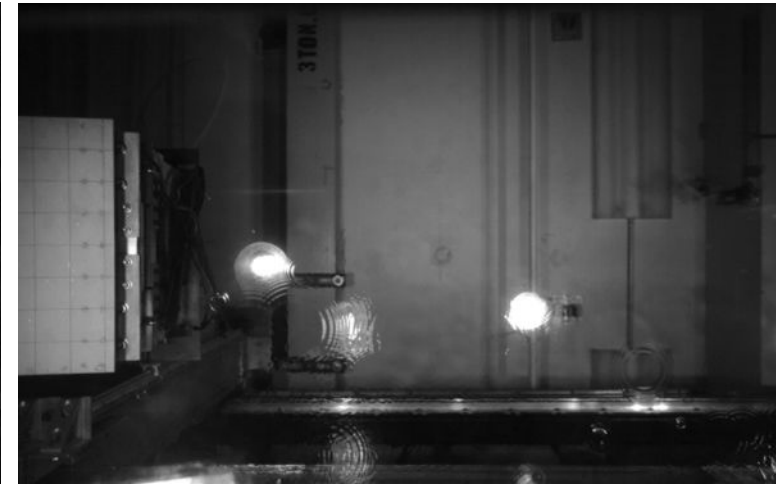
Plate		
DS1 candidate 1W	-22%	-15%
DS2 candidate 1W	-36%	-45%
DS2 candidate 2W validated	-36%	-38%

- Current candidate solution
 - Eight stiffeners
 - Minimum plate thickness: 1.7 mm
 - Maximum plate thickness: 2 mm
- The thickness of the optimized plate grows from the external sides towards the midline
- The stiffeners thickness grows from the trailing edge to the leading edge



- Oblique slam with large horizontal velocity component
- 7.5 deg pitch angle
 - Previous was 10 deg
- During slam, air entrainment can significantly affect the resultant forces acting on a structure [1]
 - Positive effect: it can distribute the loads over a larger surface
 - Negative effect: it can increase the duration of the loading, and cause additional damage to a solid structure by an oscillating air pocket or a cavitation implosion of the trapped air

Thickness [mm]	Pitch [deg]	U [m/s]	W [m/s]	U/W	Vn [m/s] (Fr)
0.00635	7.5	7.375	0.427	17.3	1.390 (0.426)

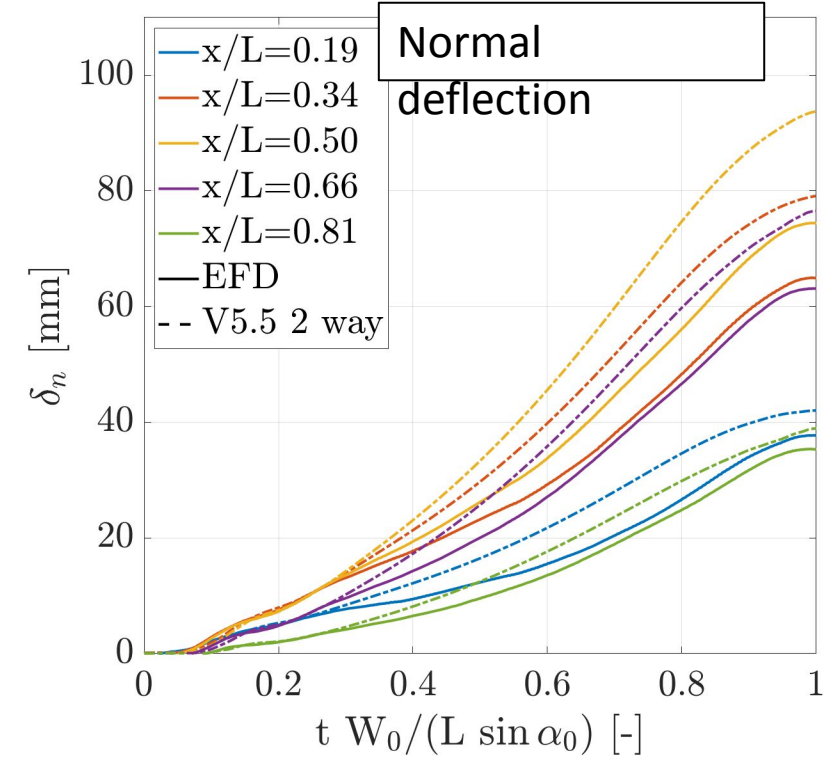
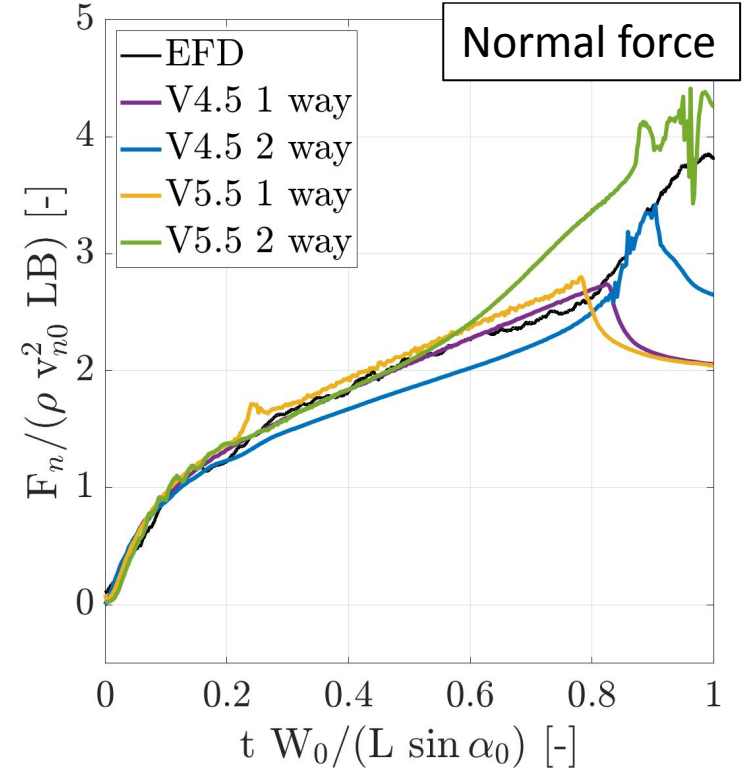


Video from UMD QPR presentation

[1] Jain, U., Vega-Martínez, P., & Van Der Meer, D. (2021). Air entrapment and its effect on pressure impulses in the slamming of a flat disc on water. *Journal of Fluid Mechanics*, 928, A31.

- During slam, air entrapment can significantly affect the resultant forces acting on a structure
- Investigate fast oblique slamming with low pitch angle using CFDShip V5.5
- Oblique slam with large horizontal velocity component is investigated
 - 7.5 deg pitch angle (previous was 10 deg)
- Preliminary results show that the trends are correct with overprediction of the normal force and structural deflection

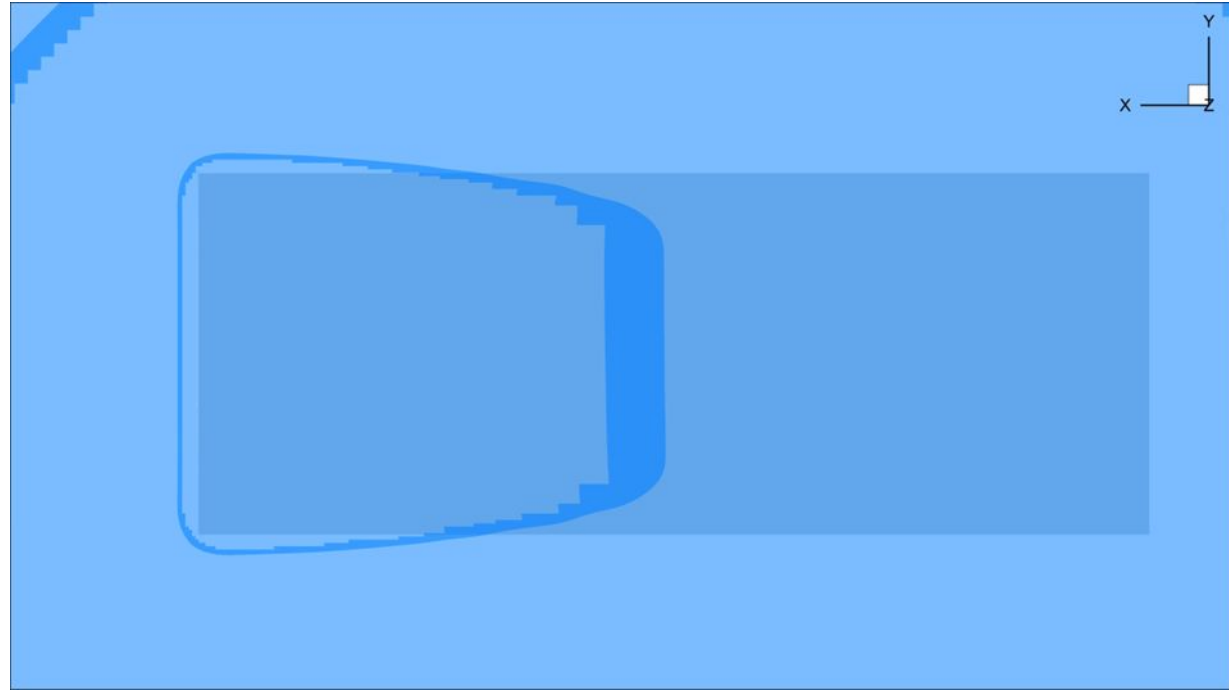
Thickness [mm]	Pitch [deg]	U [m/s]	W [m/s]	U/W	Vn [m/s] (Fr)
0.00635	7.5	7.375	0.427	17.3	1.390 (0.426)



- Although force and deflection prediction lacks high accuracy, the air entrainment is well captured



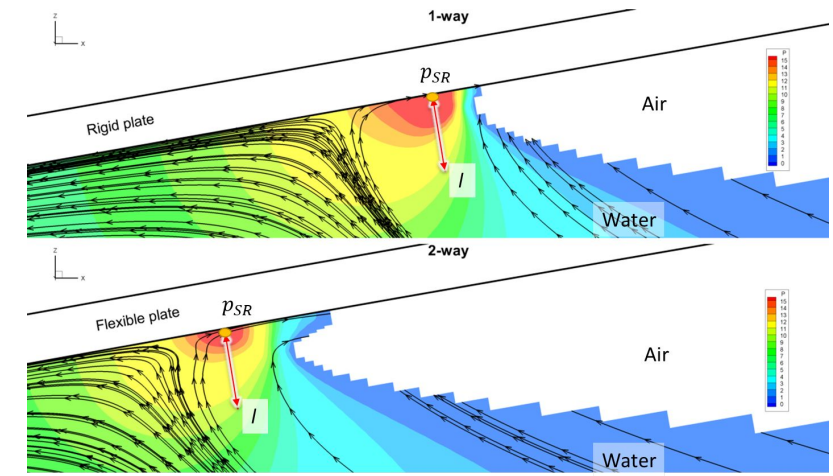
EFD: bottom view of the slam, courtesy of UMD



CFD/FSI: bottom view of the slam

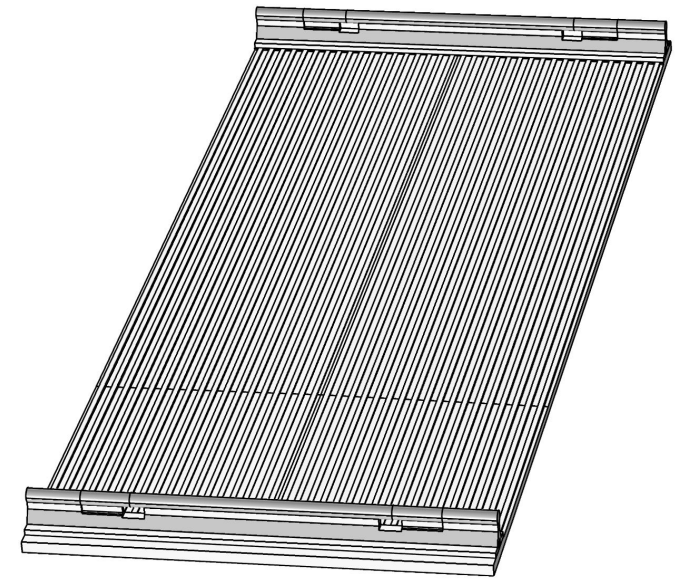
Conclusions

- Available numerical and experimental data have been exploited to investigate several aspects of the slamming of elastic plates on quiescent water
 - Methods required for the effective analysis of the slamming of elastic solids (two-way FSI with nonlinear structural solver)
 - The role of hydroelasticity and its use to select relevant cases for the effective investigation of the slamming physics
 - The use of a stagnation flow model to investigate the role of the plate kinematics on the impact force
 - The analysis of the conservation of energy to identify an effective strategy for the multidisciplinary optimization for weight reduction and safety increase
 - The effect of the plate aspect ratio on the slamming physics
- The MDO of a macroscopically orthotropic plate has been performed



Future research

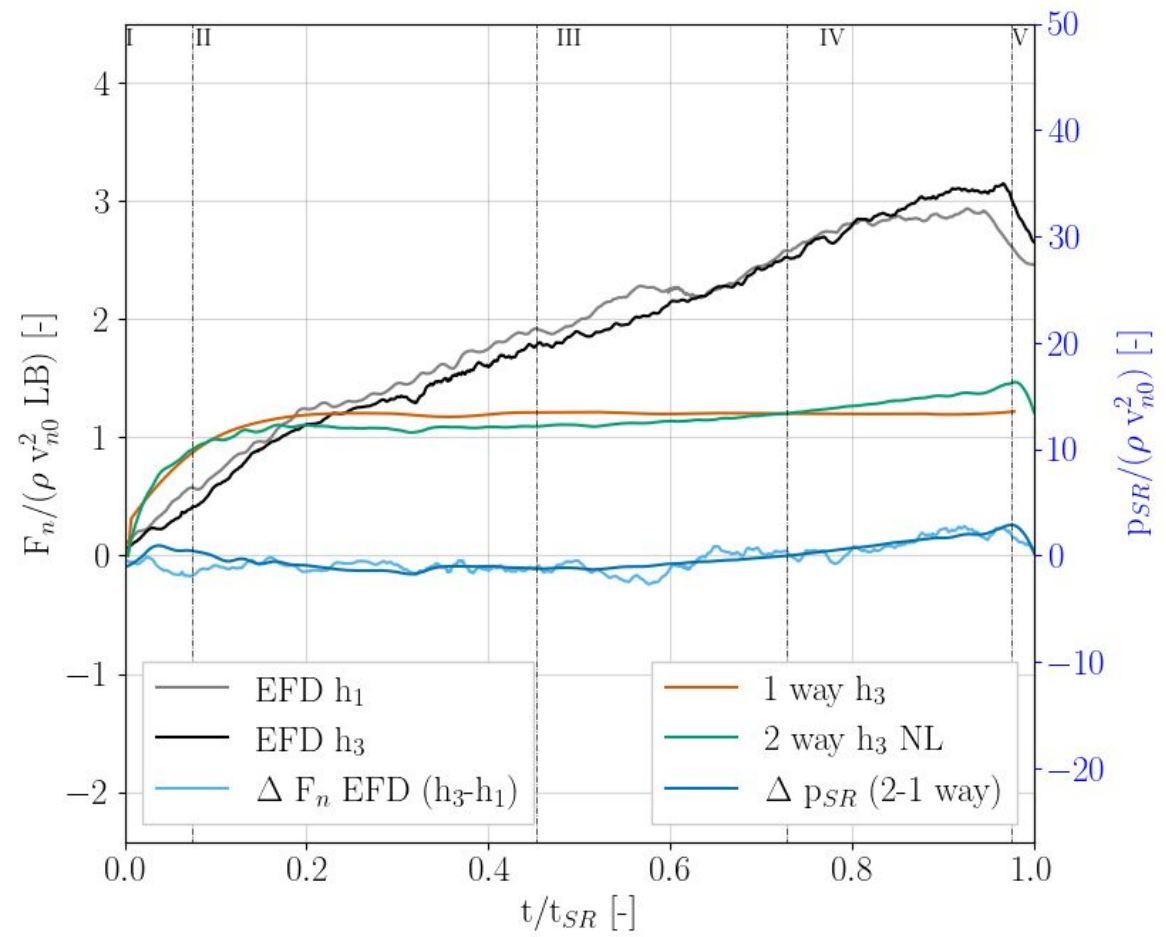
- Flexible plate slamming with waves
- High-fidelity CFD/CSD FSI multi-phase flow simulation



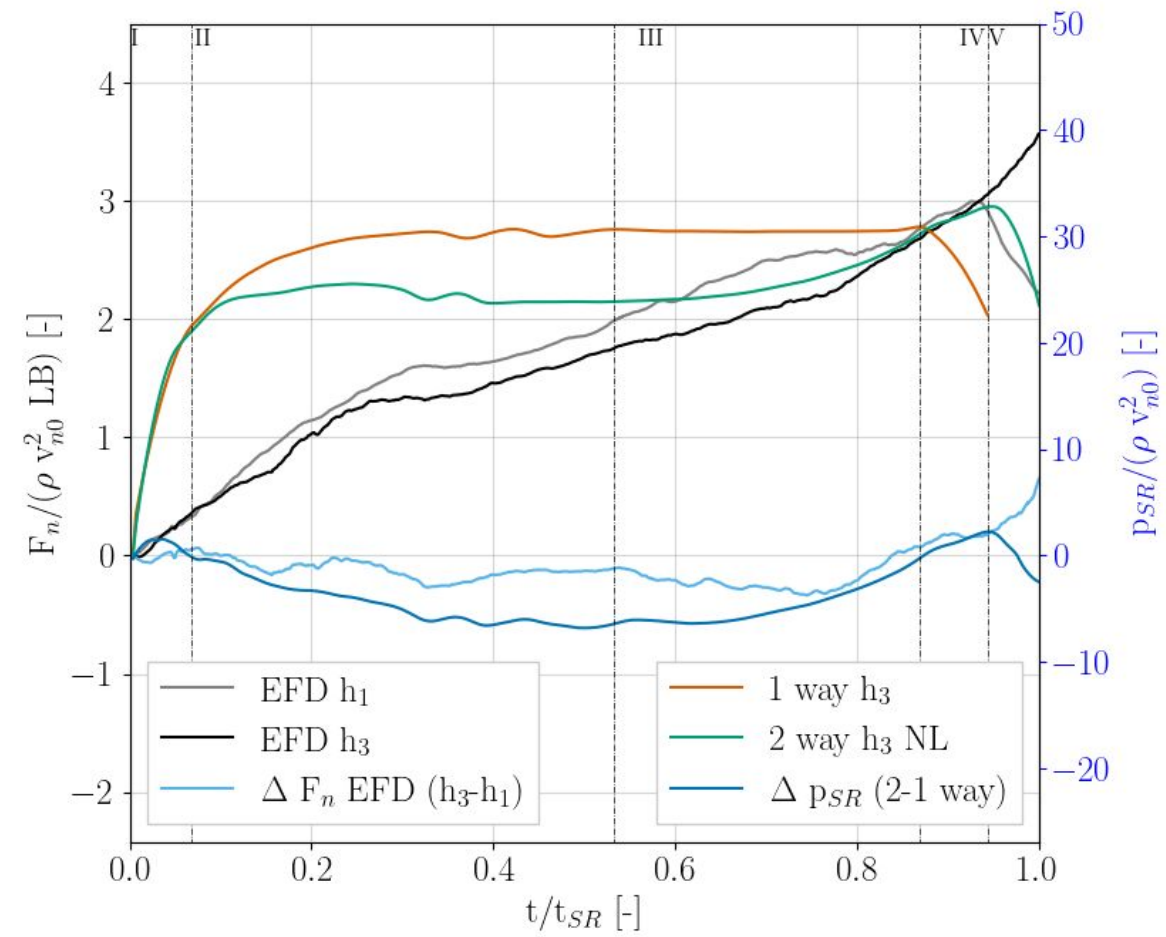
- Faltinsen, Odd M. "The effect of hydroelasticity on ship slamming." *Philosophical Transactions of the Royal Society of London. Series A: Mathematical, Physical and Engineering Sciences* 355, no. 1724 (1997): 575-591.
- Faltinsen, Odd M. "Water entry of a wedge by hydroelastic orthotropic plate theory." *Journal of ship research* 43, no. 03 (1999): 180-193.
- Faltinsen, Odd M., Jan Kvålsvold, and Jan V. Aarsnes. "Wave impact on a horizontal elastic plate." *Journal of Marine Science and Technology* 2, no. 2 (1997): 87-100.
- Huang, J., Carrica, P.M., and Stern, F., 2008. "Semi-coupled air/water immersed boundary approach for curvilinear dynamic overset grids with application to ship hydrodynamics". *International Journal for Numerical Methods in Fluids*, 58(6): 591-624.
- Iafrati, A. "Experimental investigation of the water entry of a rectangular plate at high horizontal velocity." *Journal of Fluid Mechanics* 799 (2016): 637-672.
- Panciroli, Riccardo, and M. Porfiri. "Analysis of hydroelastic slamming through particle image velocimetry." *Journal of Sound and Vibration* 347 (2015): 63-78.
- Segletes, S.B., Walters, W.P., "A note on the application of the extended Bernoulli equation", *International Journal of Impact Engineering*, 27 (2002) 561–576.
- Stenius, Ivan, Anders Rosén, and Jacob Kutteneuler. "Explicit FE-modelling of hydroelasticity in panel-water impacts." *International Shipbuilding Progress* 54, no. 2-3 (2007): 111-127.
- Stenius, Ivan, Anders Rosén, and Jakob Kutteneuler. "Hydroelastic interaction in panel-water impacts of high-speed craft." *Ocean Engineering* 38, no. 2-3 (2011): 371-381.
- Volpi, S., M. Diez, H. Sadat-Hosseini, D-H. Kim, F. Stern, R. S. Thodal, and J. L. Grenestedt. "Composite bottom panel slamming of a fast planing hull via tightly coupled fluid-structure interaction simulations and sea trials." *Ocean Engineering* 143 (2017): 240-258.
- Wang, A., Wong, K.P., Yu, M., Kiger, K.T. and Duncan, J.H., 2022. The controlled impact of elastic plates on a quiescent water surface. *Journal of Fluid Mechanics*, 939, p.A4.
- Wang, A., Kim, H. T., Wong, K. P., Yu, M., Kiger, K. T., & Duncan, J. H. (2019). "Spray formation and structural deflection during the oblique impact of a flexible plate on a quiescent water surface". *Journal of Ship Research*, 63(3), 154-164.
- Zhao, R. and Faltinsen, O., 1993. Water entry of two-dimensional bodies. *Journal of fluid mechanics*, 246, pp.593-612.
- Wu, G.X. and Sun, S.L., 2014. Similarity solution for oblique water entry of an expanding paraboloid. *Journal of fluid mechanics*, 745, pp.398-408.
- Judge, C., Troesch, A. and Perlin, M., 2004. Initial water impact of a wedge at vertical and oblique angles. *Journal of Engineering Mathematics*, 48, pp.279-303.
- Panciroli, R., 2013. Water entry of flexible wedges: Some issues on the FSI phenomena. *Applied Ocean Research*, 39, pp.72-74.
- Khabakhpasheva, T.I. and Korobkin, A.A., 2013. Elastic wedge impact onto a liquid surface: Wagner's solution and approximate models. *Journal of Fluids and Structures*, 36, pp.32-49.
- Korobkin, A.A. and Khabakhpasheva, T.I., 2006. Regular wave impact onto an elastic plate. *Journal of engineering mathematics*, 55, pp.127-150.

Backup slides – Bernulli equation

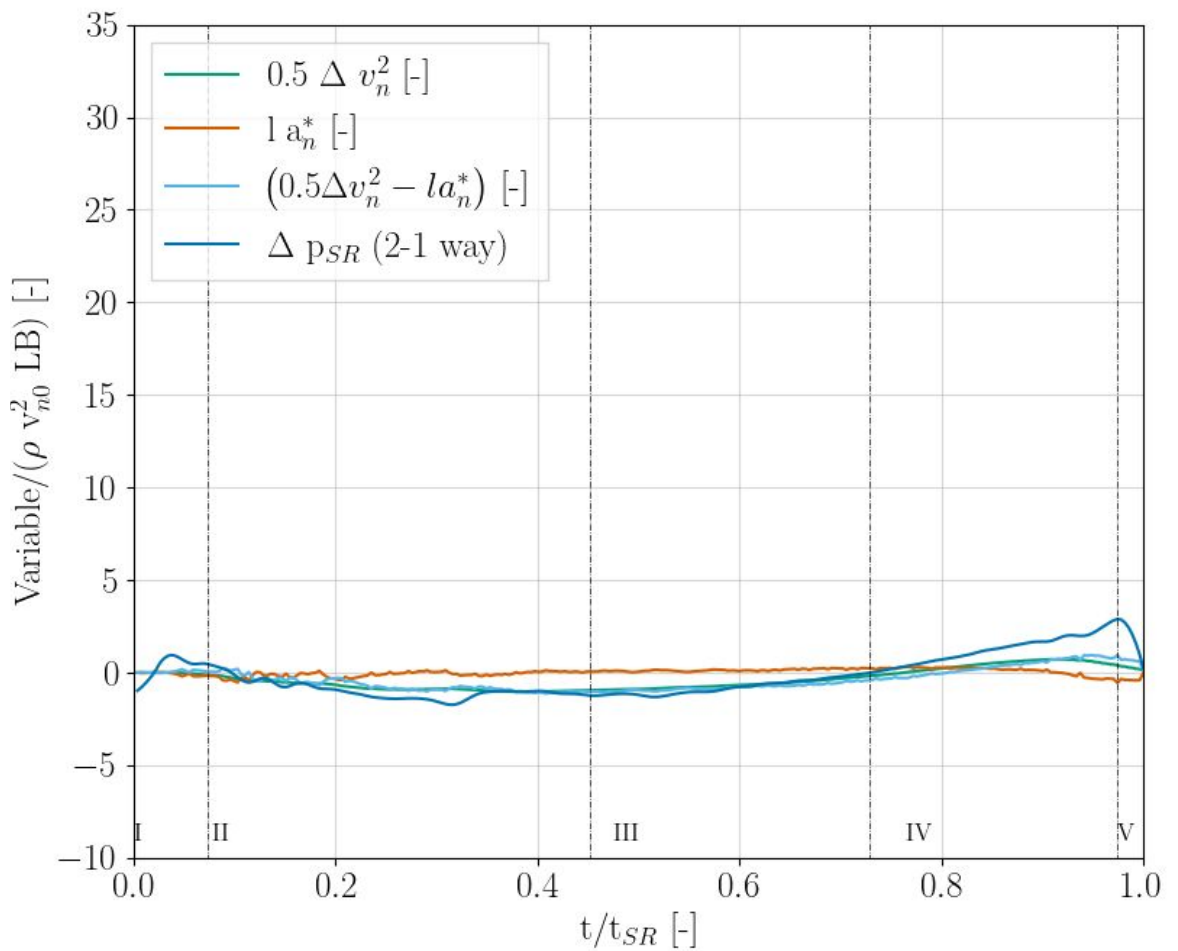
0.269



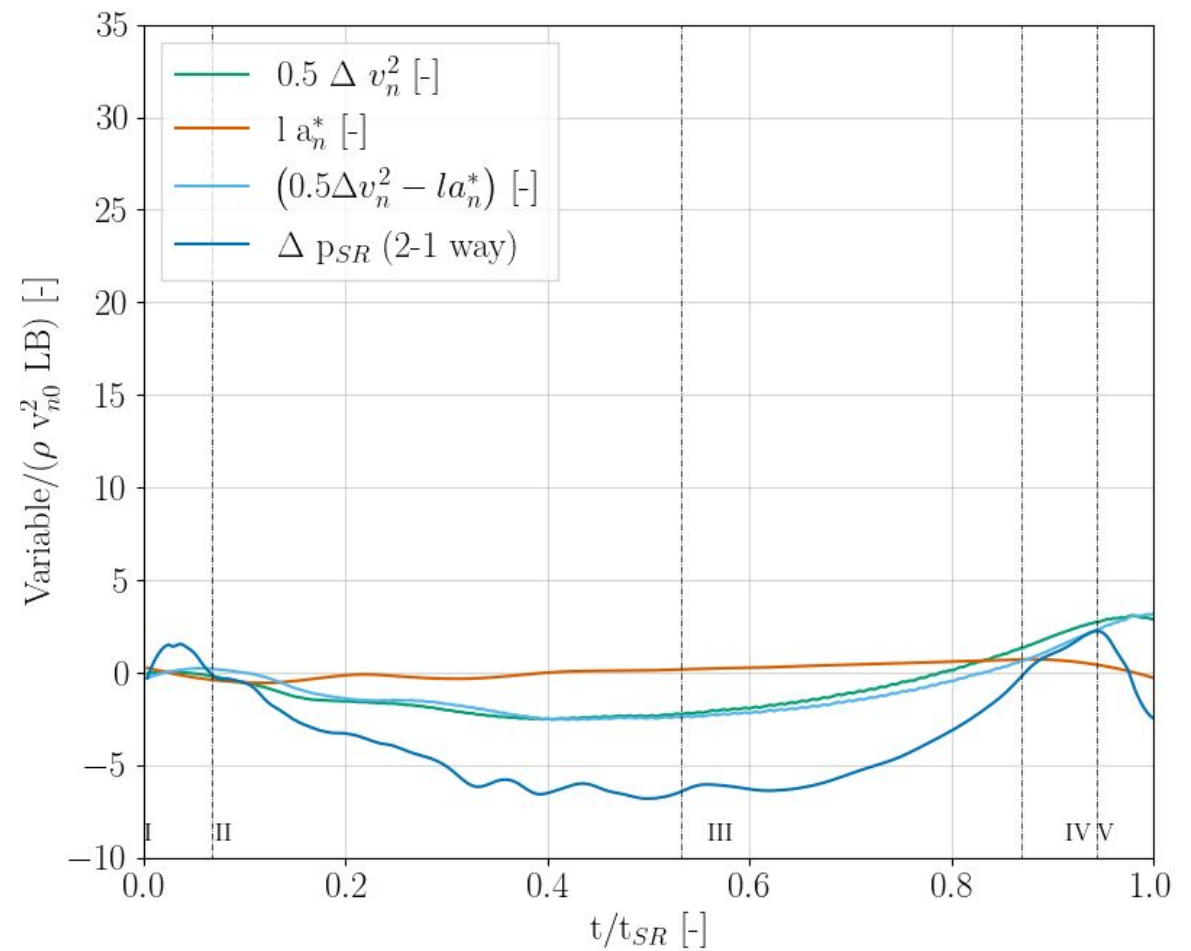
0.403



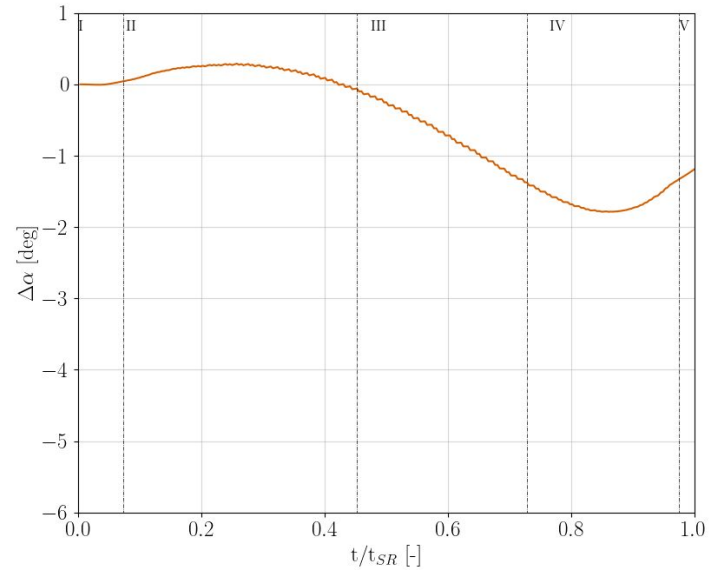
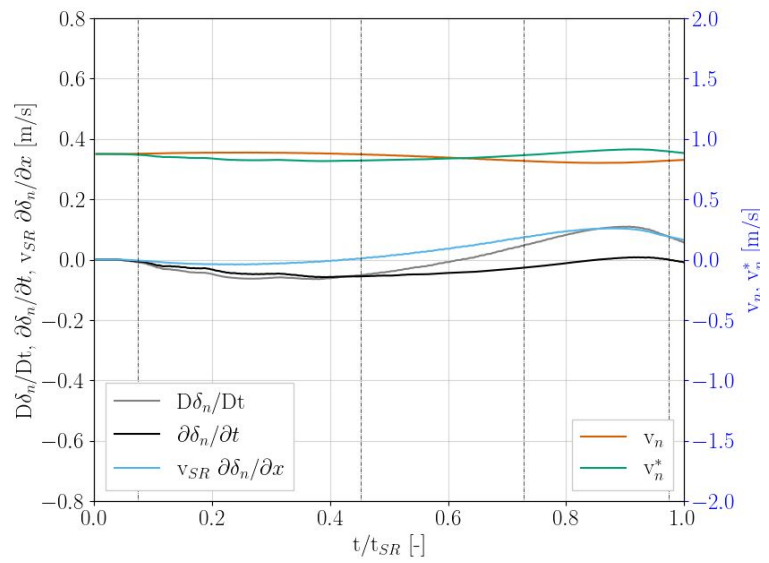
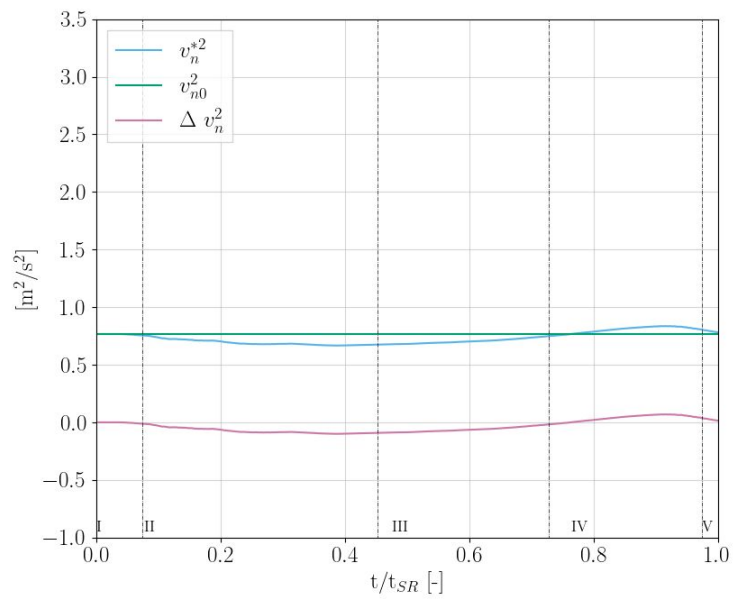
0.269



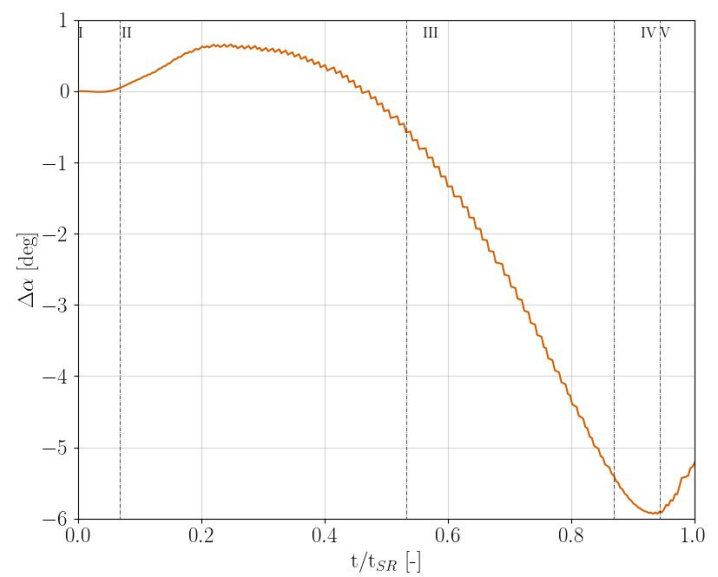
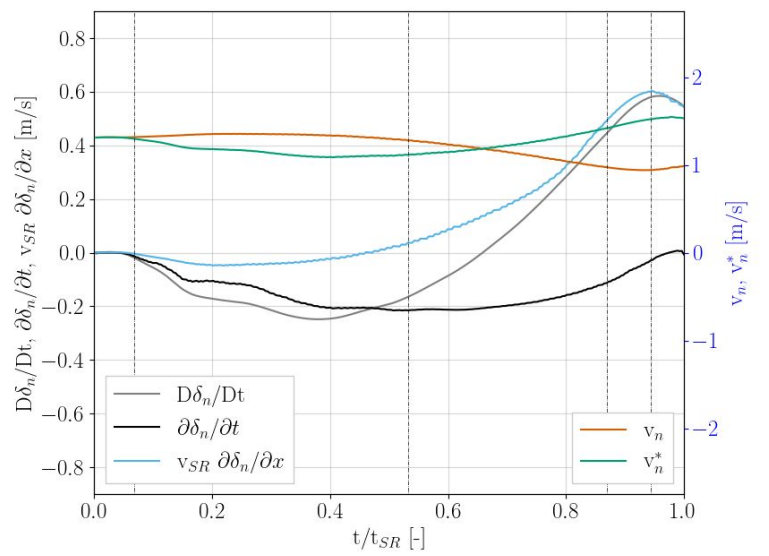
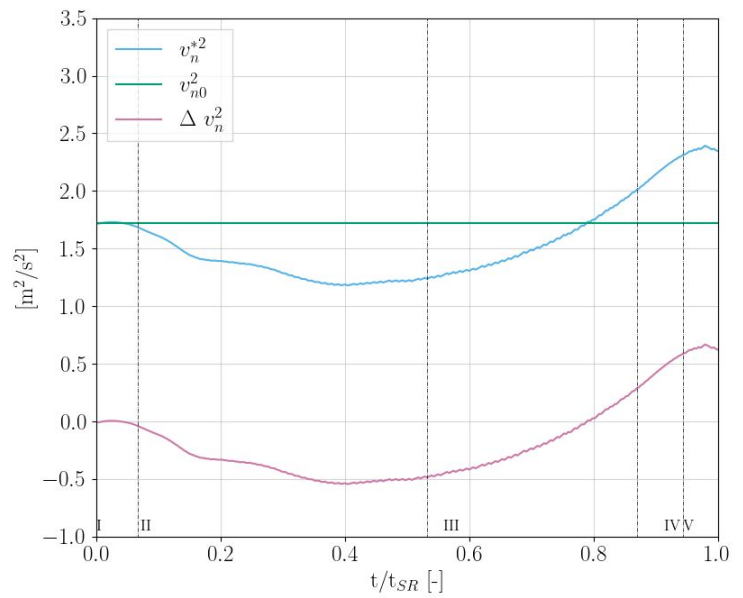
0.403



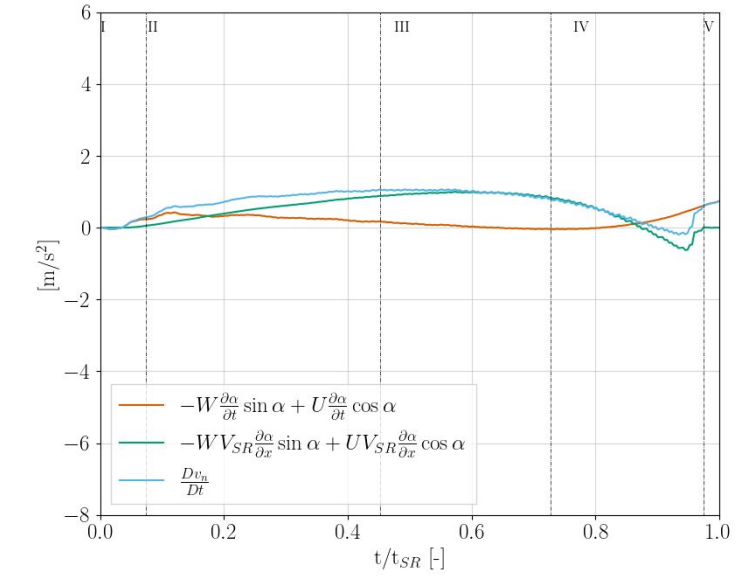
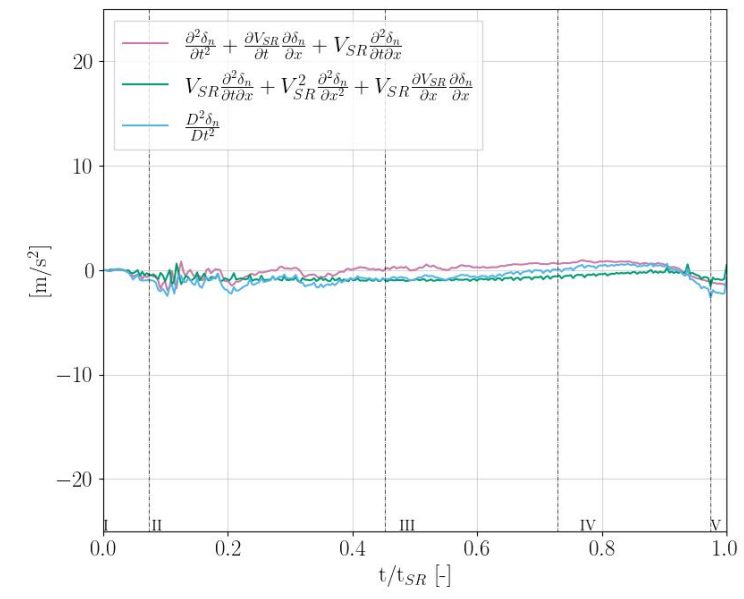
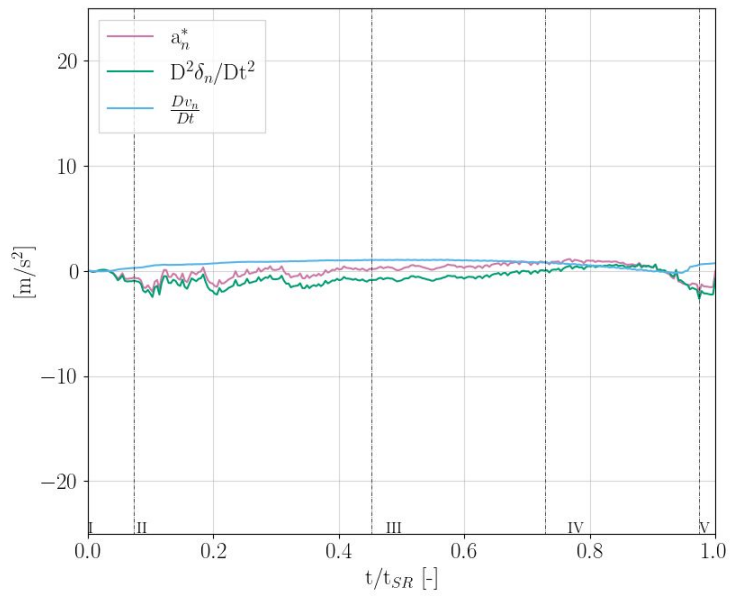
0.269



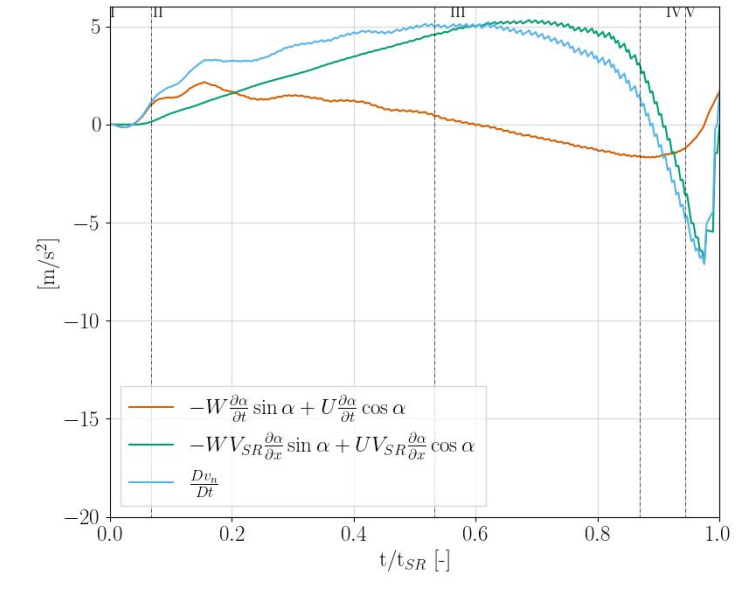
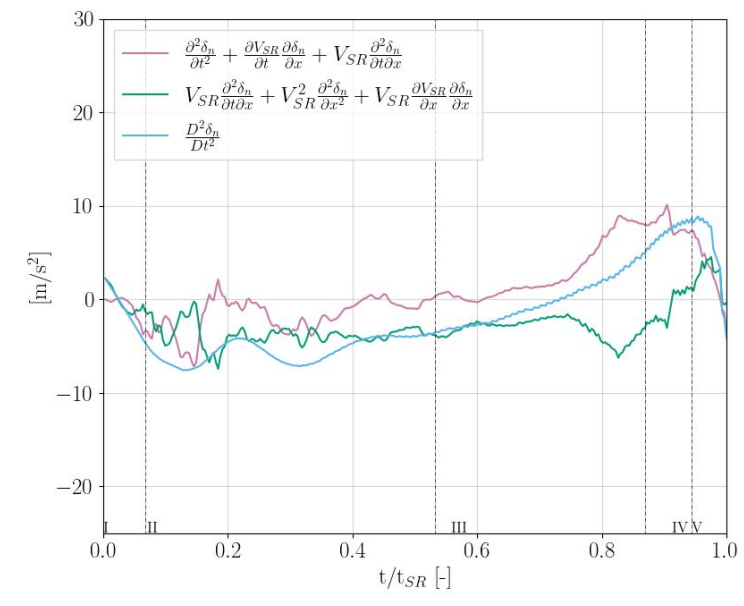
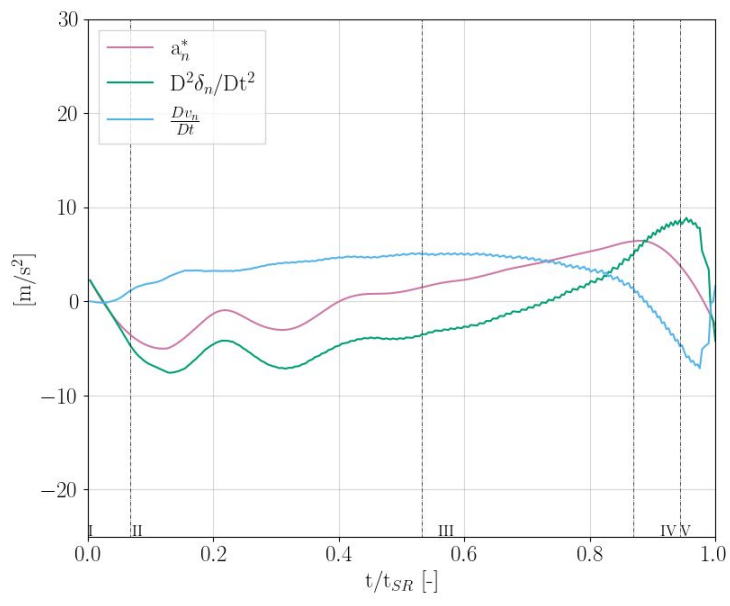
0.403



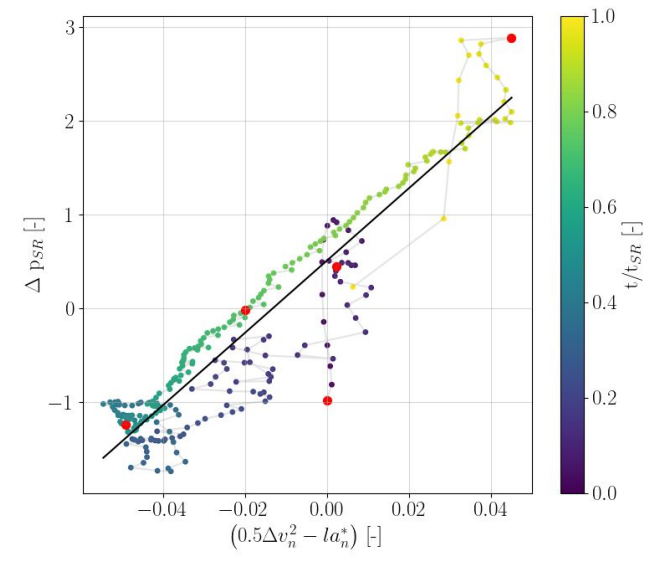
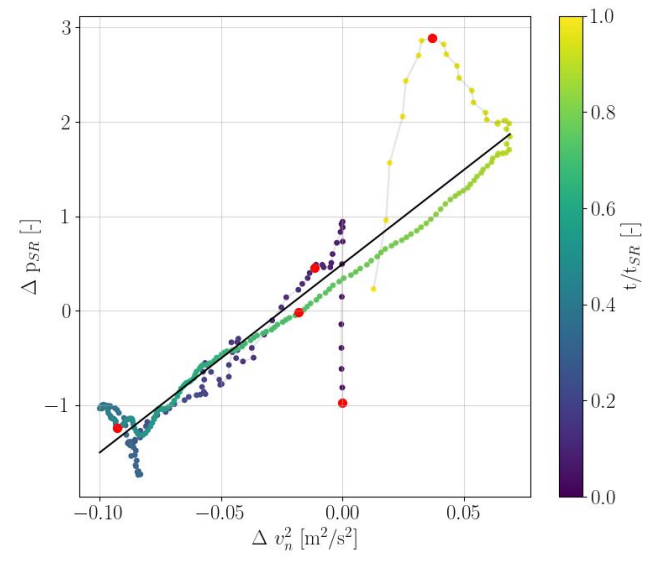
0.269



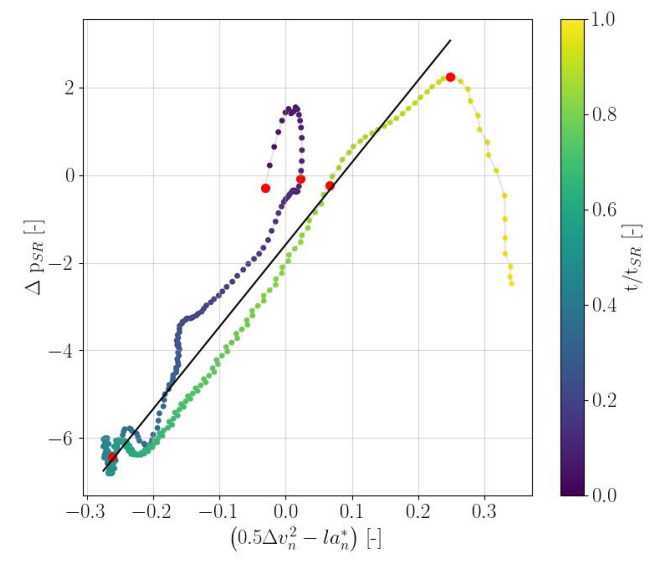
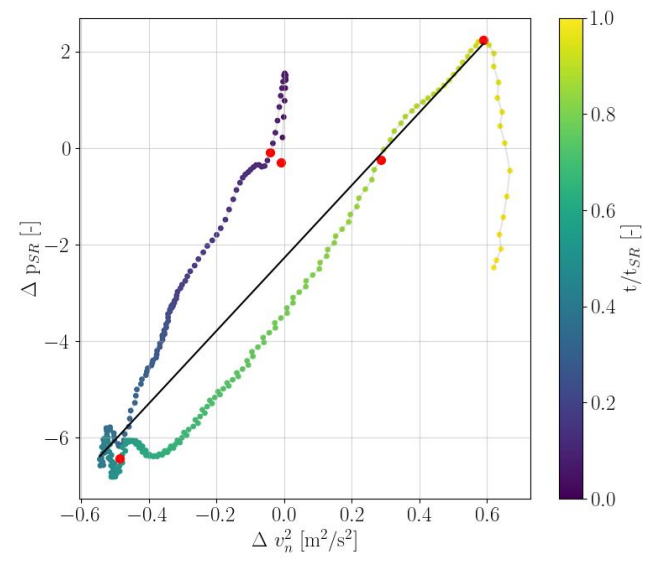
0.403



0.269



0.403



Backup slides – Energy conservation O.269

- Rigid plate

- The mount transfers energy directly to the fluid since the plate is rigid

Water

$$\frac{\partial}{\partial t} \iiint_{V_w} \rho_w \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}(t)} p(v_n) dS = W_{pR}$$

Structure

$$-W_{pR} = W_M$$

- Flexible plate

- The mount transfers energy to the fluid and to the plate (kinetic and potential elastic energy) since the plate is deforming

Water

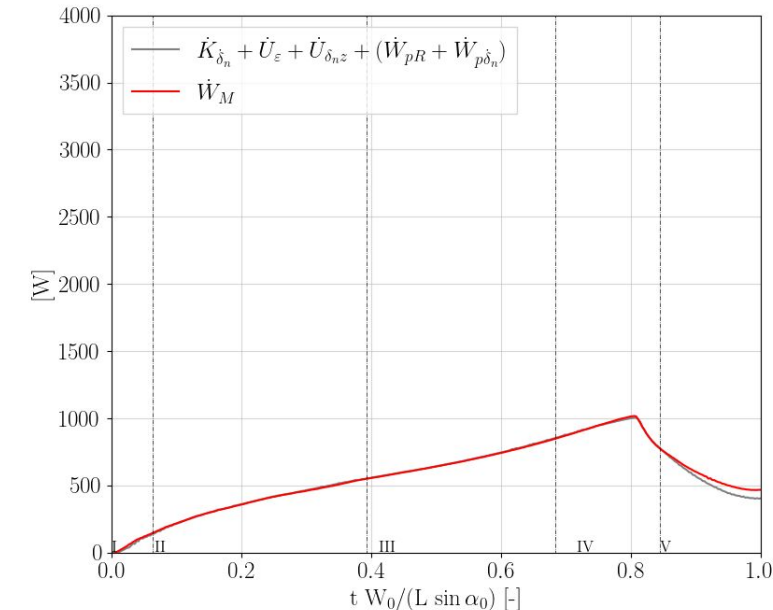
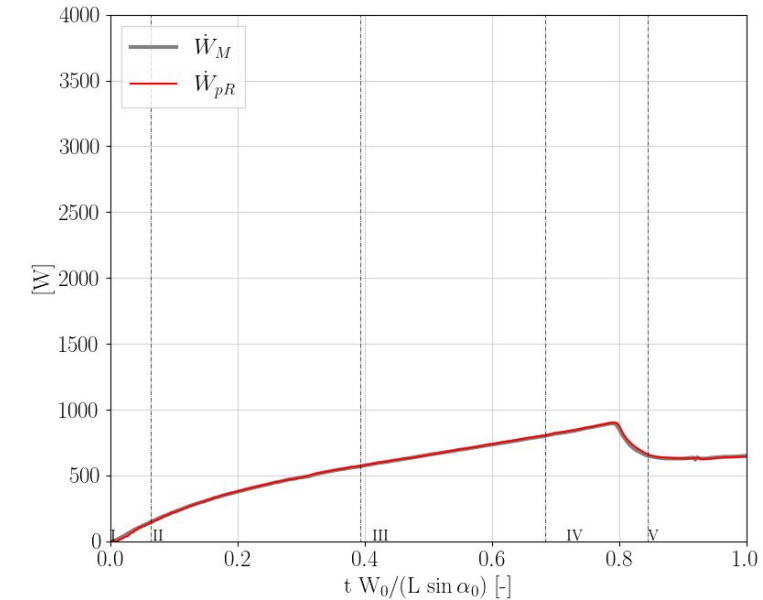
$$\frac{\partial}{\partial t} \iiint_{V_w} \rho \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}} p[v_n + \dot{\delta}_n] dS = W_{pR} + W_{p\dot{\delta}_n}$$

Structure

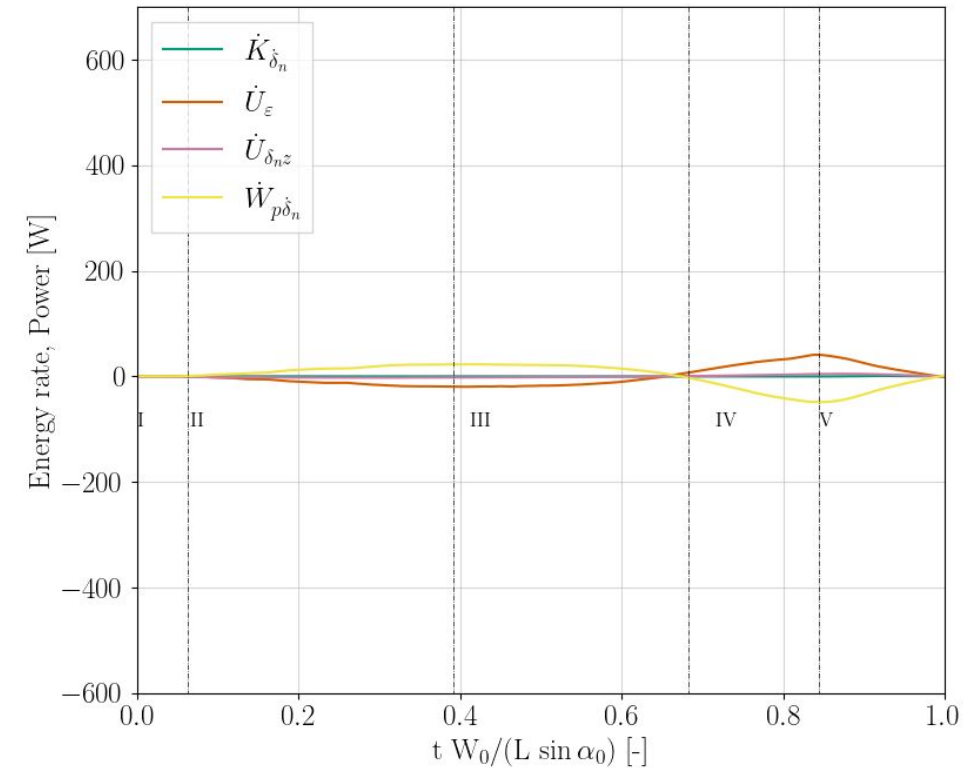
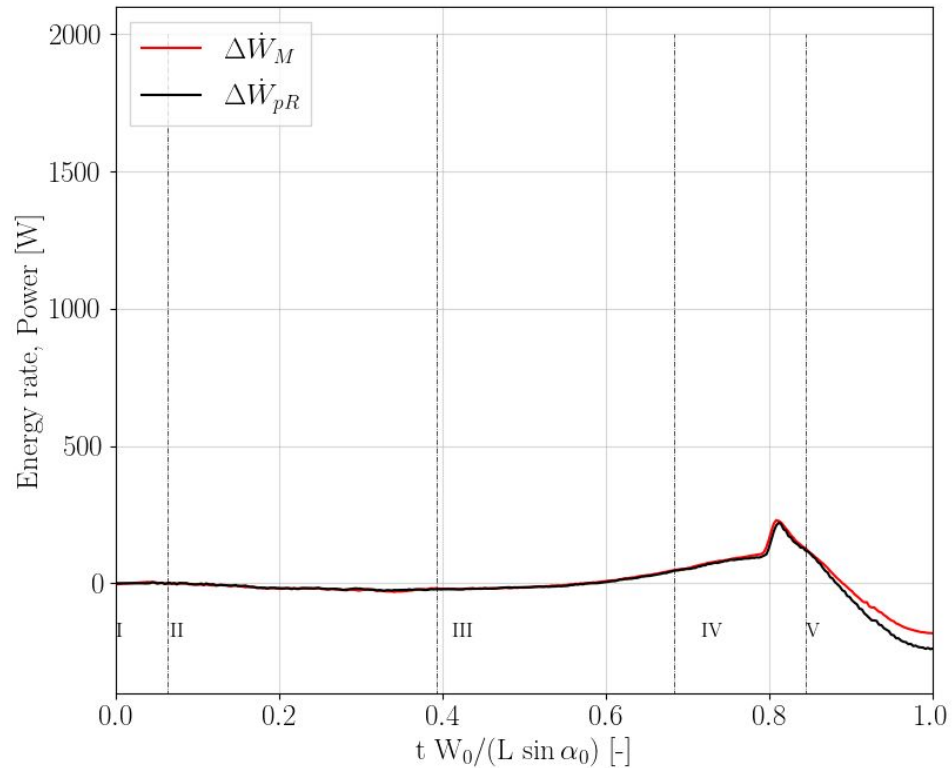
$$\dot{K}_{\dot{\delta}_n} + \dot{U}_\varepsilon - (W_{pR} + W_{p\dot{\delta}_n}) = W_M$$

$$v_n = \mathbf{v} \cdot \hat{\mathbf{n}}$$

$$\delta_n = \boldsymbol{\delta} \cdot \hat{\mathbf{n}}$$

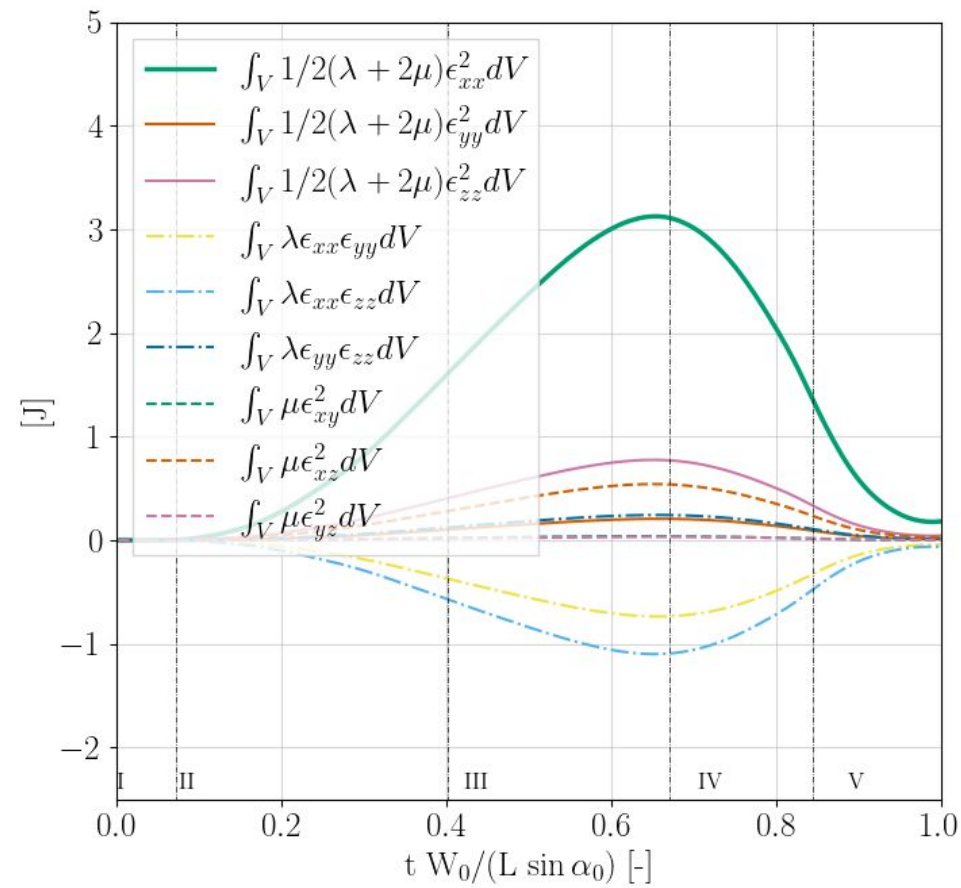


- The difference equation is satisfied
- The elastic potential energy almost equals the power provided by the pressure due to the flexibility
- The kinetic energy content is the smallest

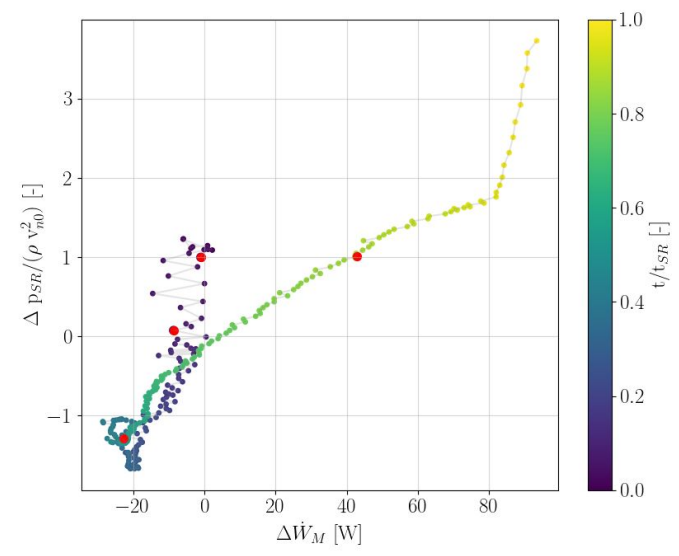
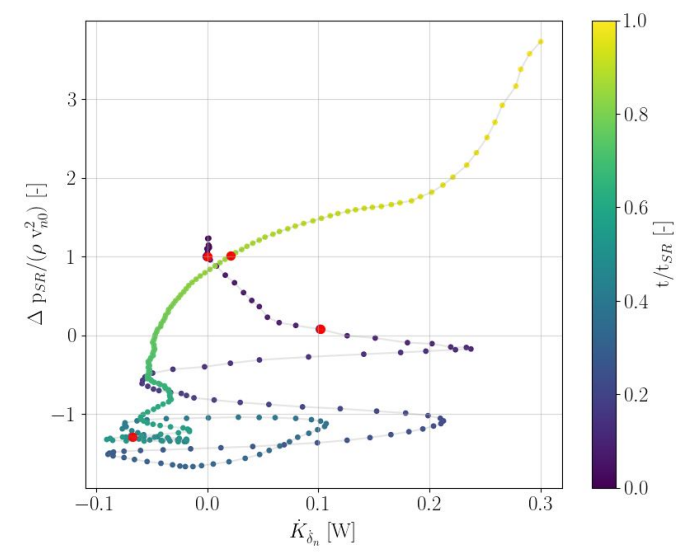
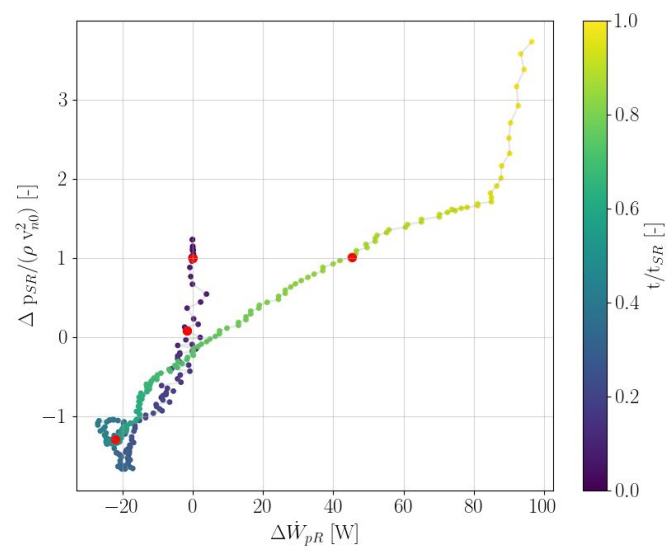
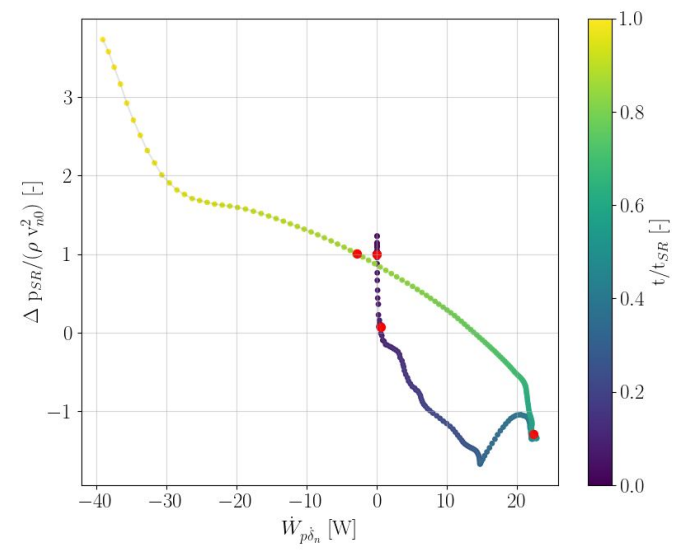
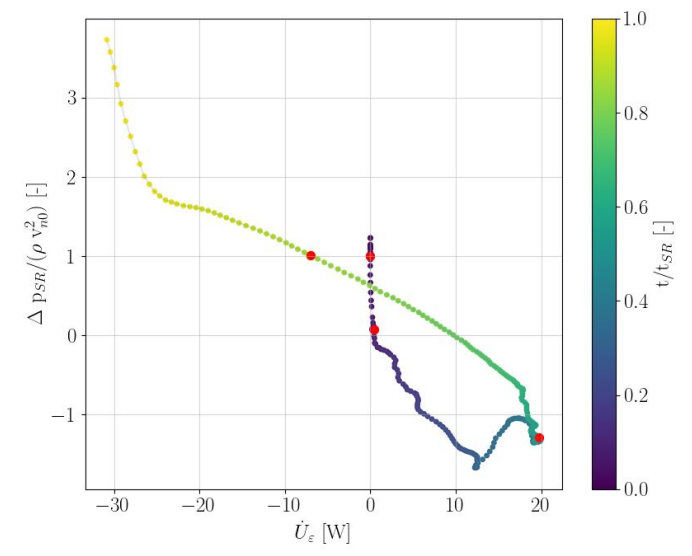


3. Conservation of Energy Analysis (10): Elastic Energy Components

- For all the cases the ϵ_{xx}^2 strain contribution is the most important.
- The second most important term is $\epsilon_{xx}\epsilon_{zz}$ since for the T-bars the ϵ_{zz} component is the most important (due to their orientation)



- A good correlation between Δp_{SR} and the elastic potential energy exists phase wise
- Δp_{SR} correlates well also with ΔW_{pR} and ΔW_M
- Δp_{SR} does not correlate well with the kinetic energy variation



Backup slides – Energy conservation O.403

- Rigid plate

- The mount transfers energy directly to the fluid since the plate is rigid

Water

$$\frac{\partial}{\partial t} \iiint_{V_w} \rho_w \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}(t)} p(v_n) dS = W_{pR}$$

Structure

$$-W_{pR} = W_M$$

- Flexible plate

- The mount transfers energy to the fluid and to the plate (kinetic and potential elastic energy) since the plate is deforming

Water

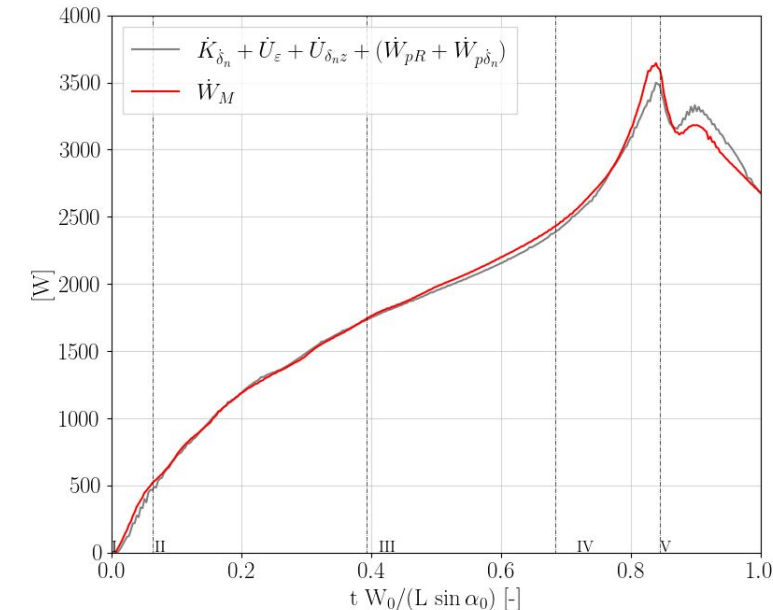
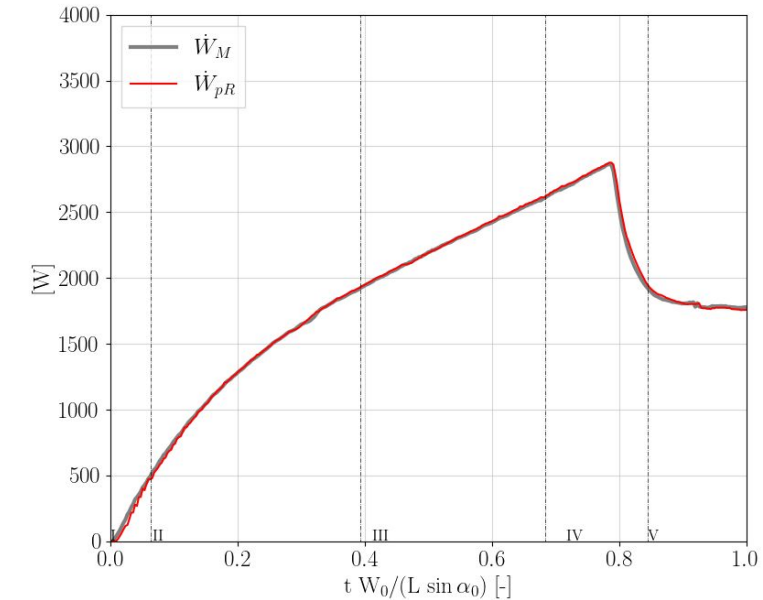
$$\frac{\partial}{\partial t} \iiint_{V_w} \rho \left(\frac{\|\mathbf{u}_w\|^2}{2} \right) dV_w = \iint_{S_{Wet}} p[v_n + \dot{\delta}_n] dS = W_{pR} + W_{p\dot{\delta}_n}$$

Structure

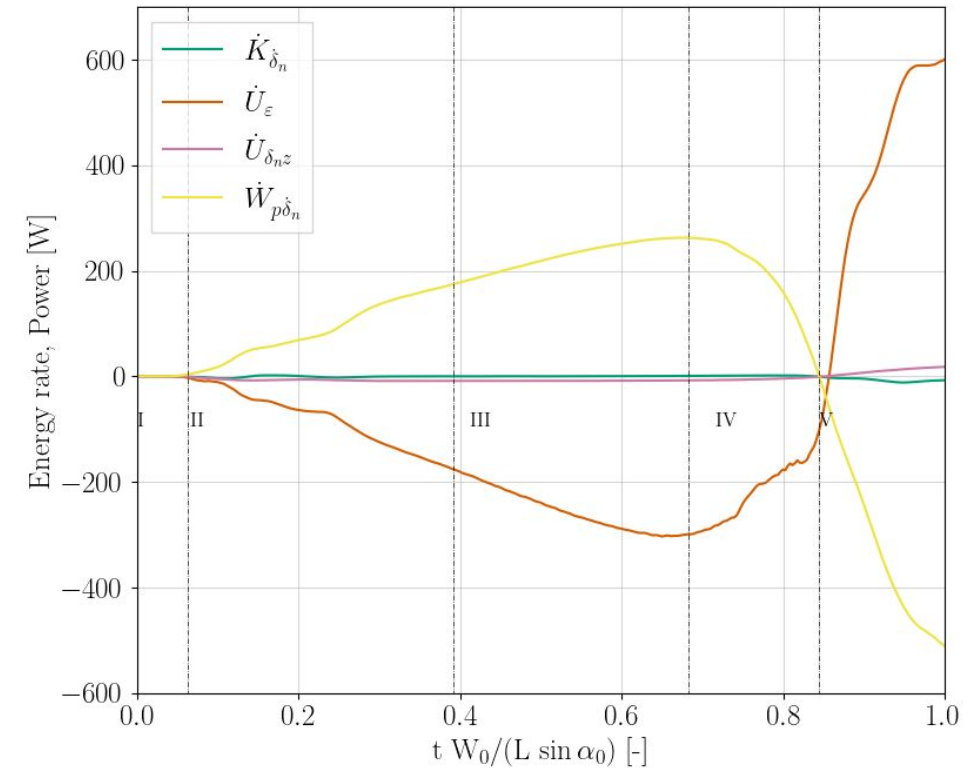
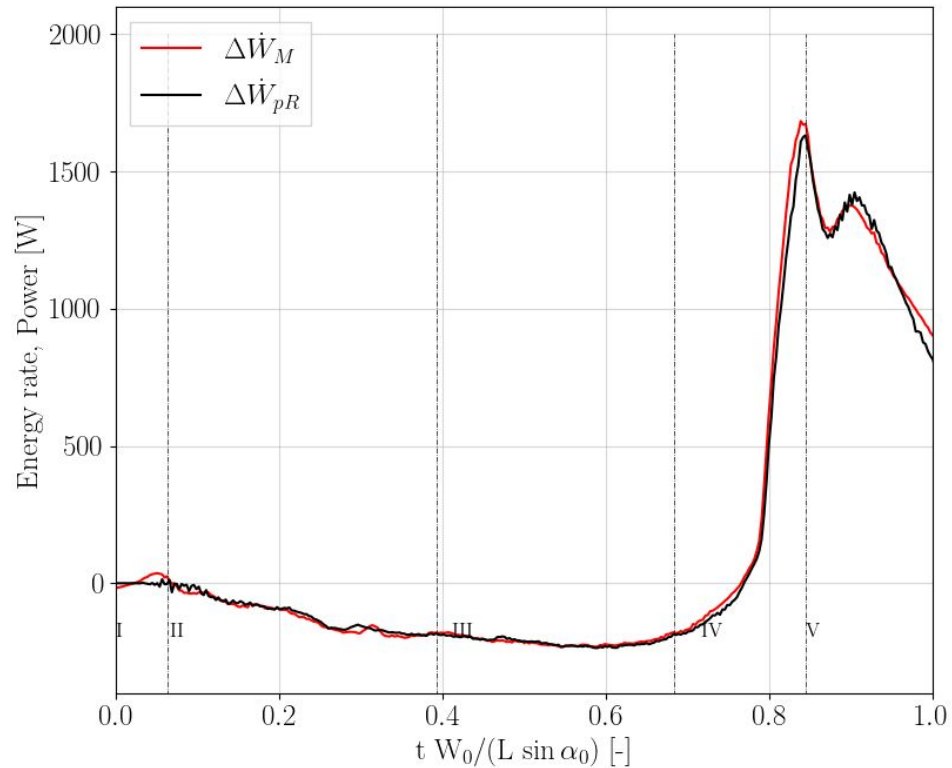
$$\dot{K}_{\dot{\delta}_n} + \dot{U}_\varepsilon - (W_{pR} + W_{p\dot{\delta}_n}) = W_M$$

$$v_n = \mathbf{v} \cdot \hat{\mathbf{n}}$$

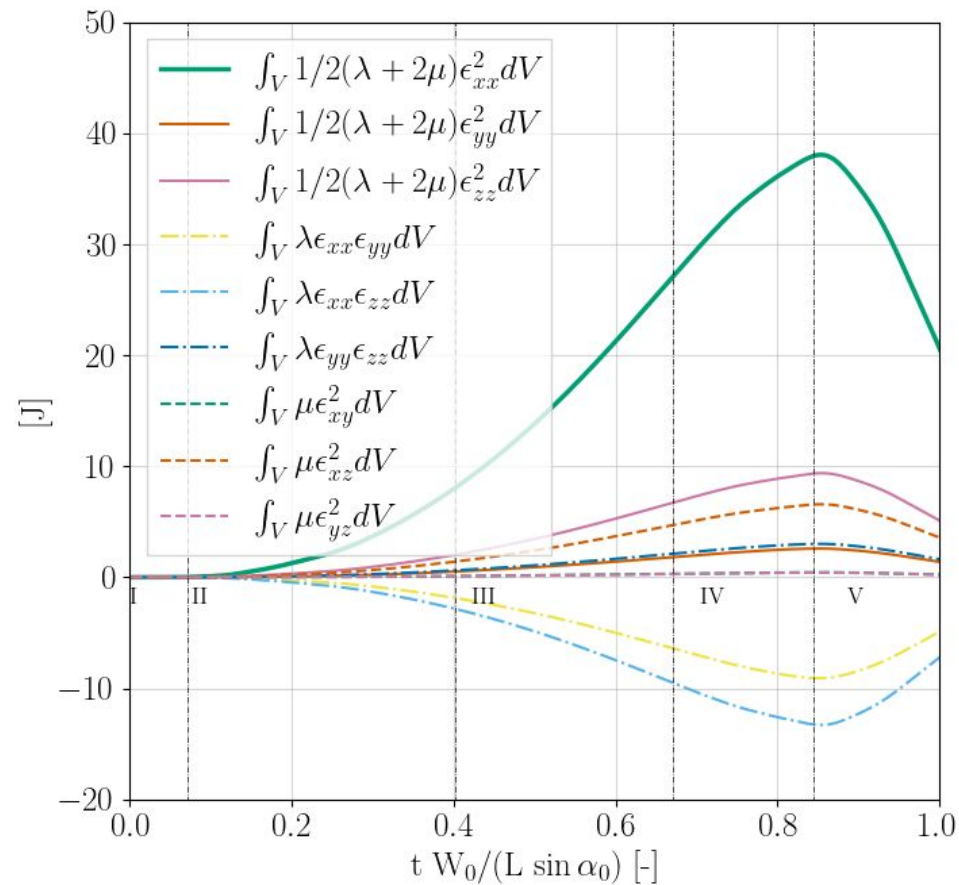
$$\delta_n = \boldsymbol{\delta} \cdot \hat{\mathbf{n}}$$



- The difference equation is satisfied
- The elastic potential energy almost equals the power provided by the pressure due to the flexibility
- The kinetic energy content is the smallest



- For all the cases the ϵ_{xx}^2 strain contribution is the most important.
- The second most important term is $\epsilon_{xx}\epsilon_{zz}$ since for the T-bars the ϵ_{zz} component is the most important (due to their orientation)



- A good correlation between Δp_{SR} and the elastic potential energy exists phase wise
- Δp_{SR} correlates well also with ΔW_{pR} and ΔW_M
- Δp_{SR} does not correlate well with the kinetic energy variation

